

Relativistic effects on LSS spectrum and bispectrum

Enea Di Dio

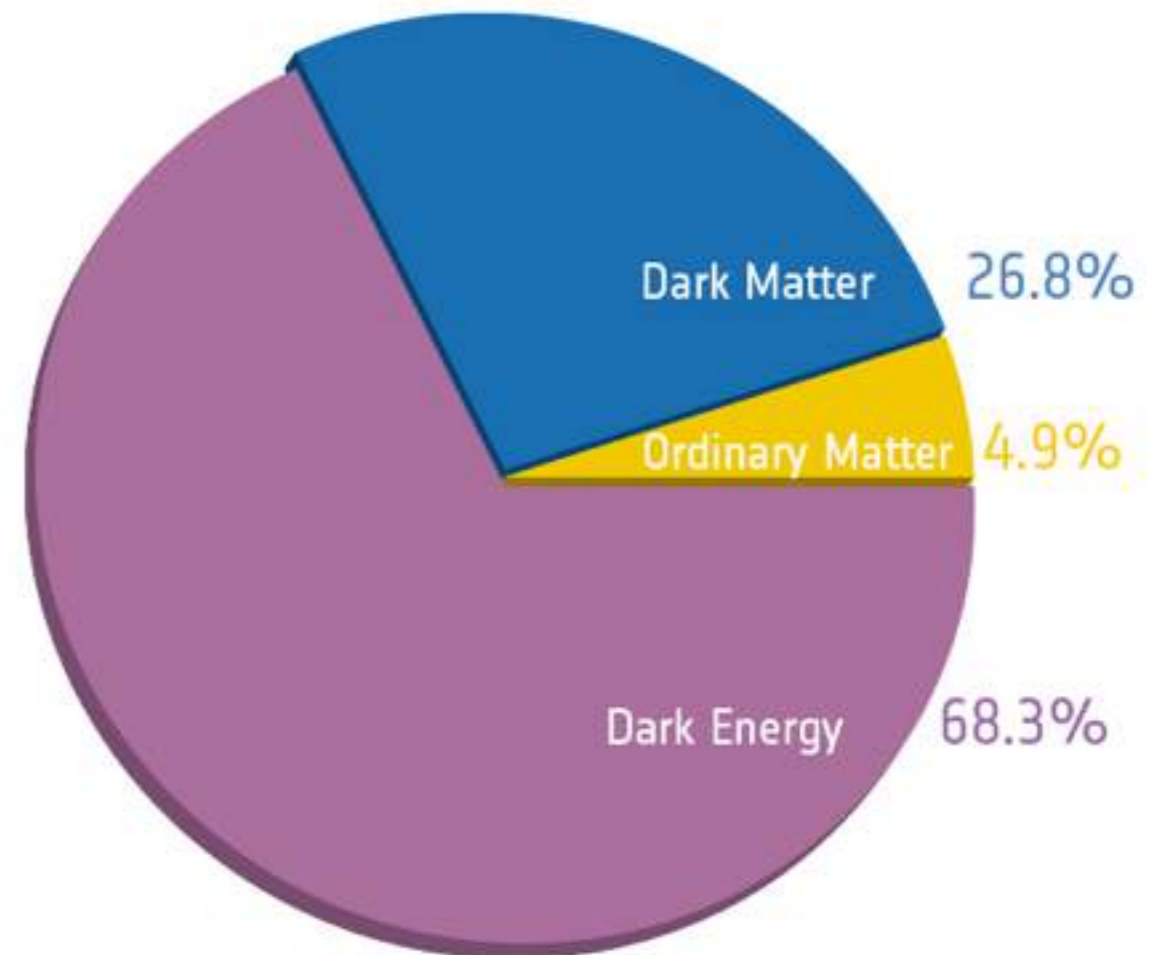
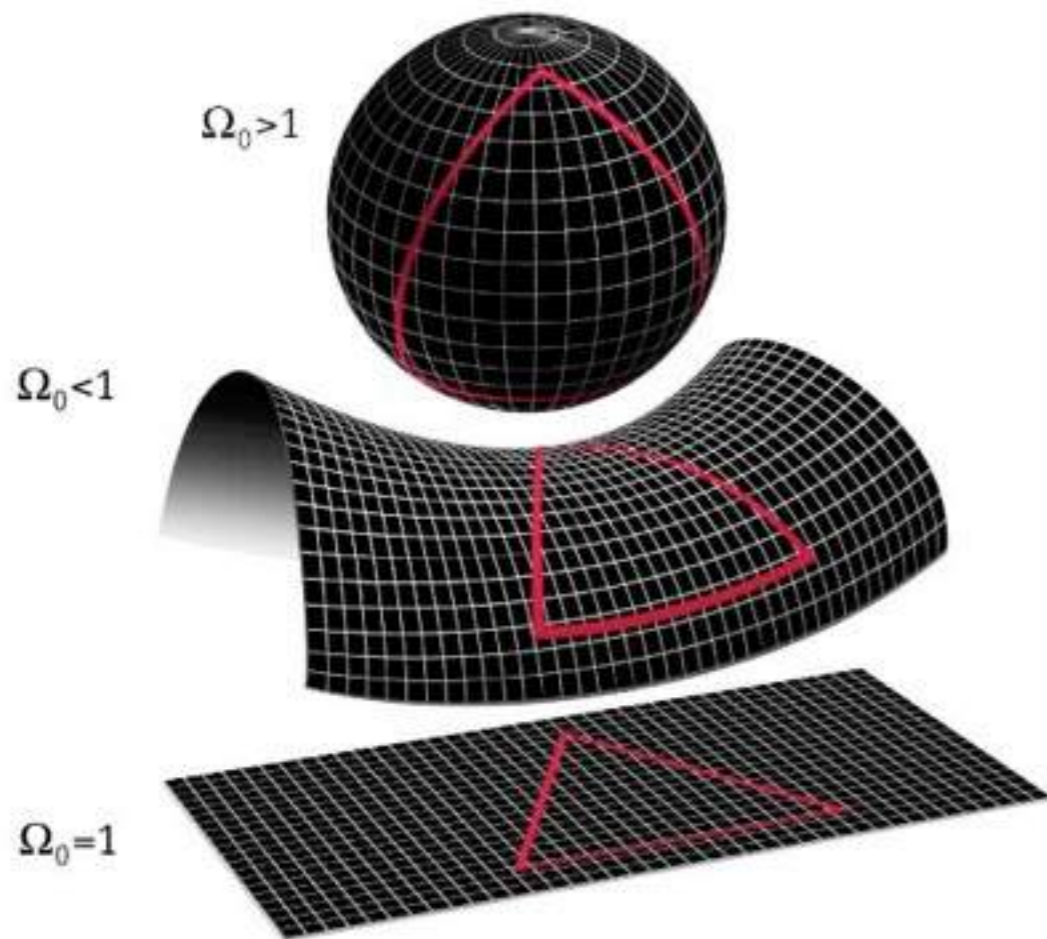
ED, Durrer, Marozzi, Montanari [arXiv:1407.0376]

ED, Durrer, Marozzi, Montanari [arXiv:1510.04202]

Irsic , ED, Viel [arXiv:1510.03436]

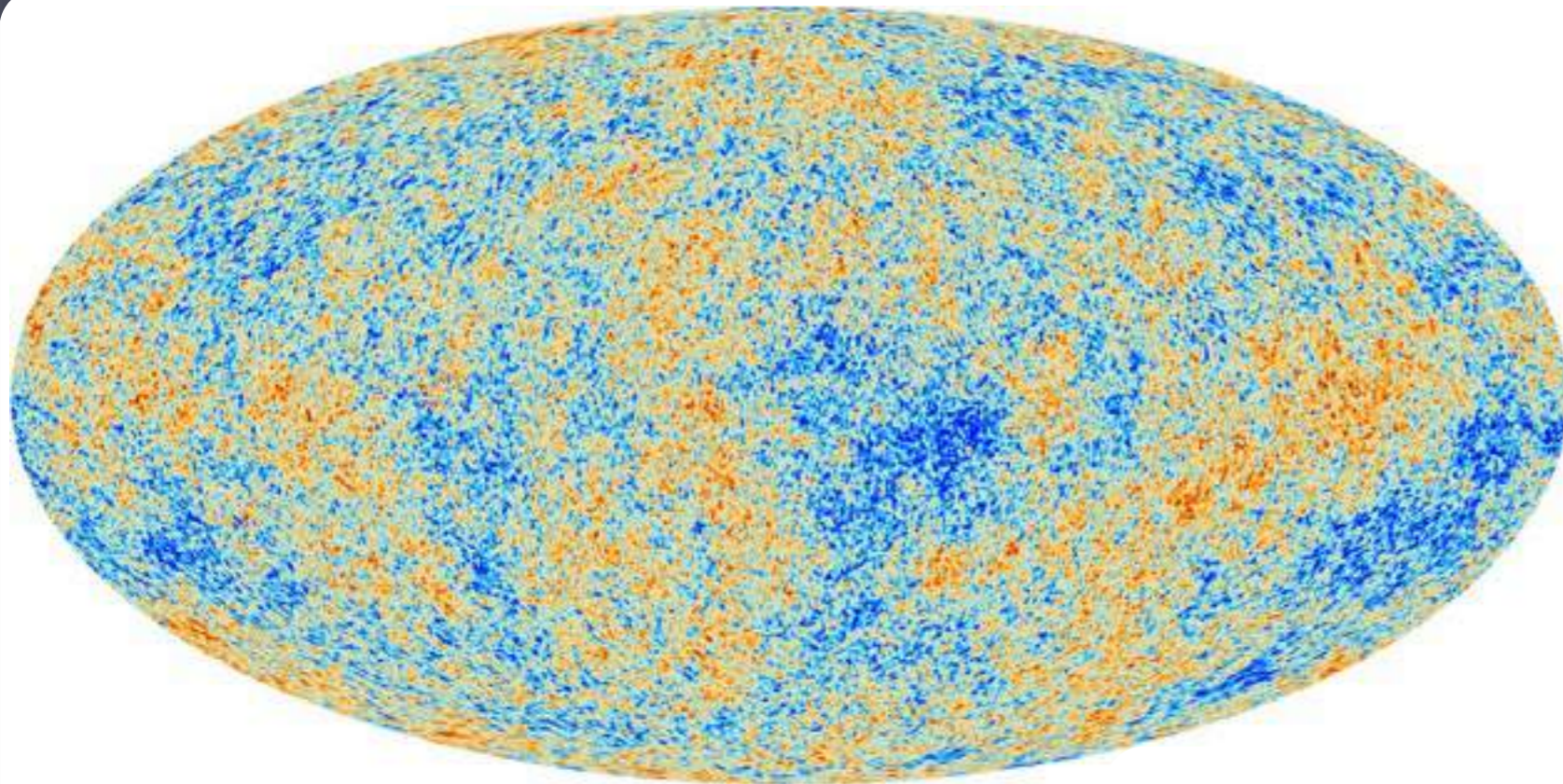
ED, Montanari, Raccanelli, Durrer, Kamionkowski, Lesgourgues [arXiv:1603.09073]

Standard Cosmological Model

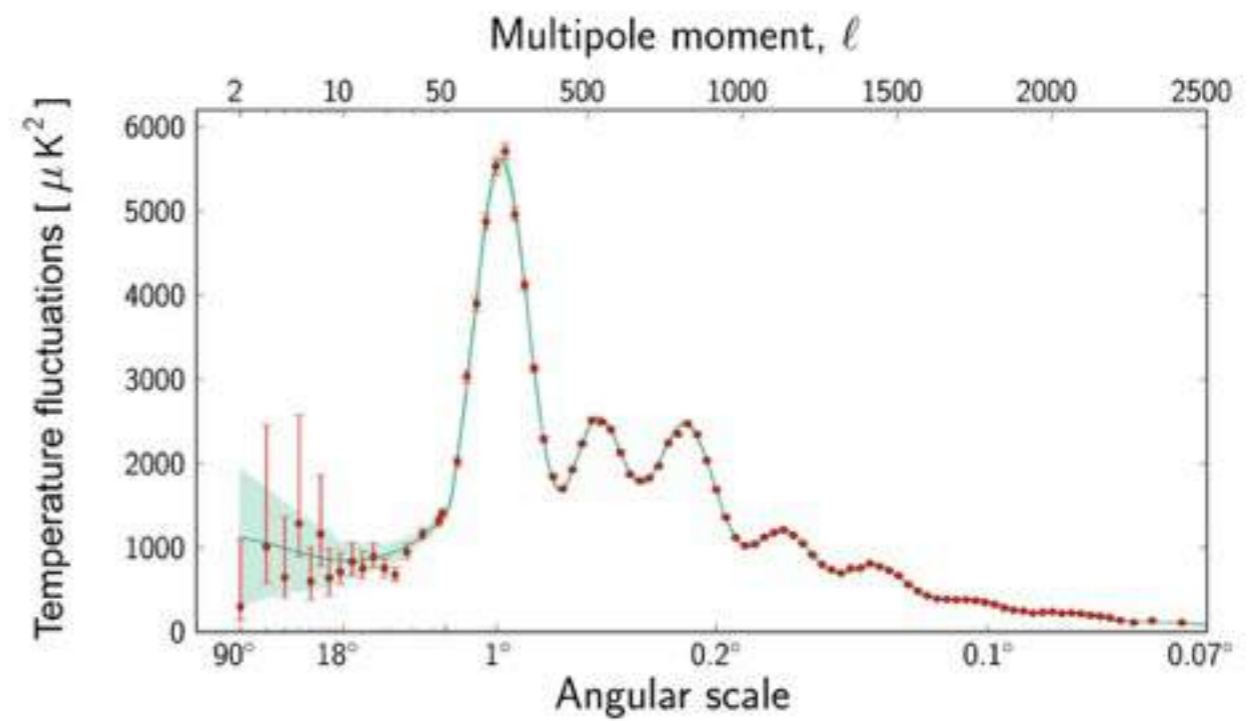


$$G_{\mu\nu} = \kappa T_{\mu\nu}$$

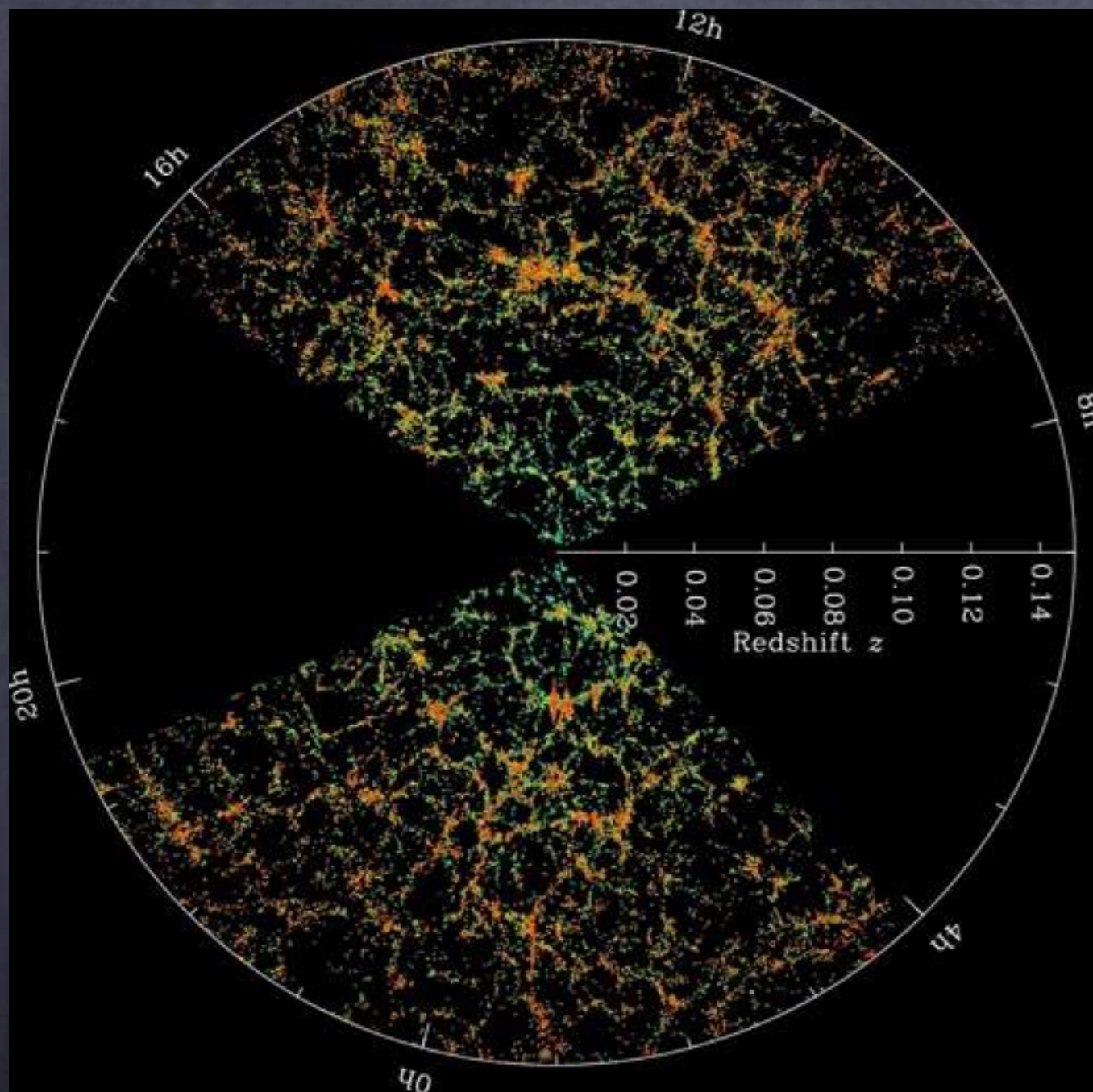
CMB



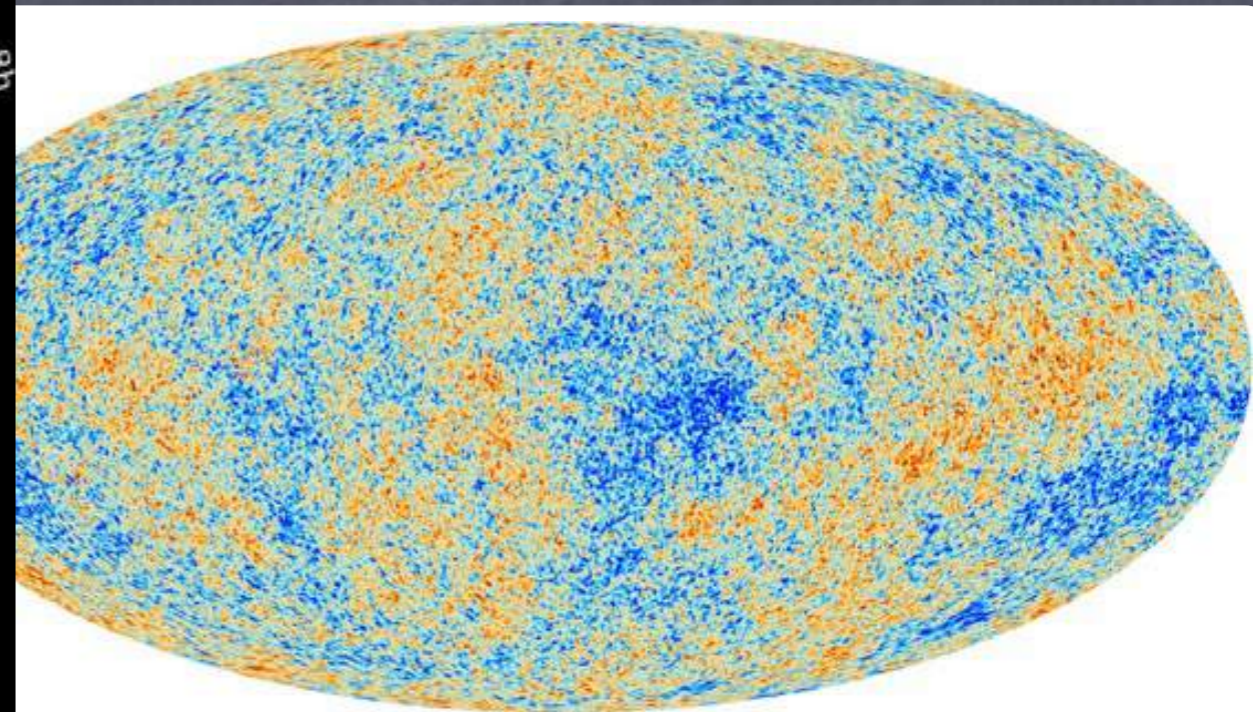
Planck Collaboration



Large Scale Structures



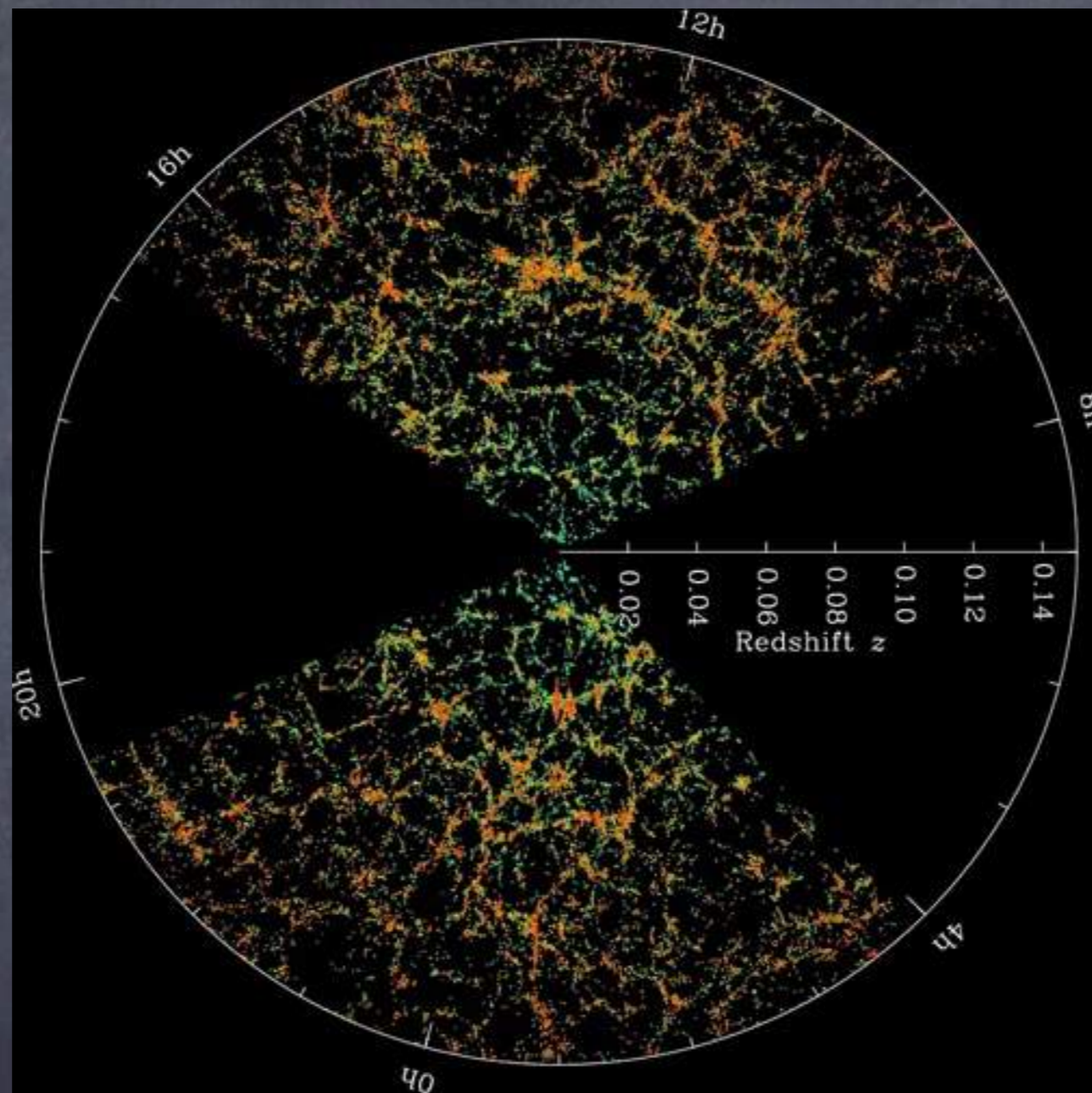
$$2 \sum_{\ell=2}^{2500} (2\ell + 1) \sim 10^7$$



Planck Collaboration

$$(3000)^3 = 2.7 \times 10^{10}$$

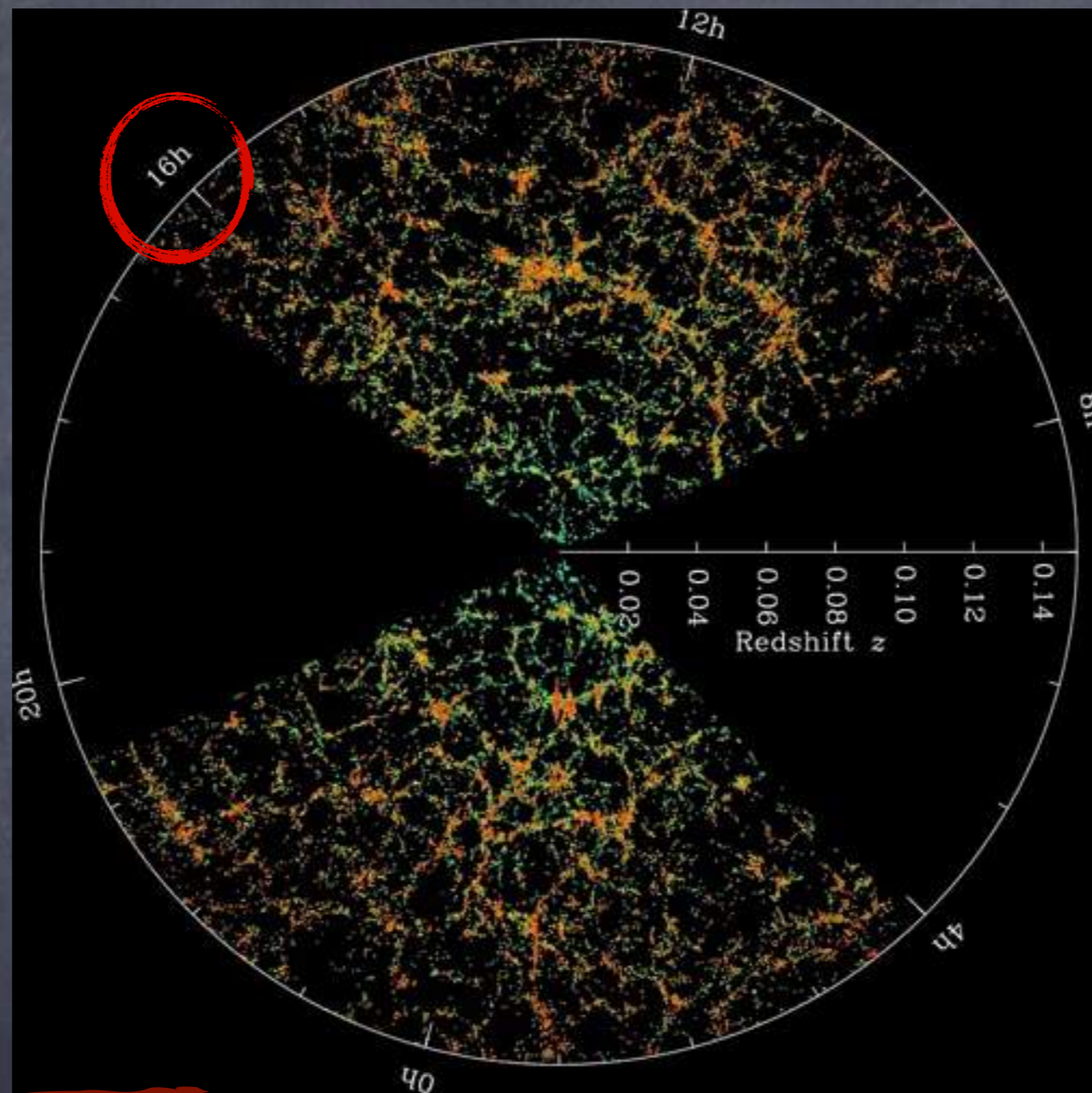
What do we really observe?



Angular position \mathbf{n} Redshift z

$$N(\mathbf{n}, z) d\Omega_{\mathbf{n}} dz$$

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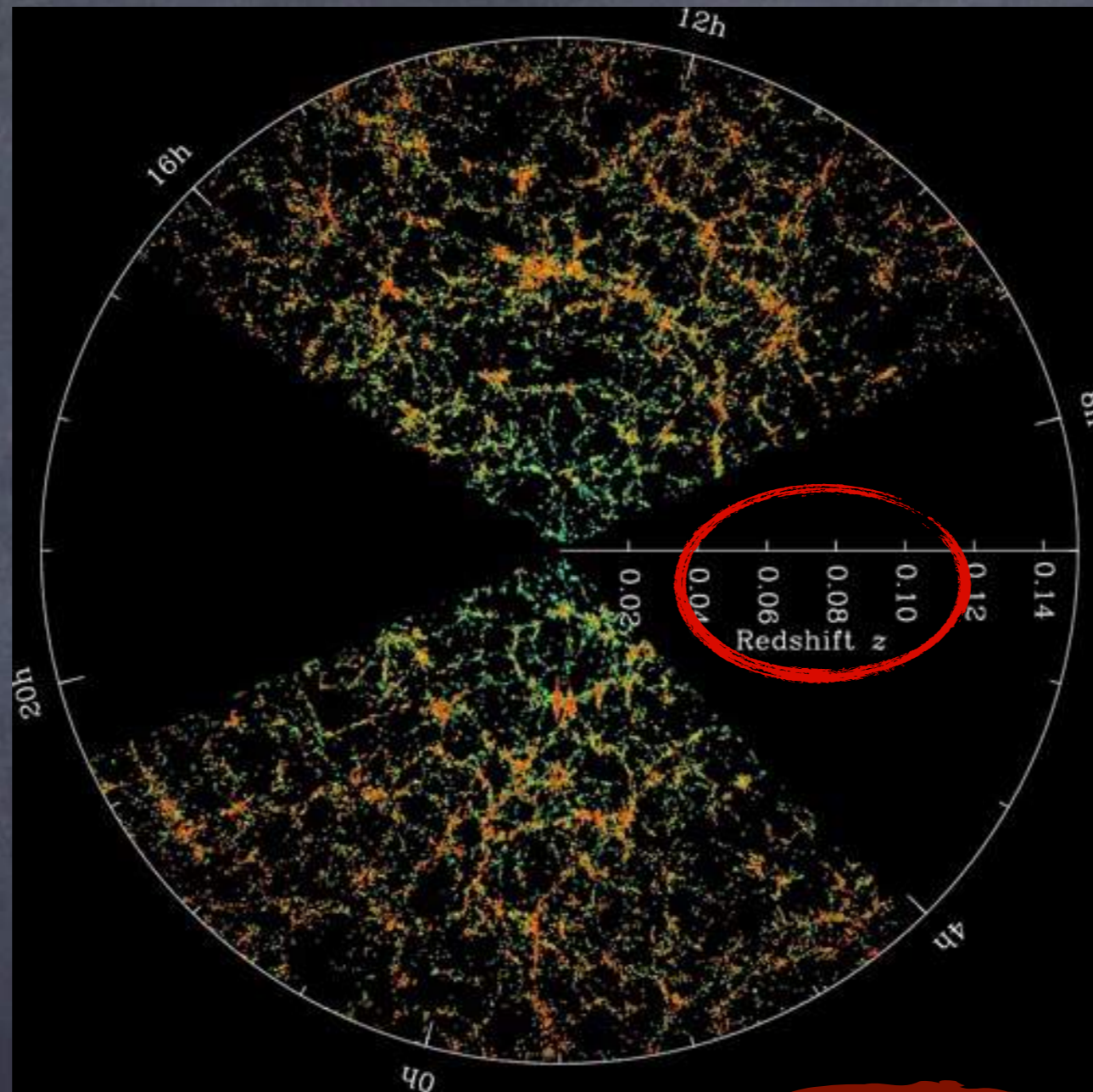


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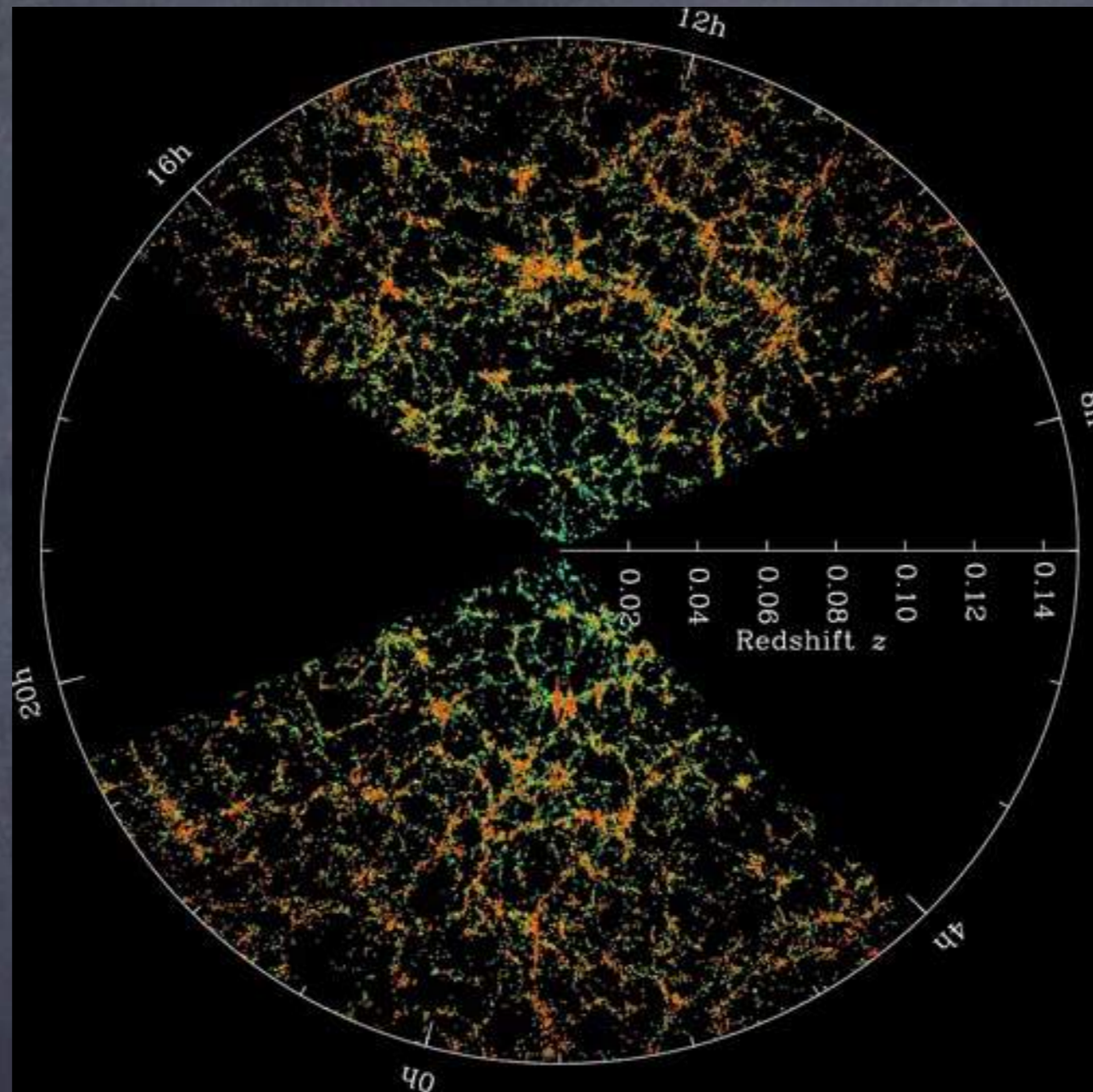


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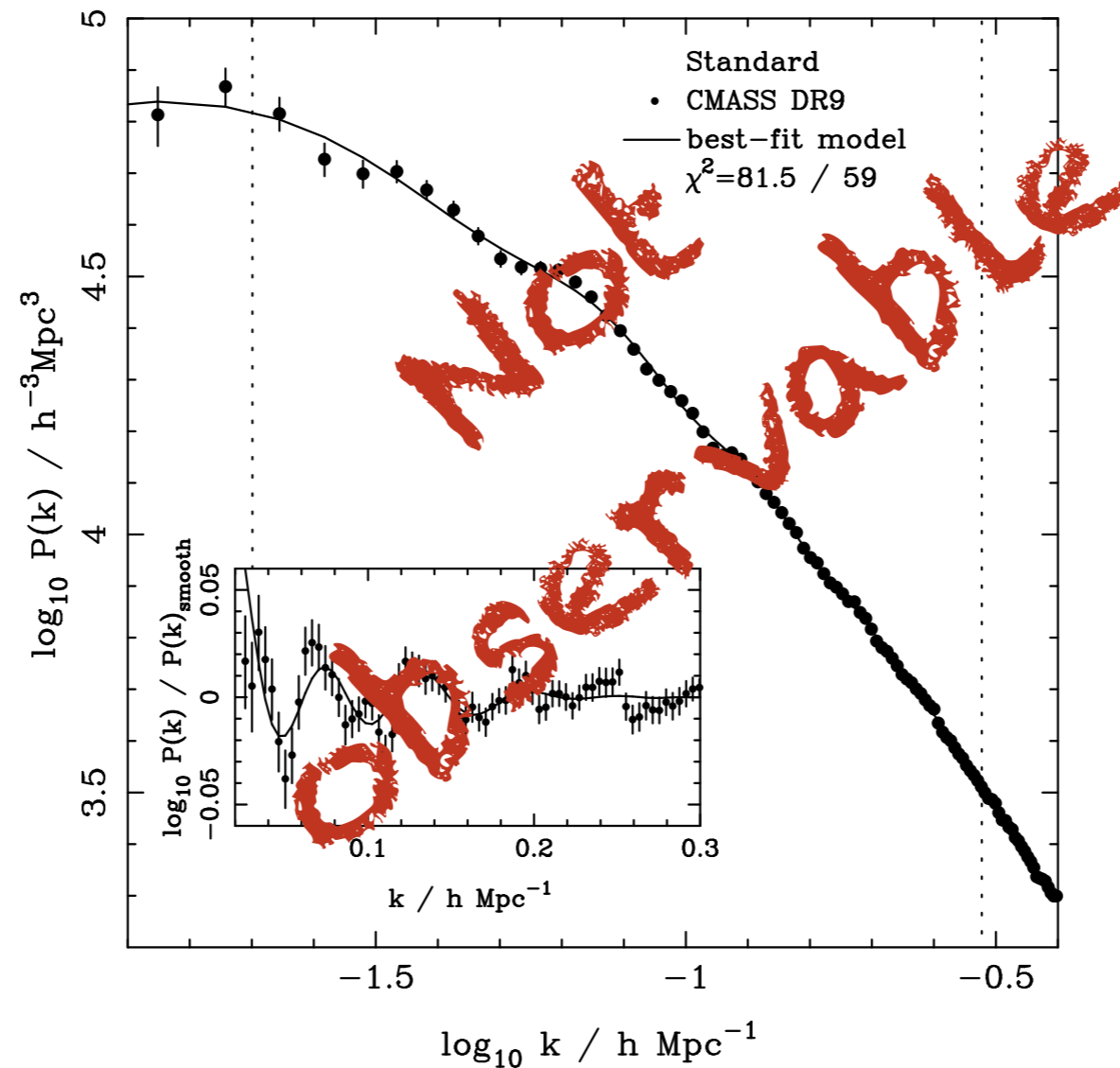


Info about
mass,
spectral type,
...

Angular position \mathbf{n} Redshift z

$$N(\mathbf{n}, z) d\Omega_{\mathbf{n}} dz$$

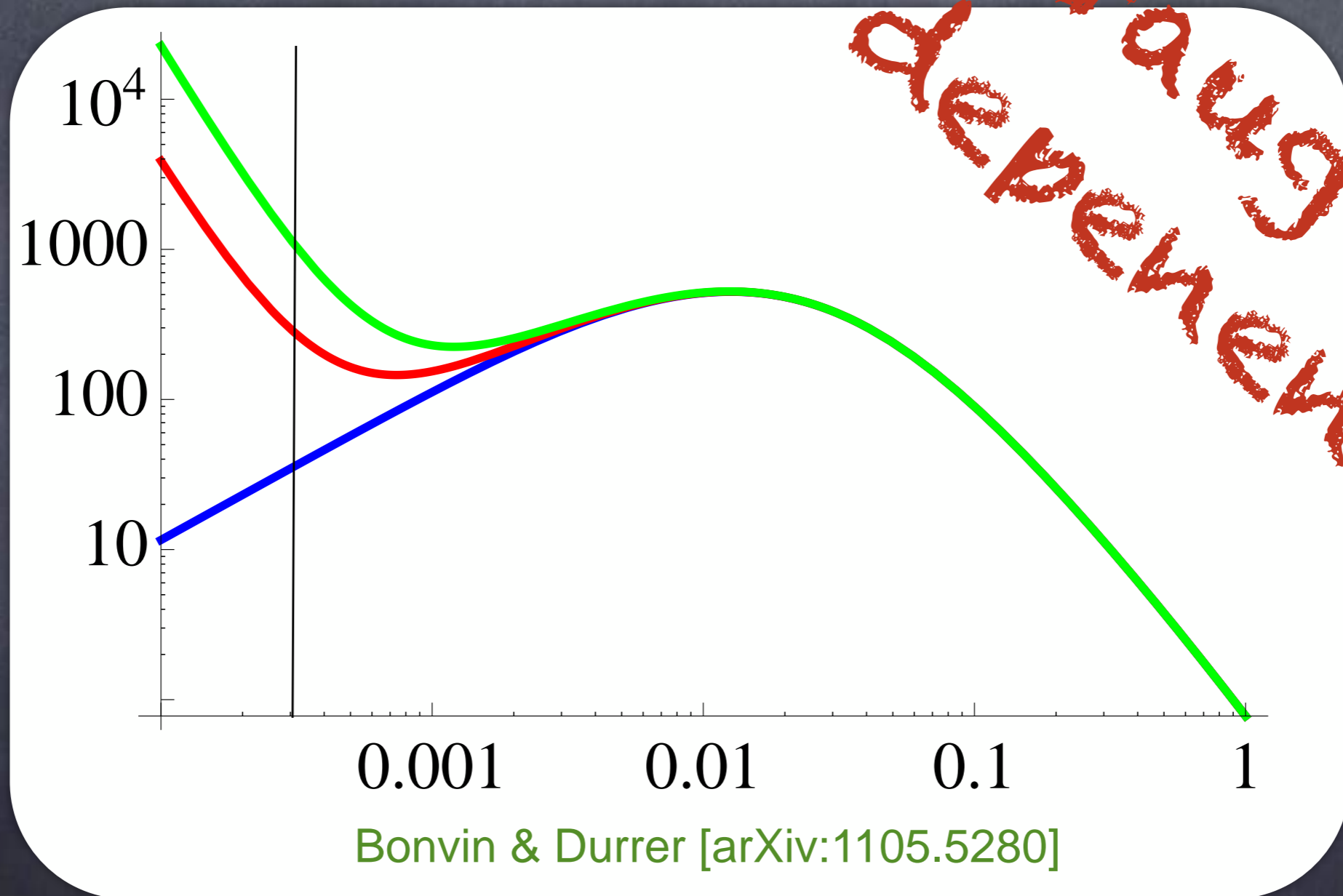
Large Scale Structures



Anderson et al '12 [arXiv:1203.6594]

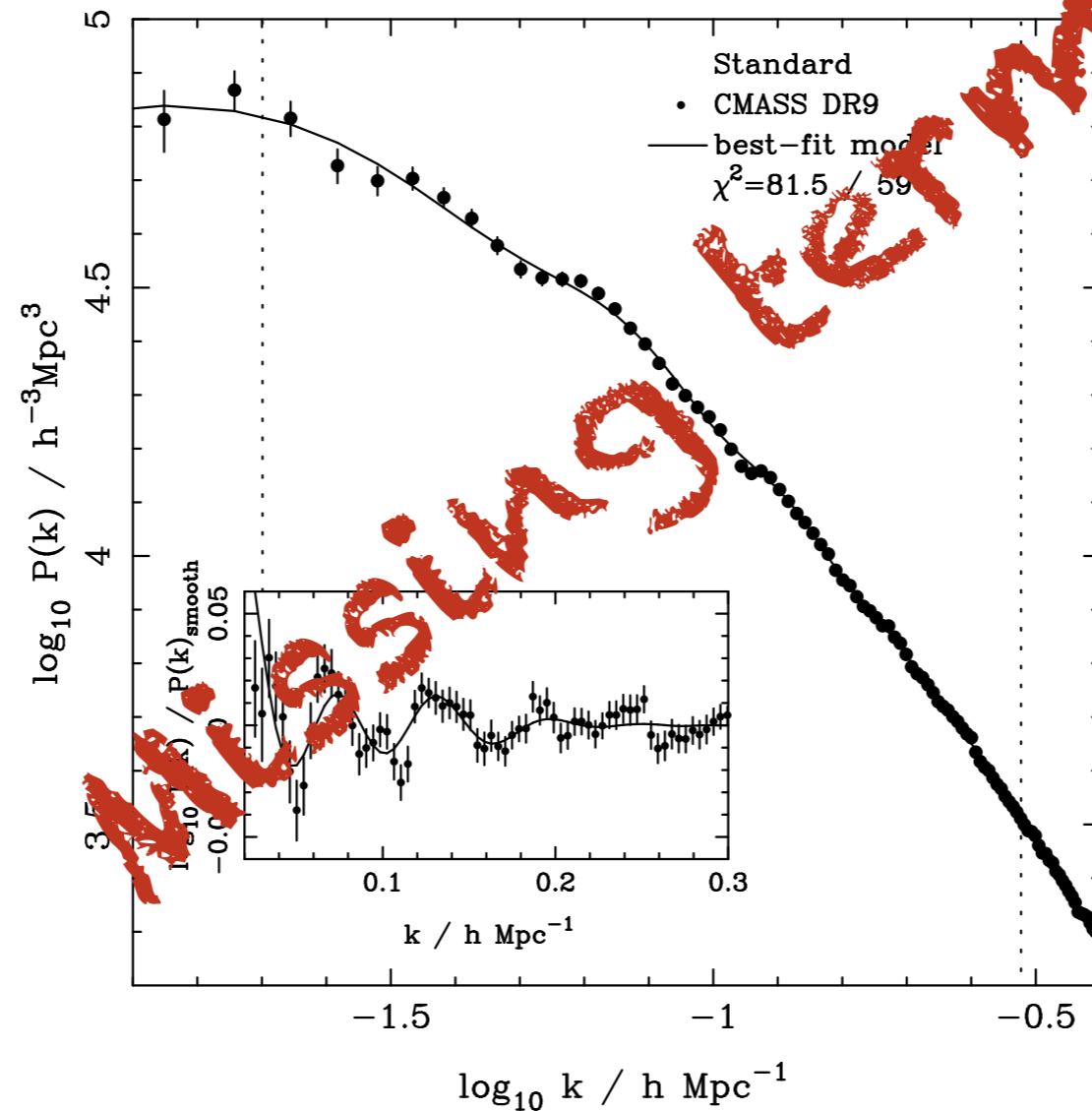
$$P_{\text{obs}}(k, \mu, z) = b(z)^2 (1 + \beta(z) \mu^2)^2 P(k, z)$$

Large Scale Structures



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Large Scale Structures



Anderson et al '12 [arXiv:1203.6594]

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Large Scale Structures

$$\Delta_N(\mathbf{n}, z, m_*) = b(z) D_{cm}(L > \bar{L}_*) + \mathcal{H}^{-1} \partial_r \mathbf{n} \cdot \mathbf{v} \quad \text{Standard}$$

$$\begin{aligned} & - \frac{2-5s}{2} \int_0^r \frac{r-r'}{rr'} \Delta_\Omega (\Psi + \Phi) dr' \\ & + \left(5s + \frac{2-5s}{\mathcal{H}r} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}}^N \right) \mathbf{n} \cdot \mathbf{v} + (f_{\text{evo}}^N - 3) \mathcal{H}v \\ & + \left(5s + \frac{2-5s}{\mathcal{H}r} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}}^N \right) \int_0^r (\dot{\Psi} + \dot{\Phi}) dr' \\ & + \left(5s + \frac{2-5s}{\mathcal{H}r} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}}^N \right) \Psi \\ & + \frac{2-5s}{r} \int_0^r (\Psi + \Phi) dr \\ & + (5s - 2) \Phi + \Psi + \mathcal{H}^{-1} \dot{\Phi} \end{aligned}$$

Relativistic
Effects

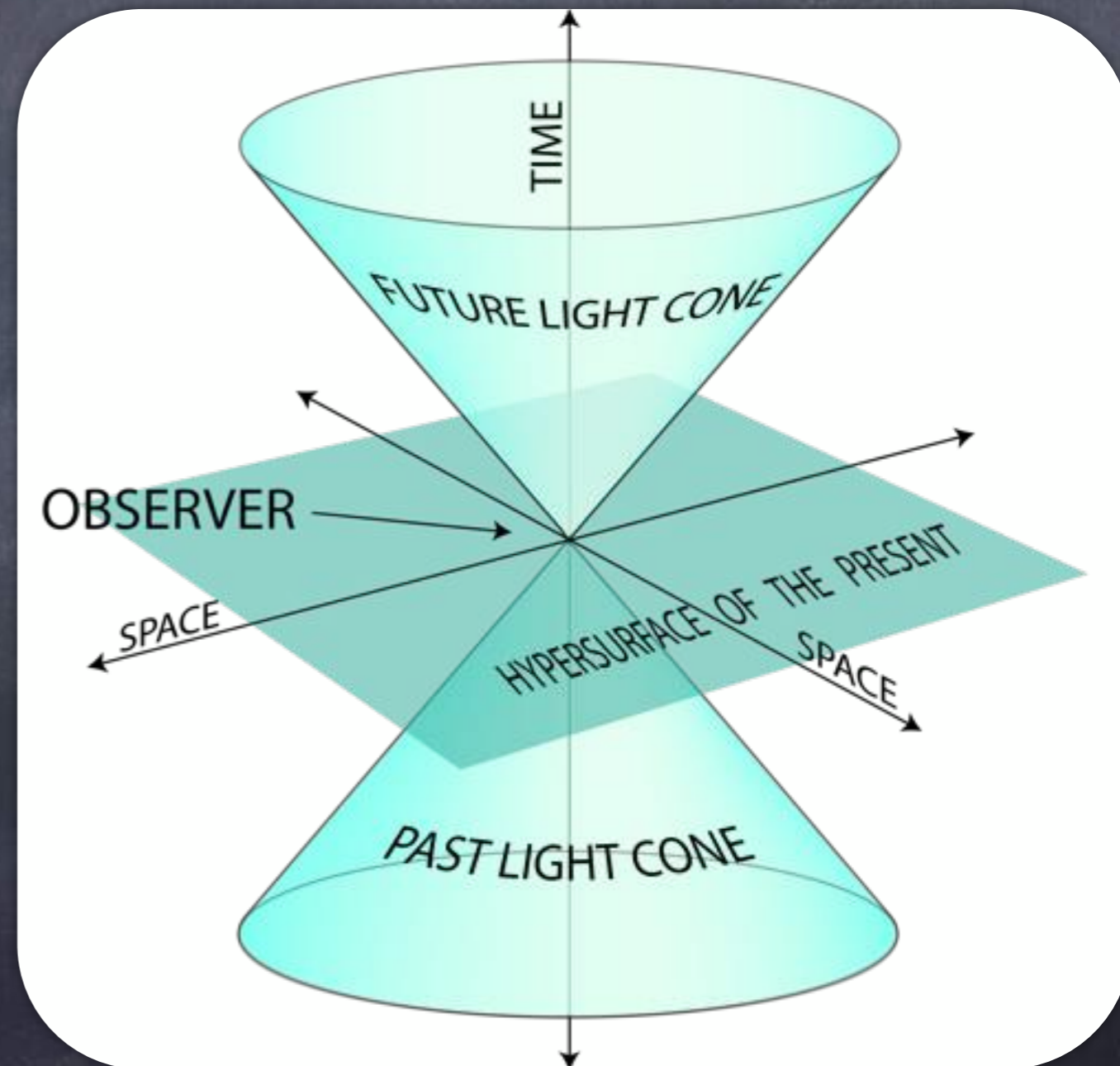
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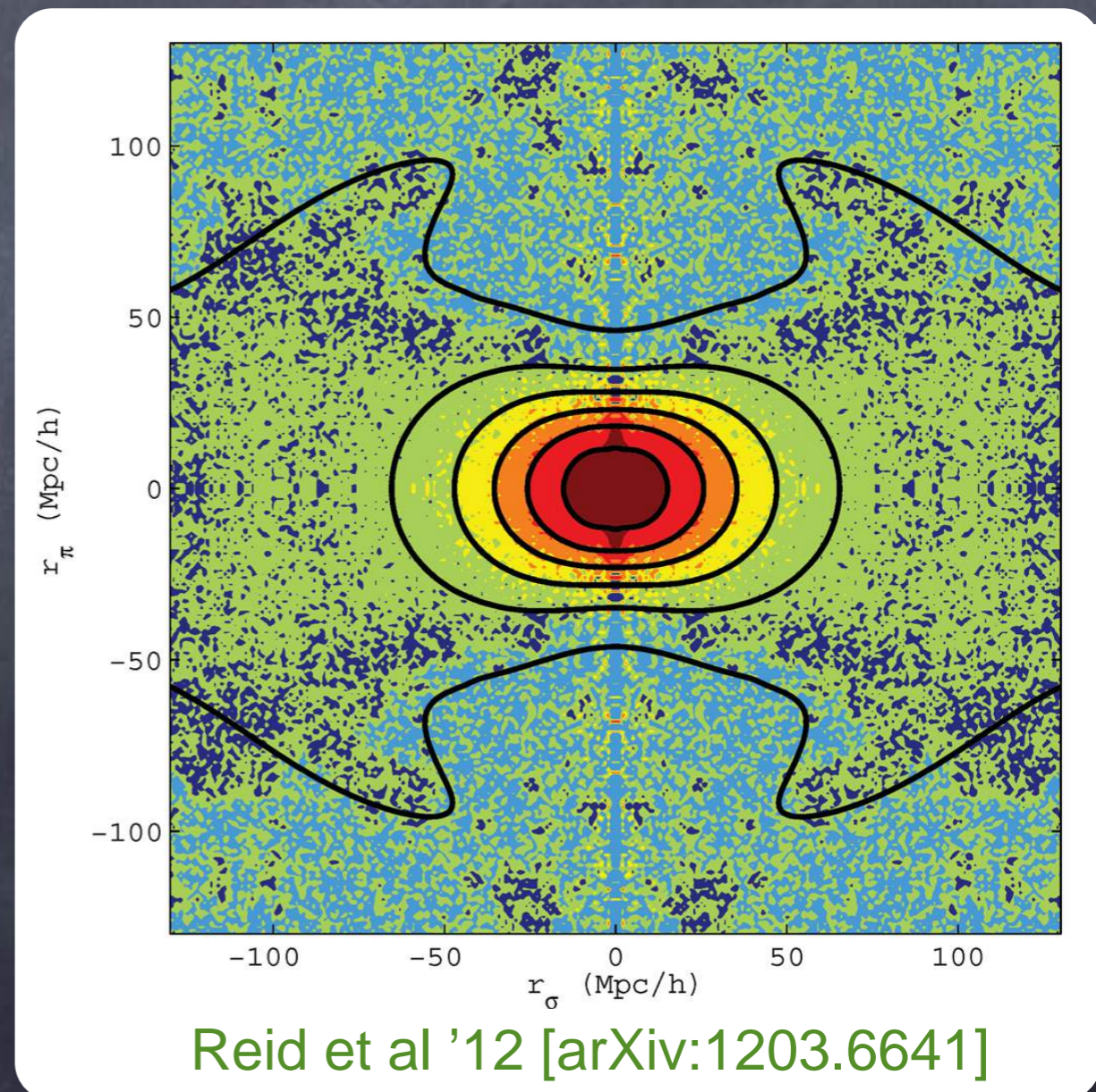
- observation on the past lightcone



What do we really observe?

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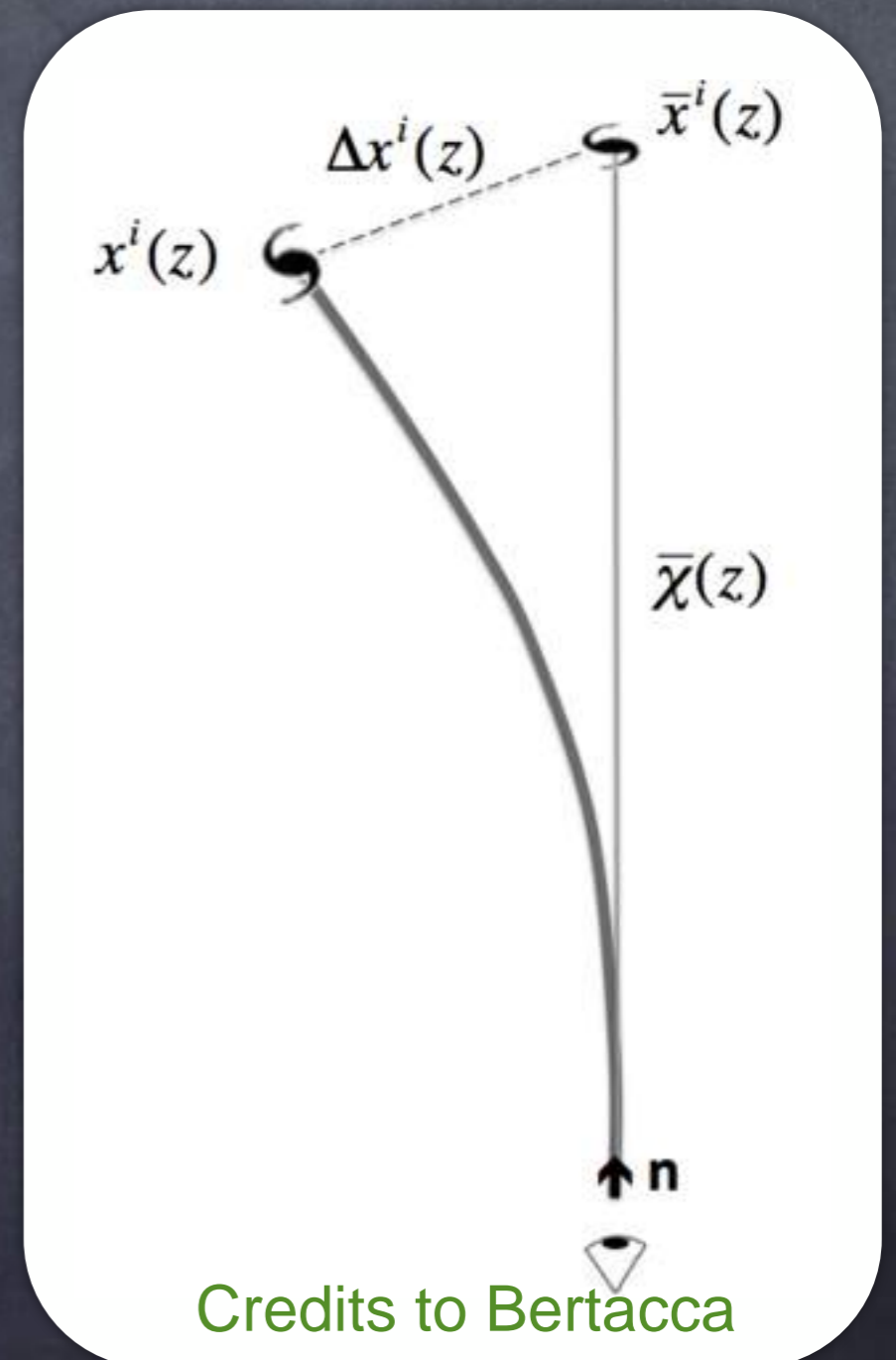
- observation on the past lightcone
- redshift perturbed by peculiar velocity



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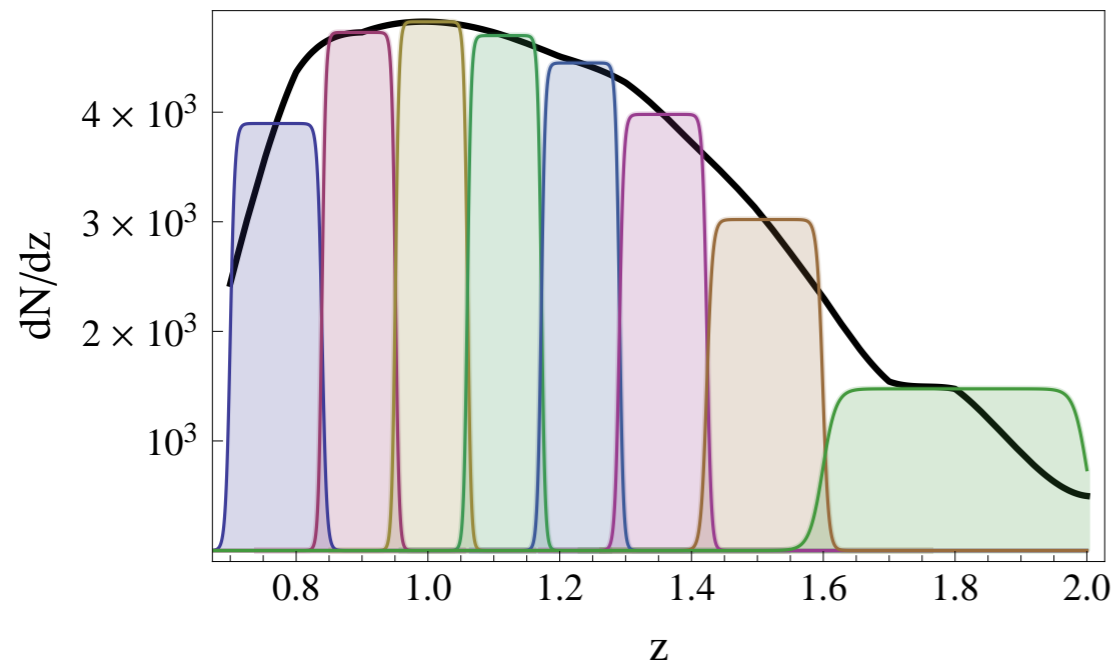


$c_\ell(z_1, z_2)$

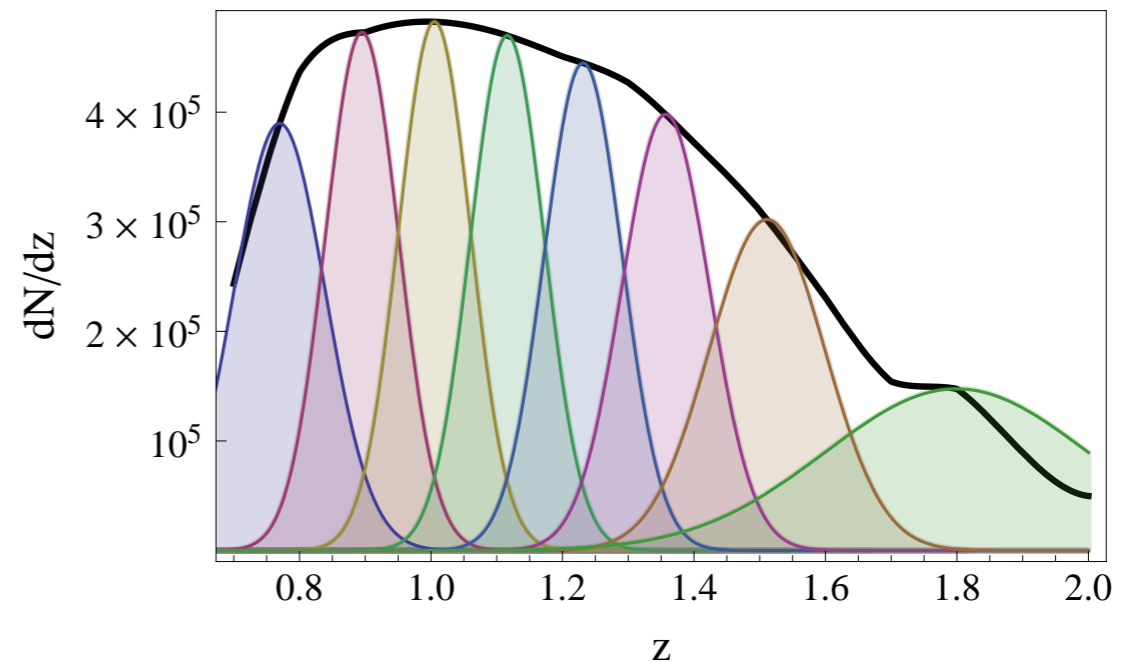
Binning Strategy

To recover the 3D information we need to split the redshift range in many bins and to consider the cross-correlations between different redshift bins.

Spectroscopic Survey



Photometric Survey

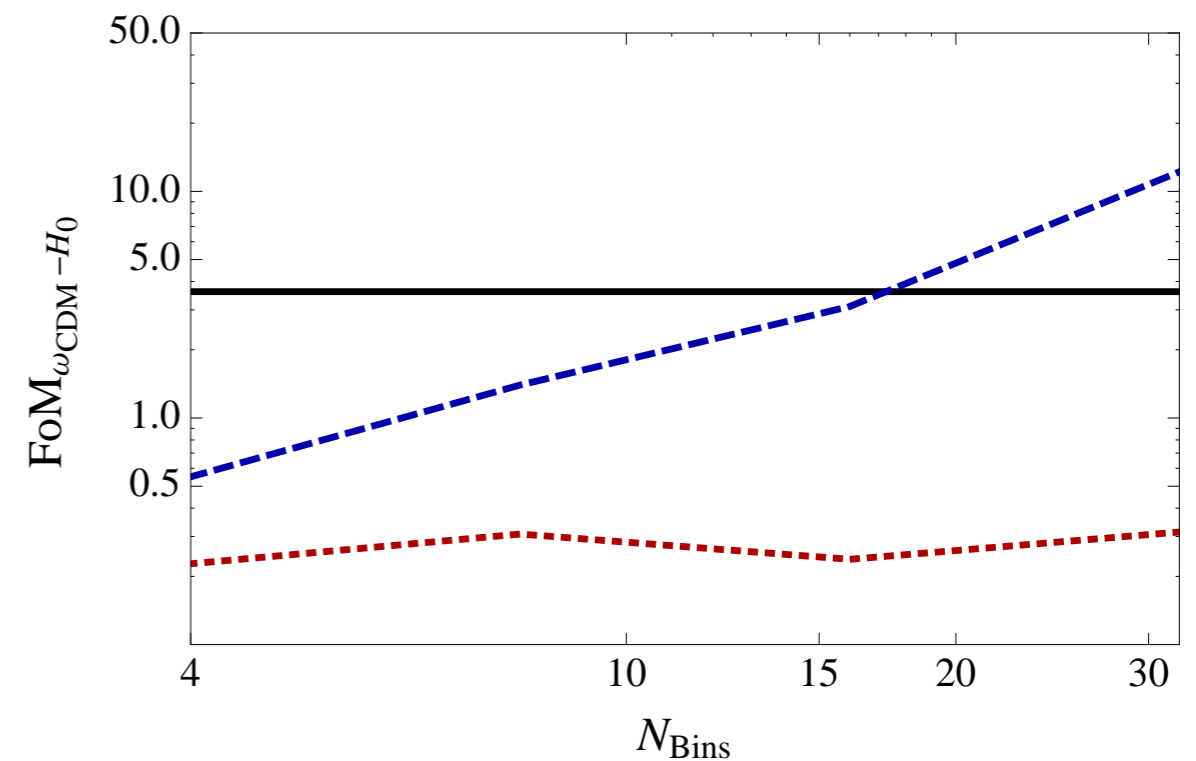
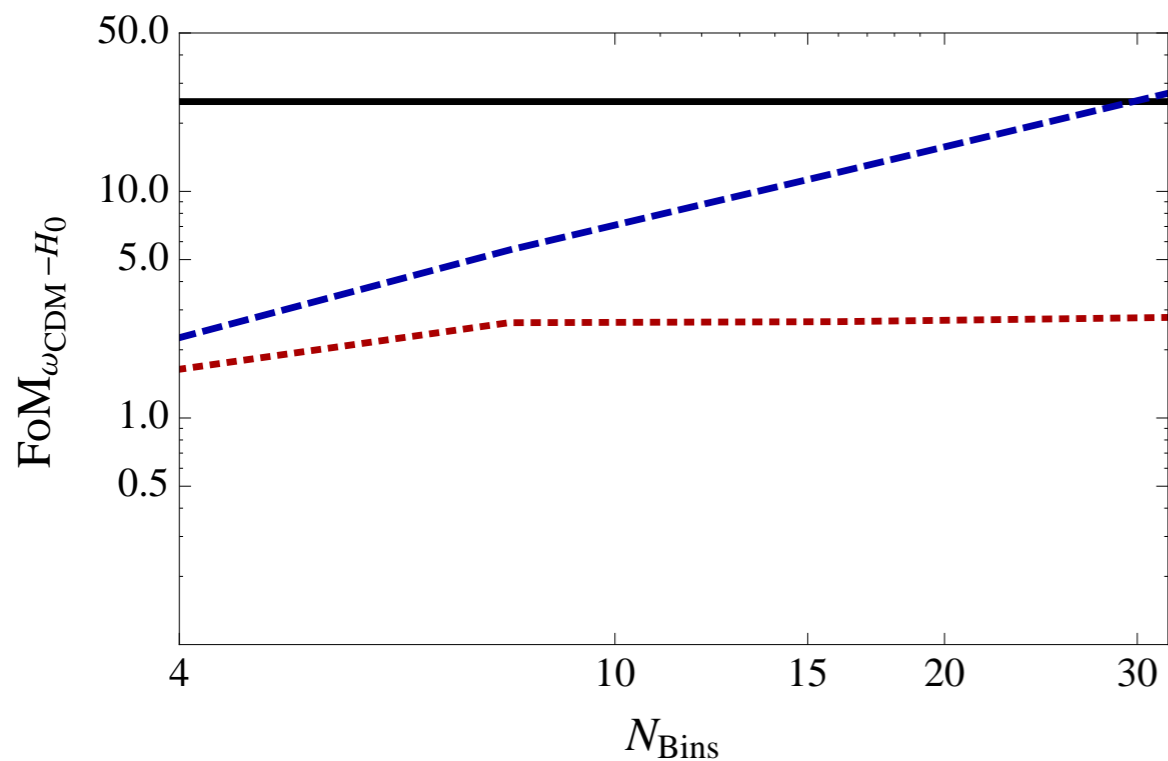


Cosmological Parameter Forecast

spectroscopic DES-like

$$\lambda_{\min} = 34 \text{ Mpc}/h$$

$$\lambda_{\min} = 68 \text{ Mpc}/h$$



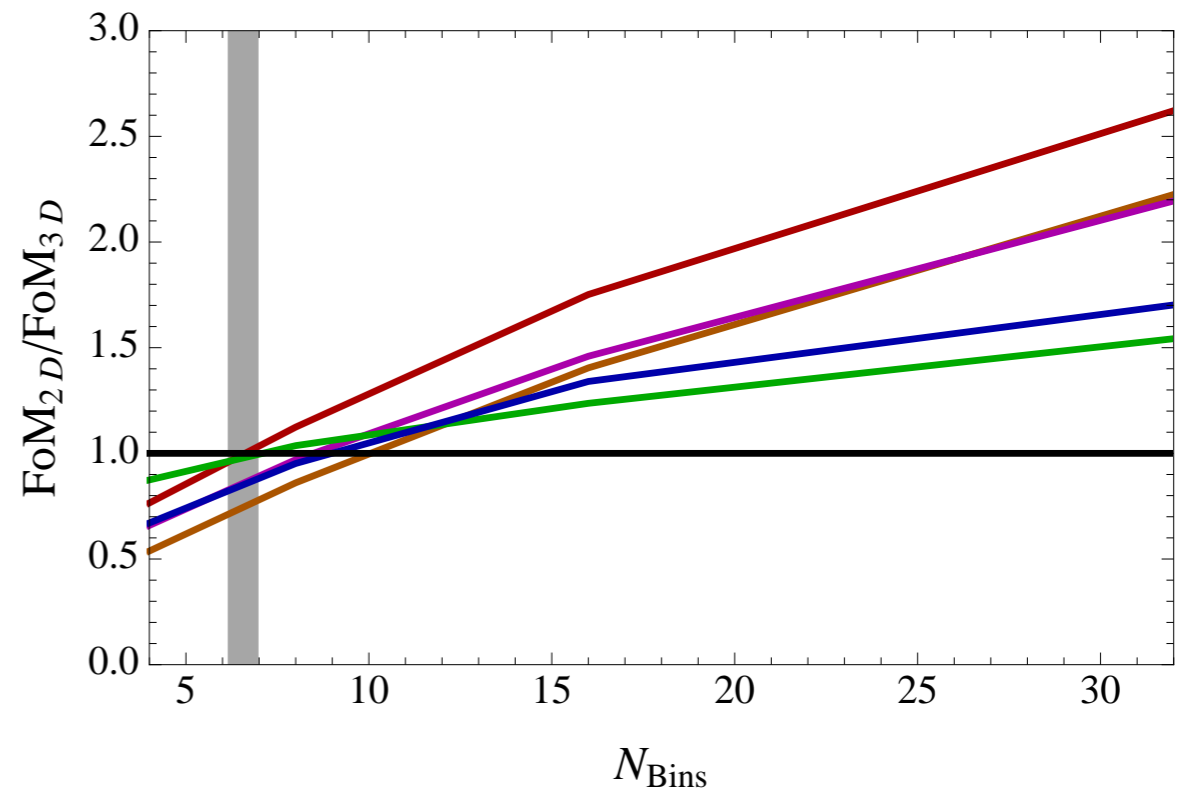
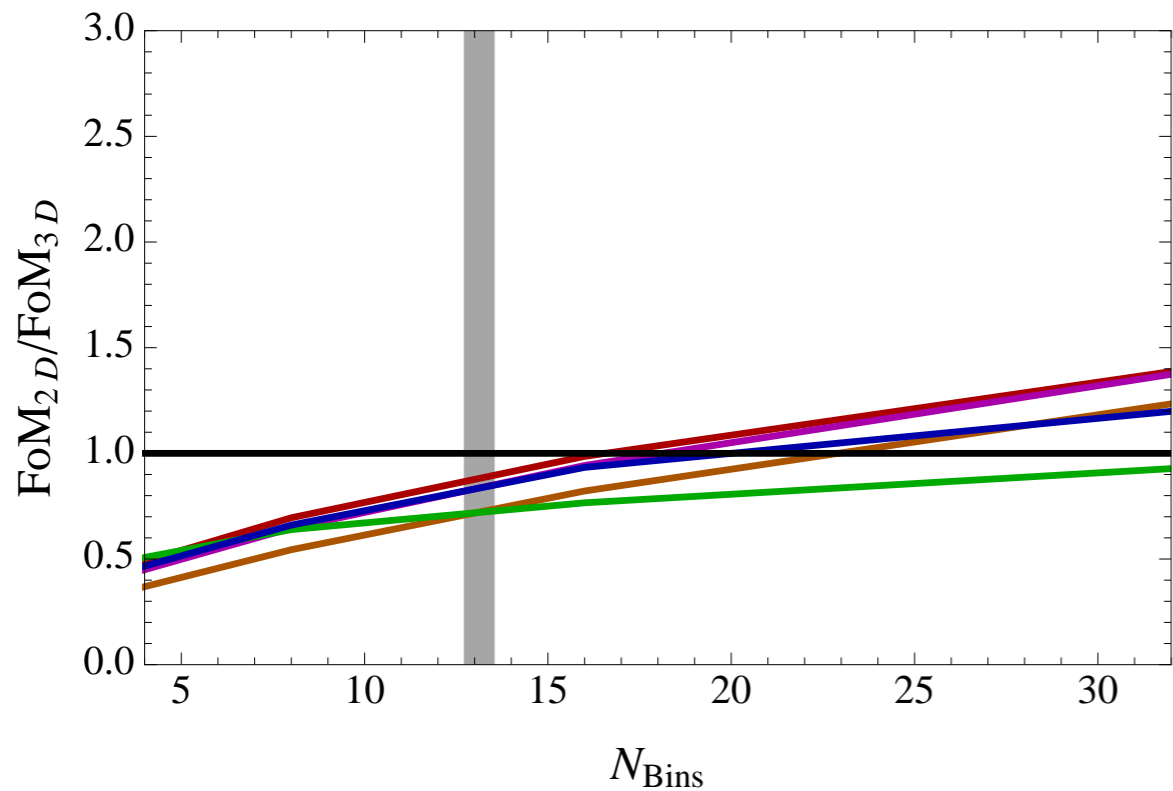
only redshift bins auto-correlations
with redshift bins cross-correlations

Cosmological Parameter Forecast

2D vs 3D

$$\lambda_{\min} = 34 \text{ Mpc}/h$$

$$\lambda_{\min} = 68 \text{ Mpc}/h$$



ω_b ω_{CDM} n_s H_0 A_s

Large Scale Structures

$$\begin{aligned}
 \Delta_N(\mathbf{n}, z, m_*) &= b(z) D_{cm}(L > \bar{L}_*) + \mathcal{H}^{-1} \partial_r \mathbf{n} \cdot \mathbf{v} \\
 &\sim D \\
 &- \frac{2-5s}{2} \int_0^r \frac{r-r'}{rr'} \Delta_\Omega(\Psi + \Phi) dr' \\
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 \end{aligned}$$

Bonvin & Durrer [arXiv:1105.5280],
 Challinor & Lewis [arXiv:1105.5292],
 Yoo [arXiv:1009.3021]

Large Scale Structures

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$$\sim \frac{\mathcal{H}}{k} D$$

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$\sim \left(\frac{\mathcal{H}}{k} \right)^2 D$

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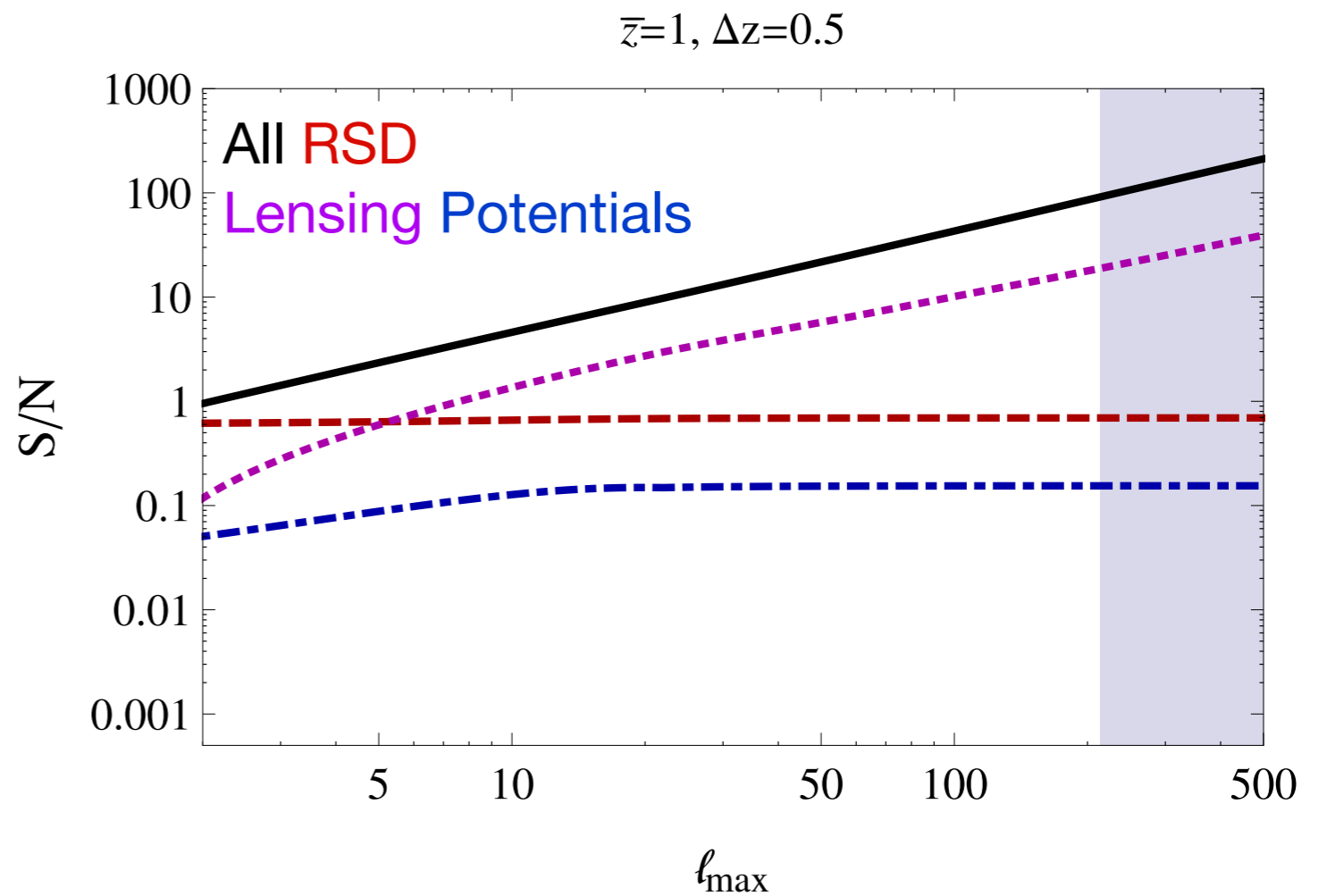
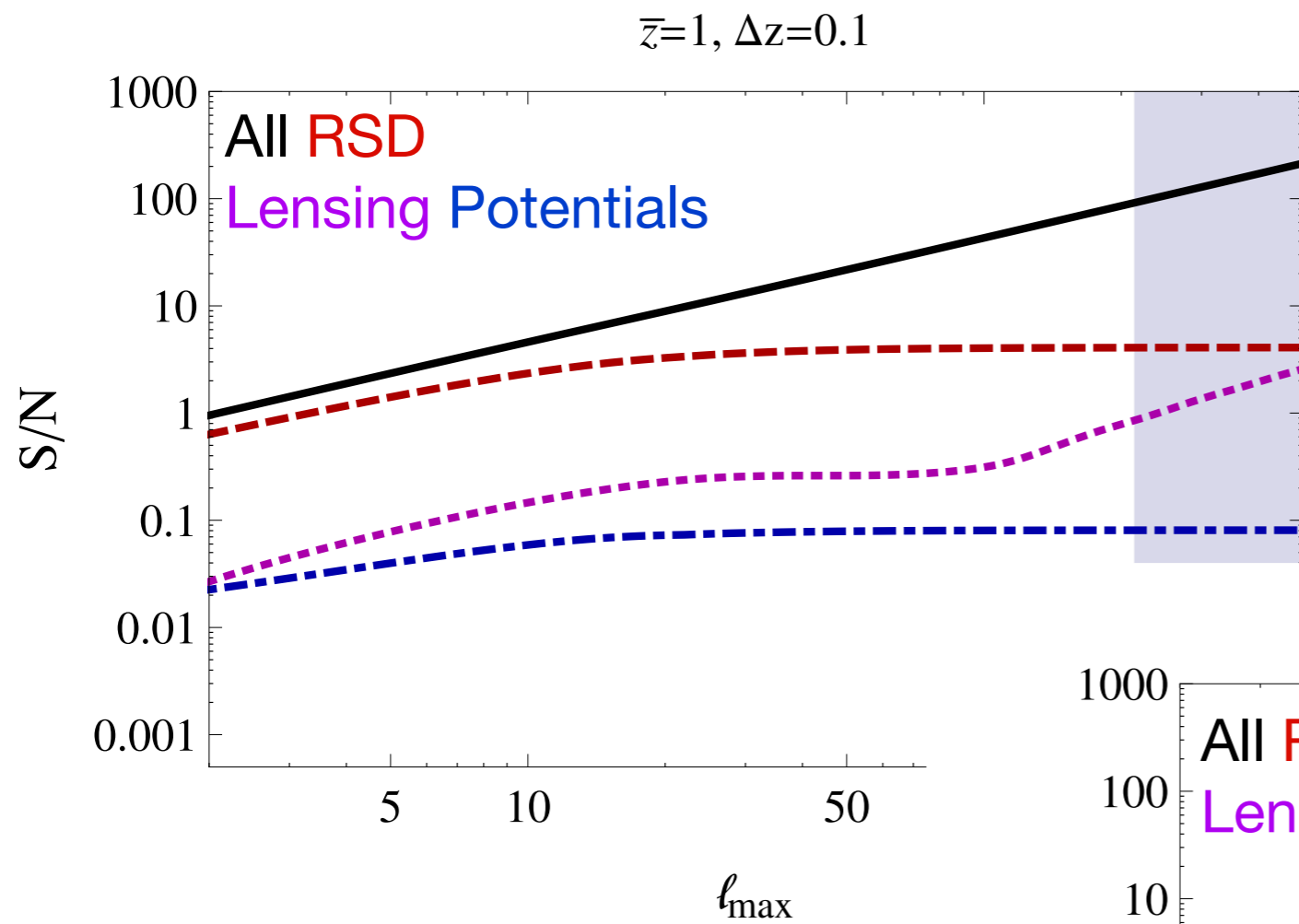
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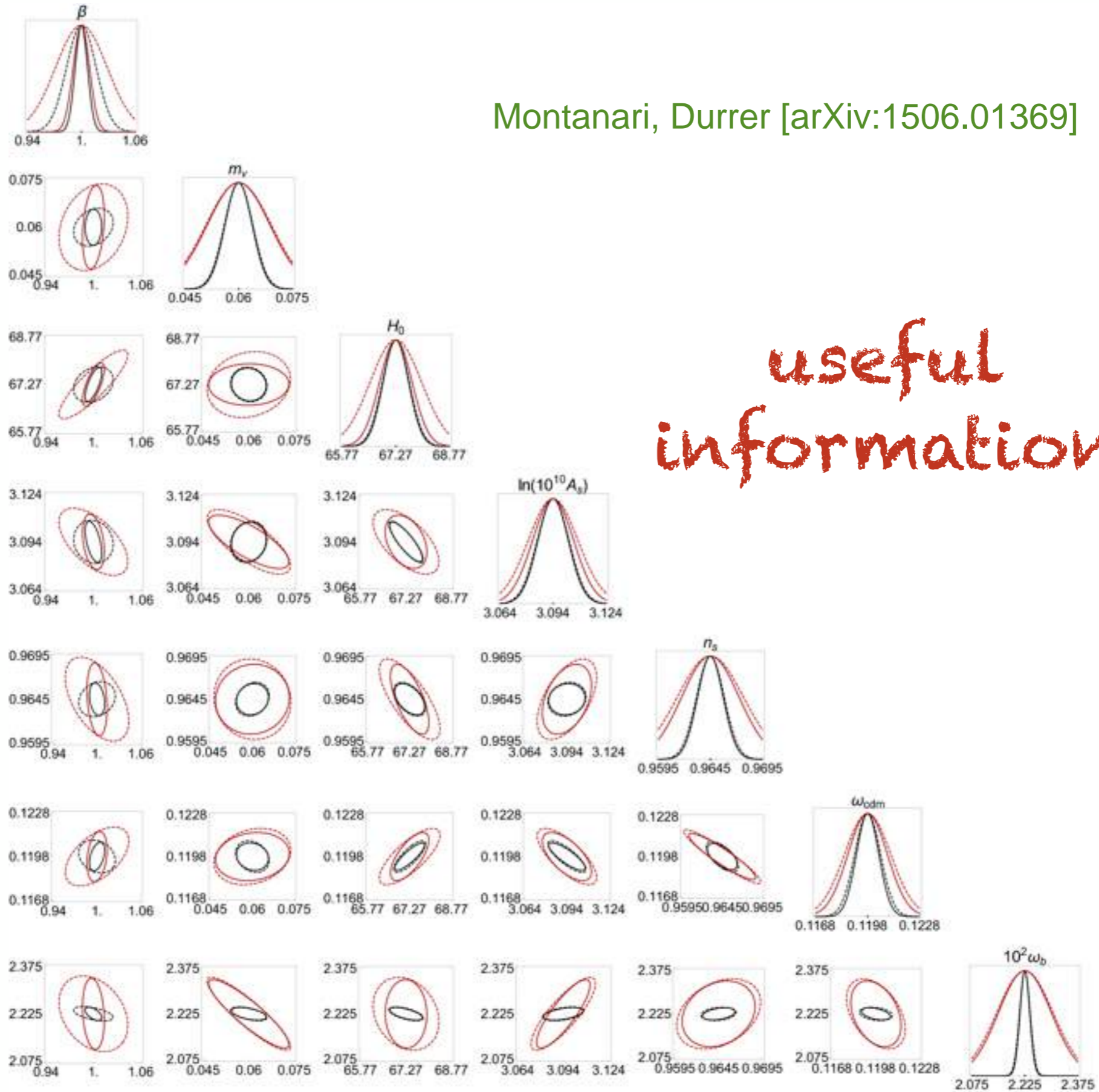
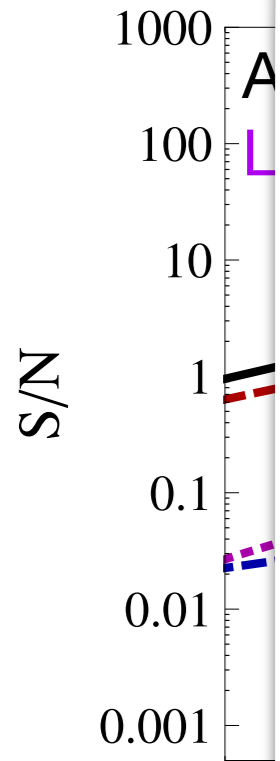
Lensing Potential



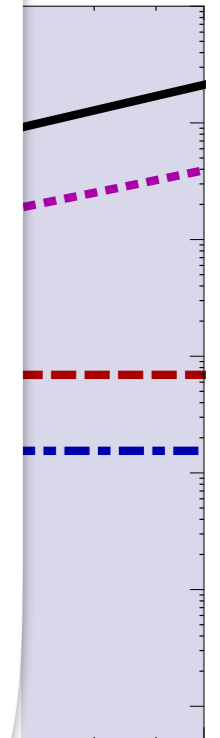
ED, Montanari, Durrer, Lesgourgues,
[arXiv:1308.6186]

Montanari, Durrer [arXiv:1506.01369]

useful
information

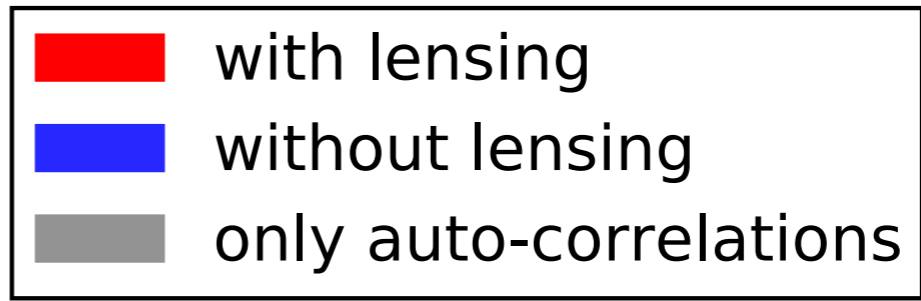


ED, Montanari
[arXiv:1308.0001]



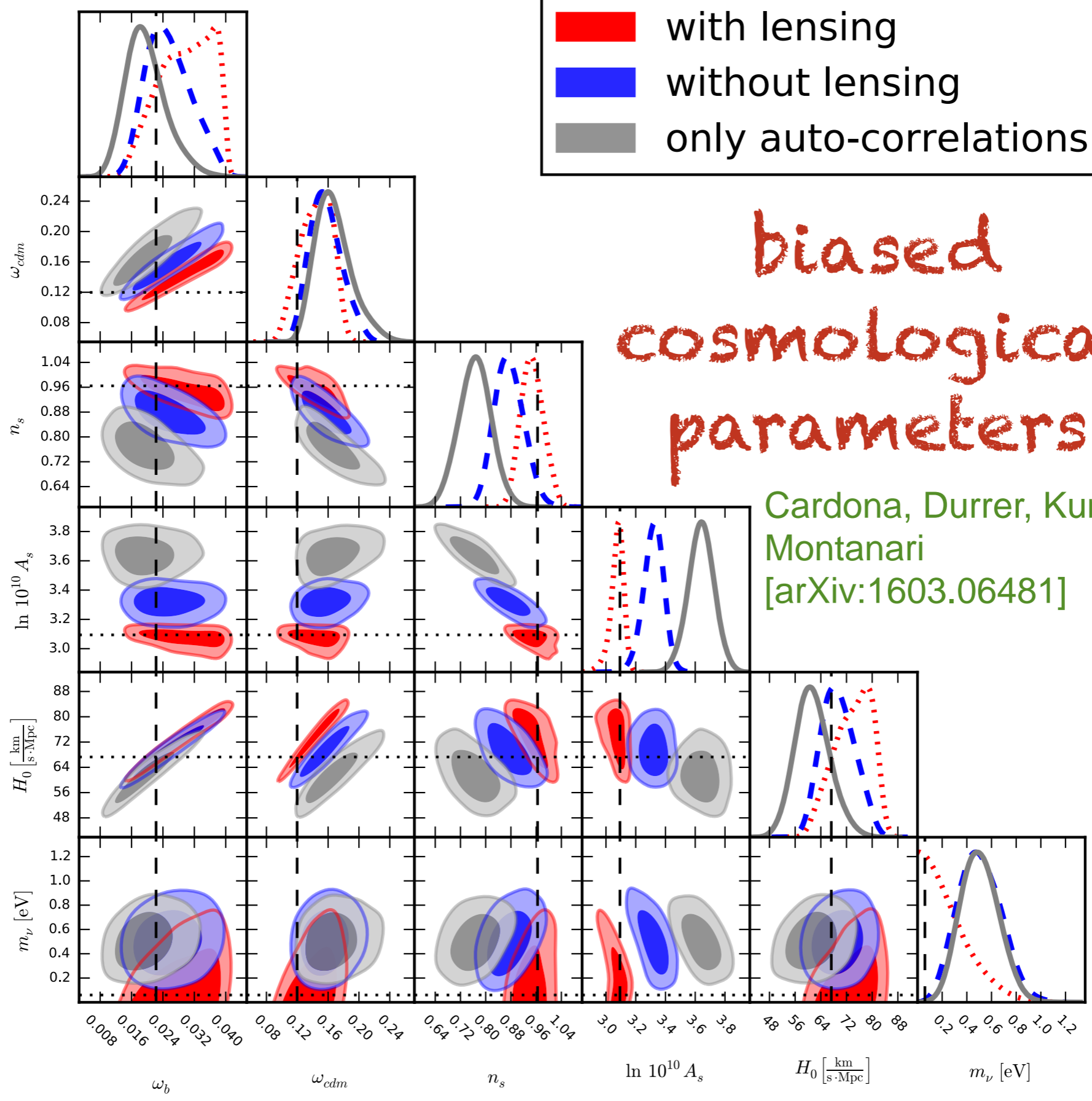
S/N

1000
100
10
1
0.1
0.01
0.001

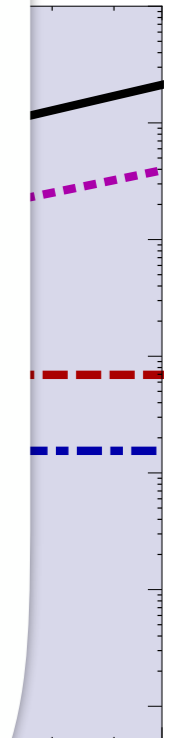


biased
cosmological
parameters

Cardona, Durrer, Kunz,
Montanari
[arXiv:1603.06481]



ED, Mont
[arXiv:13



Large Scale Structures

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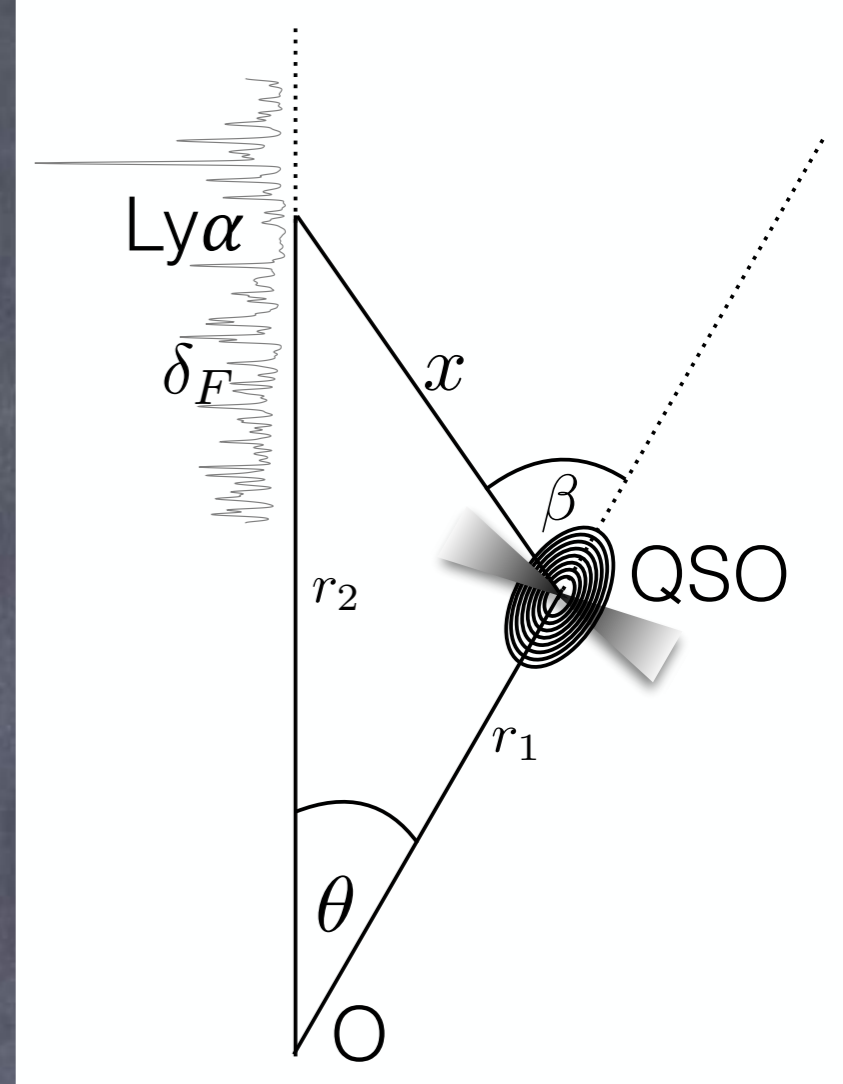
Bonvin & Durrer [arXiv:1105.5280],
 Challinor & Lewis [arXiv:1105.5292],
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Quasars x Ly-alpha correlation

Irsic , ED, Viel [arXiv:1510.03436]

$$\xi_{Q\alpha}(z_1, z_2, \theta) = \langle \Delta_Q(\mathbf{n}_1, z_1) \delta_F(\mathbf{n}_2, z_2) \rangle$$

Large bias factor difference

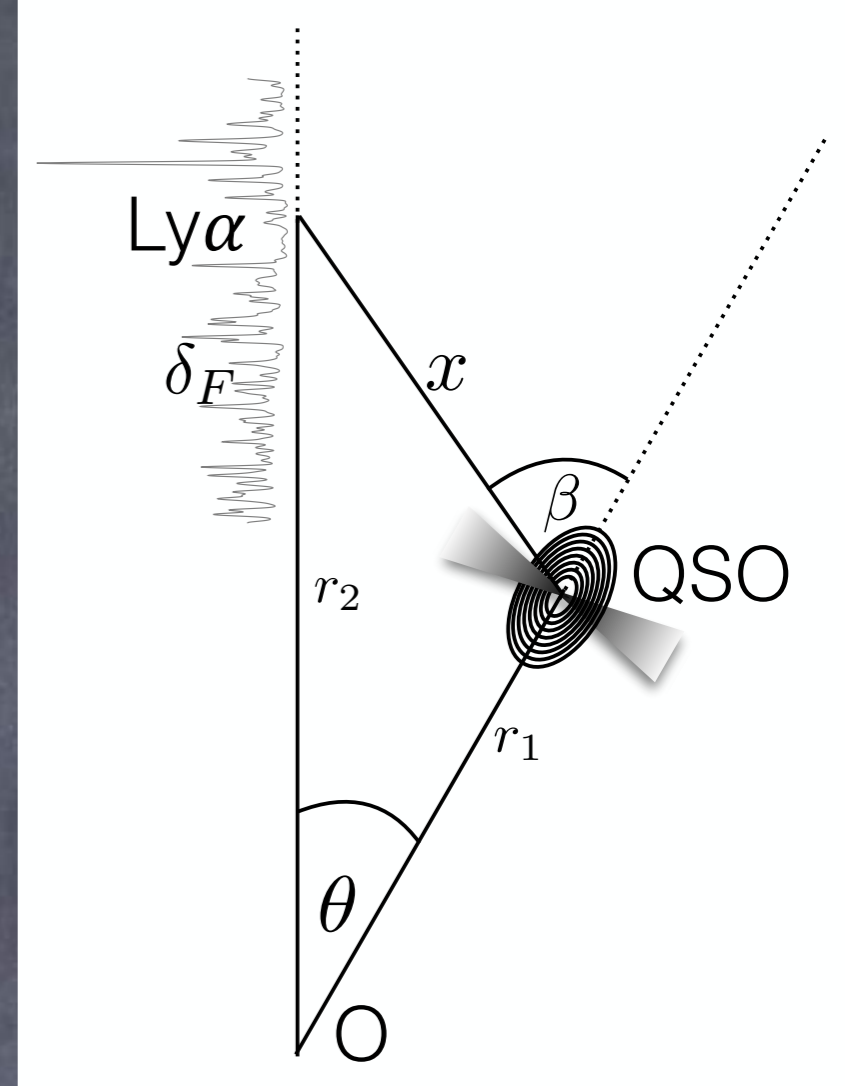


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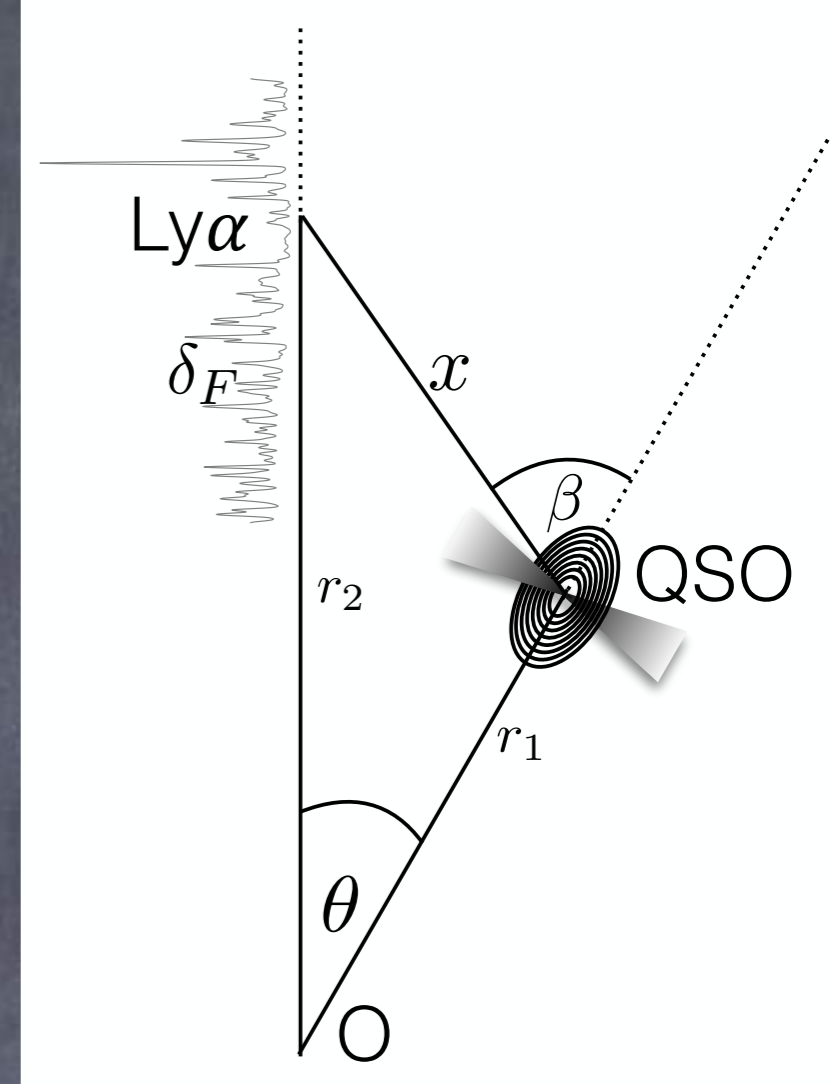
$$\delta_F(\mathbf{n}, z) = b_\alpha \delta^{\text{sync}} + b_v \mathcal{H}^{-1} \partial_r \mathbf{n} \cdot \mathbf{v} - \bar{\tau}(z) \left[- \left(2 + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} \right) \frac{\delta z}{1+z} + \mathbf{n} \cdot \mathbf{v} + \Psi + \mathcal{H}^{-1} \dot{\Phi} - 3\mathcal{H}v \right]$$

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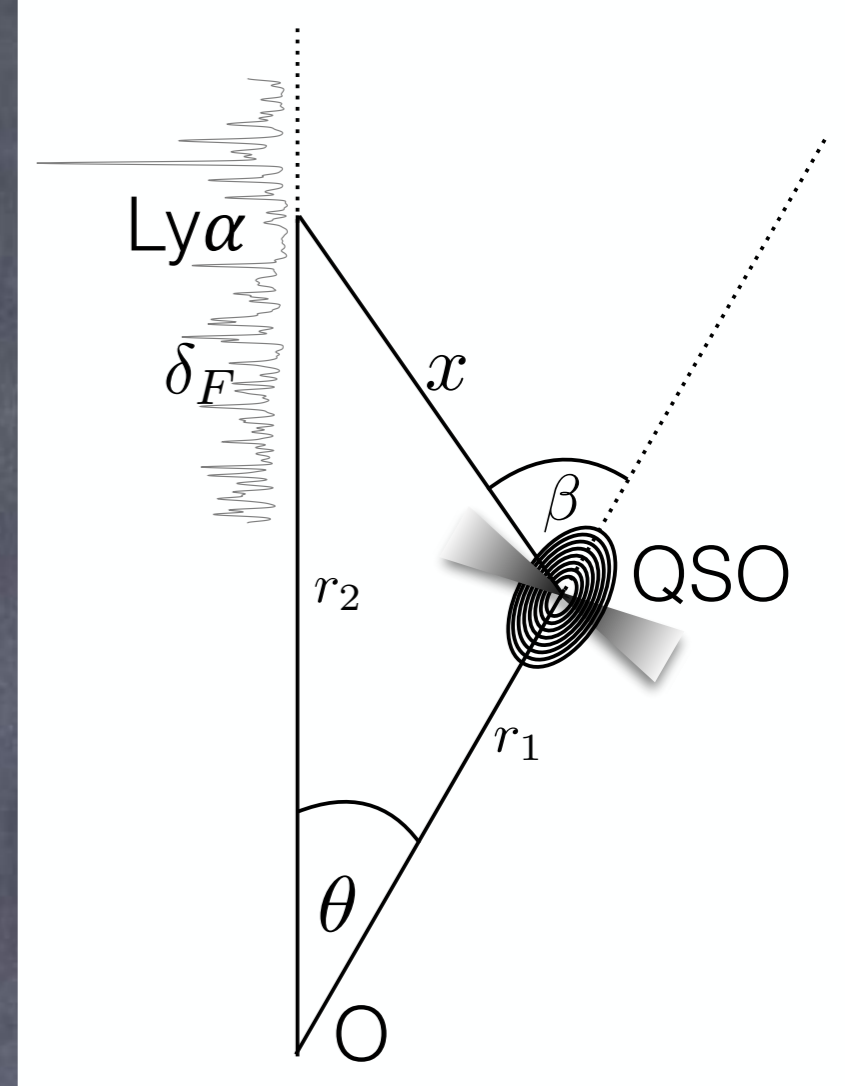
Standard

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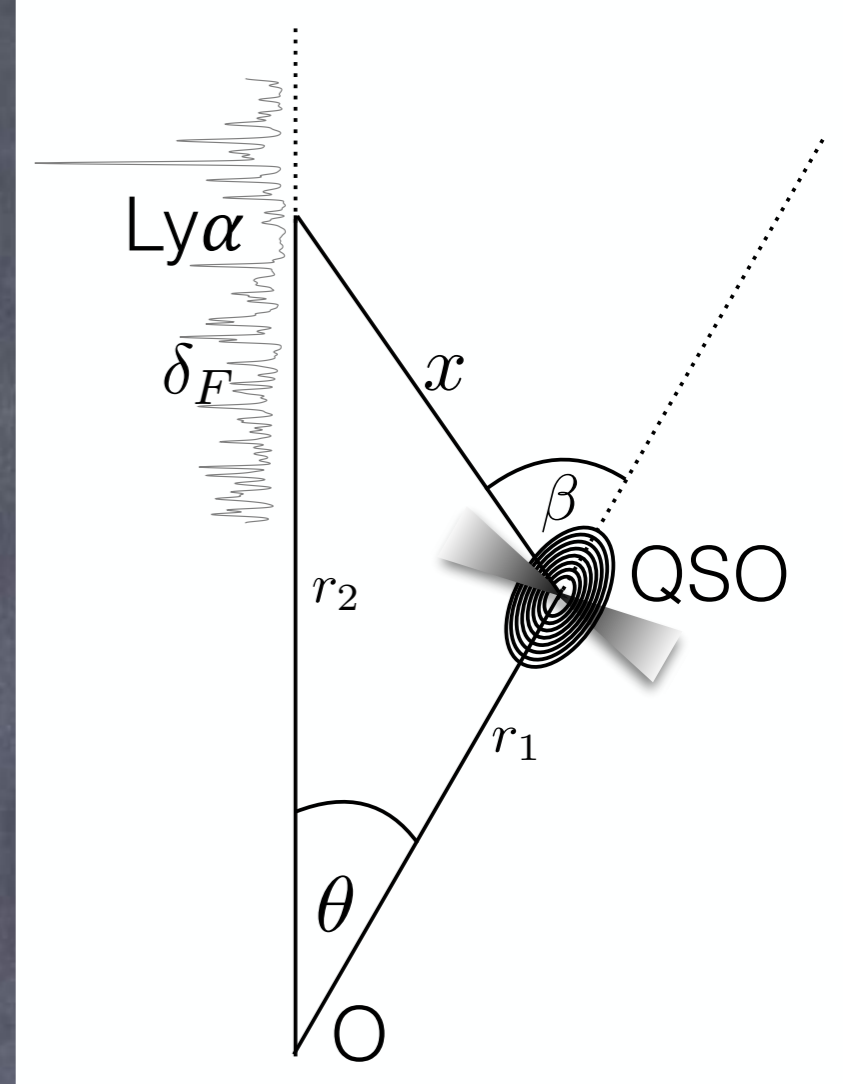
Relativistic

Quasars x Ly-alpha correlation

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$$\xi_{Q\alpha} = \xi_{Q\alpha}^{\text{newt}} + \xi_{Q\alpha}^{\text{magnification}} + \xi_{Q\alpha}^{\text{relativistic}}$$



Quasars x Ly-alpha correlation

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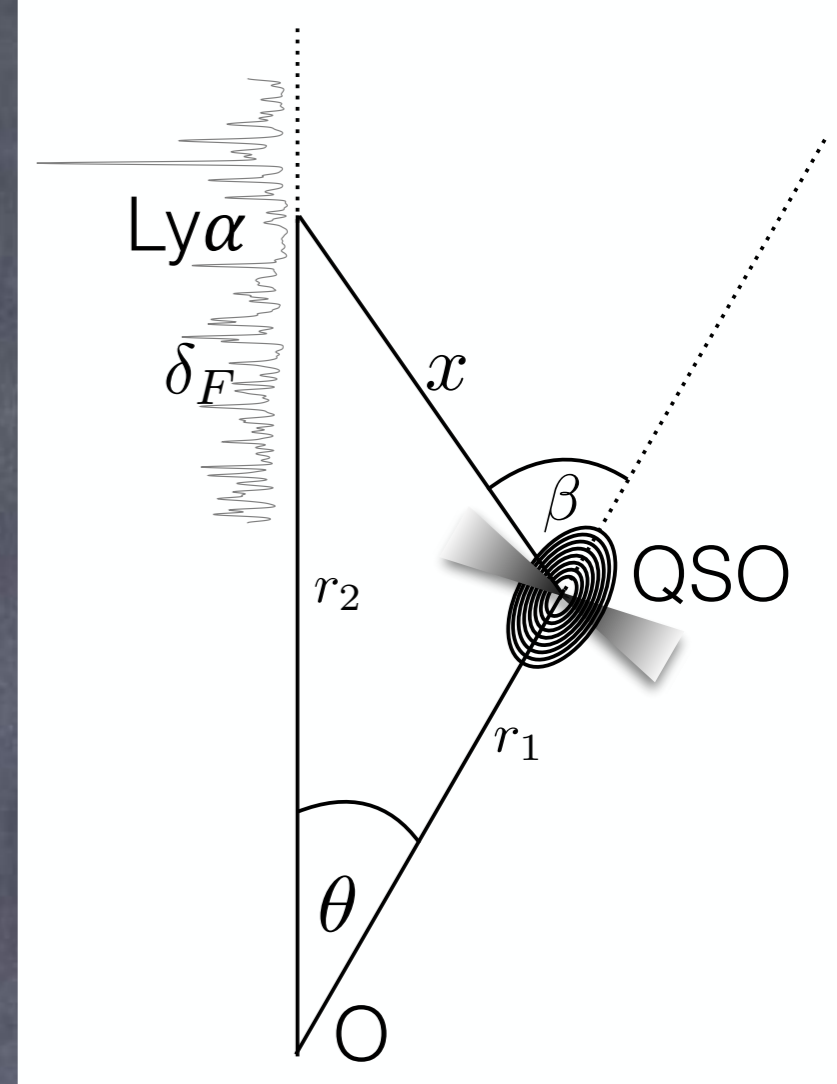
$$\xi_{Q\alpha}^{\text{newt}} \sim b_Q b_\alpha \int \frac{dk}{2\pi^2} k^2 P(k) j_0(kx)$$

$$+ b_v \int \frac{dk}{2\pi^2} k^2 f^2 P(k) \left[\frac{1}{5} j_0(kx) - \frac{4}{7} j_2(kx) + \frac{8}{35} j_4(kx) \right]$$

$$+ (b_Q b_v + b_\alpha) \int \frac{dk}{2\pi^2} k^2 f P(k) \left[\frac{1}{3} j_0(kx) - \frac{2}{3} j_2(kx) \right]$$

Order $\mathcal{O}(1)$

Even spherical Bessel functions



Quasars x Ly-alpha correlation

Irsic , ED, Viel [arXiv:1510.03436]

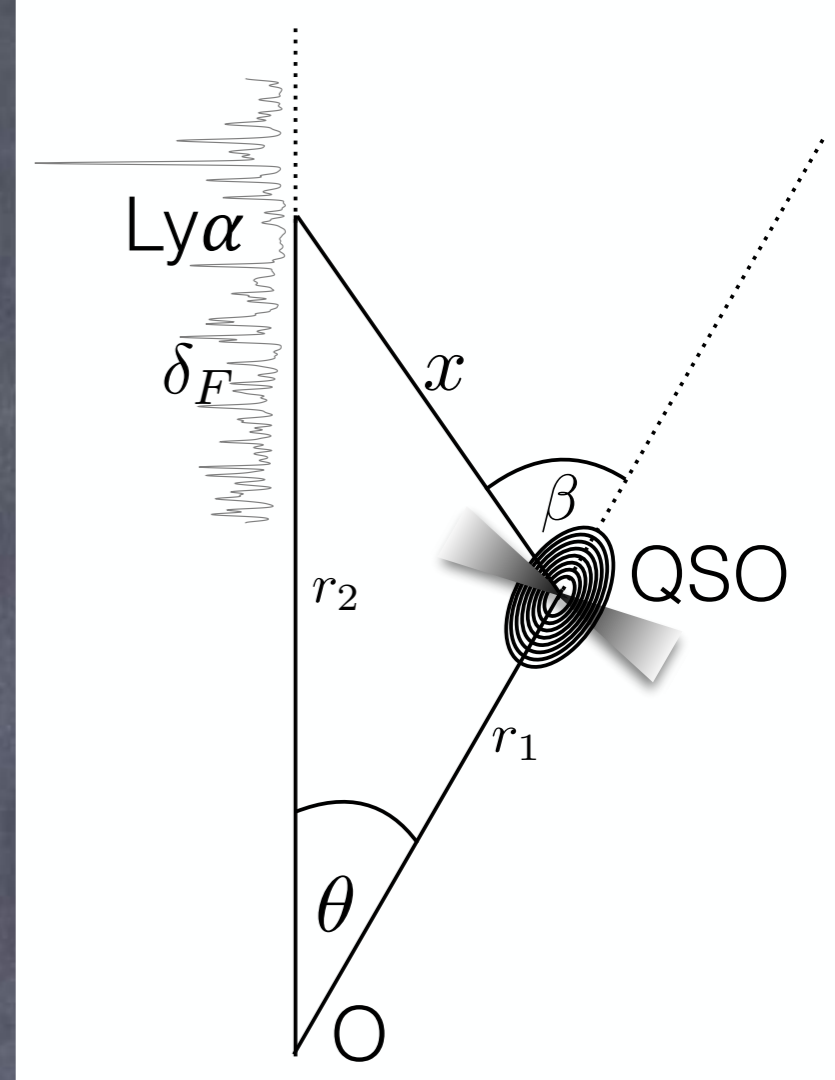
$$\xi_{Q\alpha}(z_1, z_2, \theta) = \langle \Delta_Q(\mathbf{n}_1, z_1) \delta_F(\mathbf{n}_2, z_2) \rangle$$

$$\xi_{Q\alpha} = \xi_{Q\alpha}^{\text{newt}} + \xi_{Q\alpha}^{\text{magnification}} + \xi_{Q\alpha}^{\text{relativistic}}$$

$$\begin{aligned} \xi_{Q\alpha}^{\text{Doppler}} \sim & (-b_Q \mathcal{R}_\alpha + \mathcal{R}_Q b_\alpha) \int \frac{dk}{2\pi^2} k^2 f P(k) j_1(kx) \frac{\mathcal{H}}{k} \\ & + (-\mathcal{R}_\alpha + \mathcal{R}_Q b_v) \int \frac{dk}{2\pi^2} k^2 f^2 P(k) \left[\frac{3}{5} j_1(kx) - \frac{2}{5} j_3(kx) \right] \frac{\mathcal{H}}{k} \\ & + \mathcal{R}_\alpha \mathcal{R}_Q \int \frac{dk}{2\pi^2} k^2 f^2 P(k) \left[\frac{1}{3} j_0(kx) - \frac{2}{3} j_2(kx) \right] \left(\frac{\mathcal{H}}{k} \right)^2 \end{aligned}$$

Order $\mathcal{O}(\mathcal{H}/k)$

Odd spherical Bessel functions



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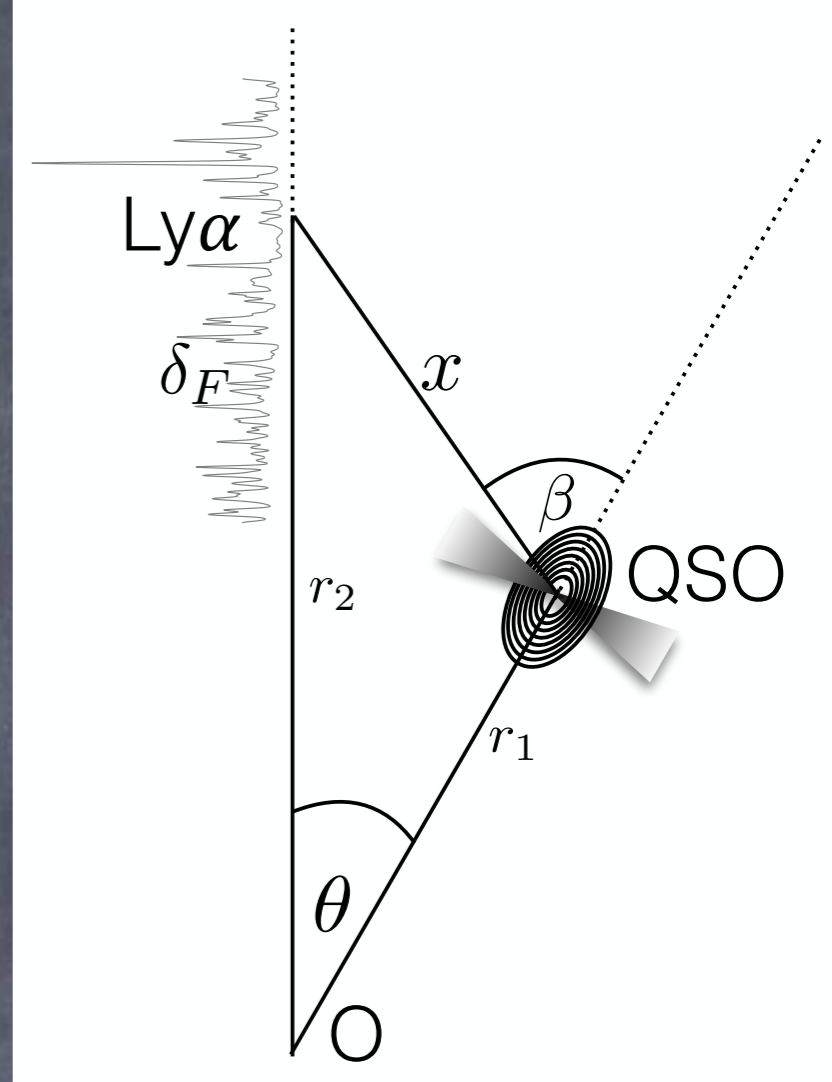
single tracer

$$\begin{aligned} \xi_{Q\alpha}^{\text{Doppler}} \sim & \int \frac{dk}{2\pi^2} k^2 f P(k) j_1(kx) \frac{\mathcal{H}}{k} \\ & + \int \frac{dk}{2\pi^2} k^2 f^2 P(k) \left[\frac{3}{5} j_1(kx) - \frac{2}{5} j_3(kx) \right] \frac{\mathcal{H}}{k} \\ & + \mathcal{R}_\alpha \mathcal{R}_Q \int \frac{dk}{2\pi^2} k^2 f^2 P(k) \left[\frac{1}{3} j_0(kx) - \frac{2}{3} j_2(kx) \right] \left(\frac{\mathcal{H}}{k} \right)^2 \end{aligned}$$

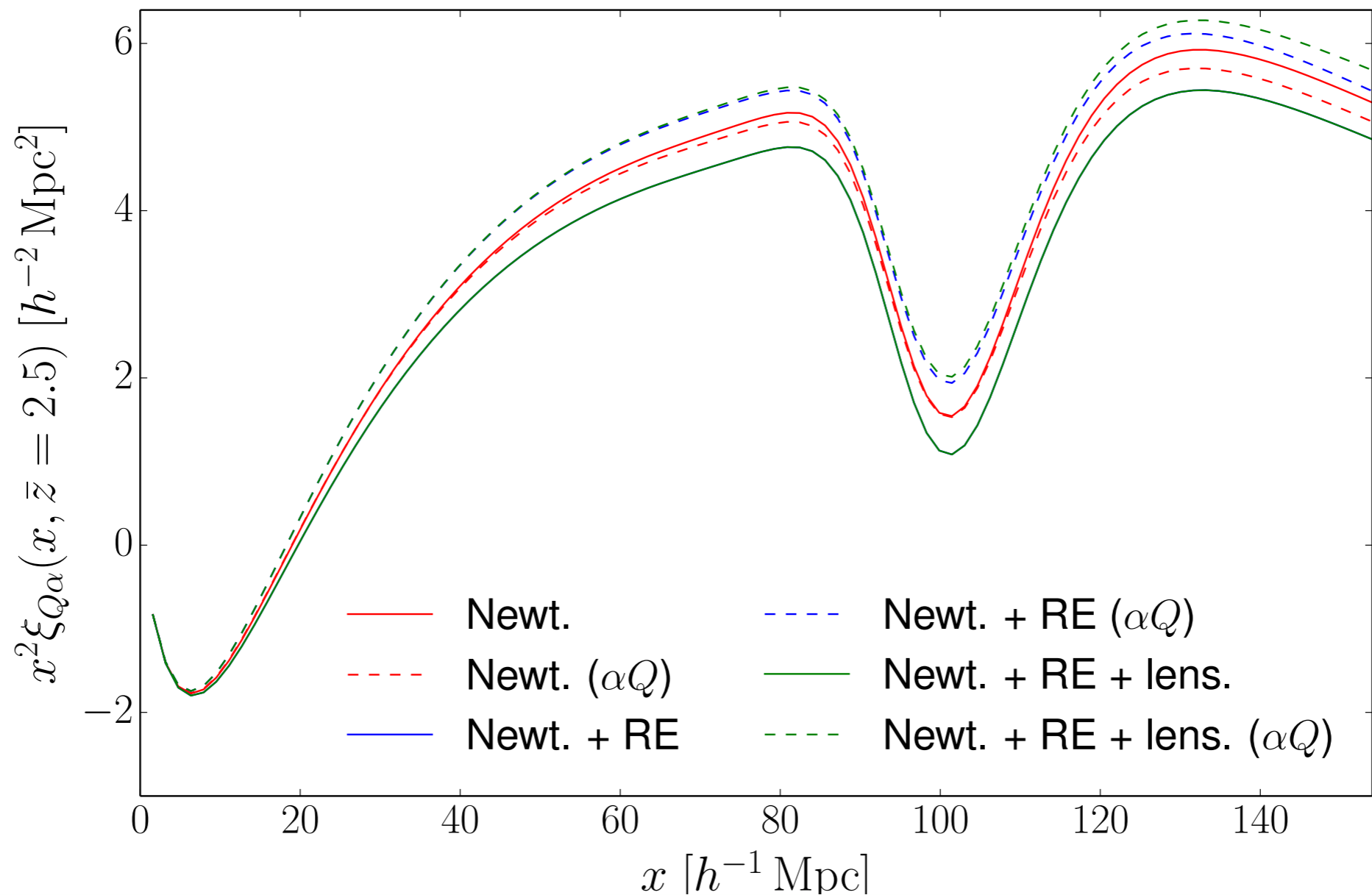
Order $\mathcal{O}(\mathcal{H}/k)$ $\mathcal{O}(\mathcal{H}^2/k^2)$

even

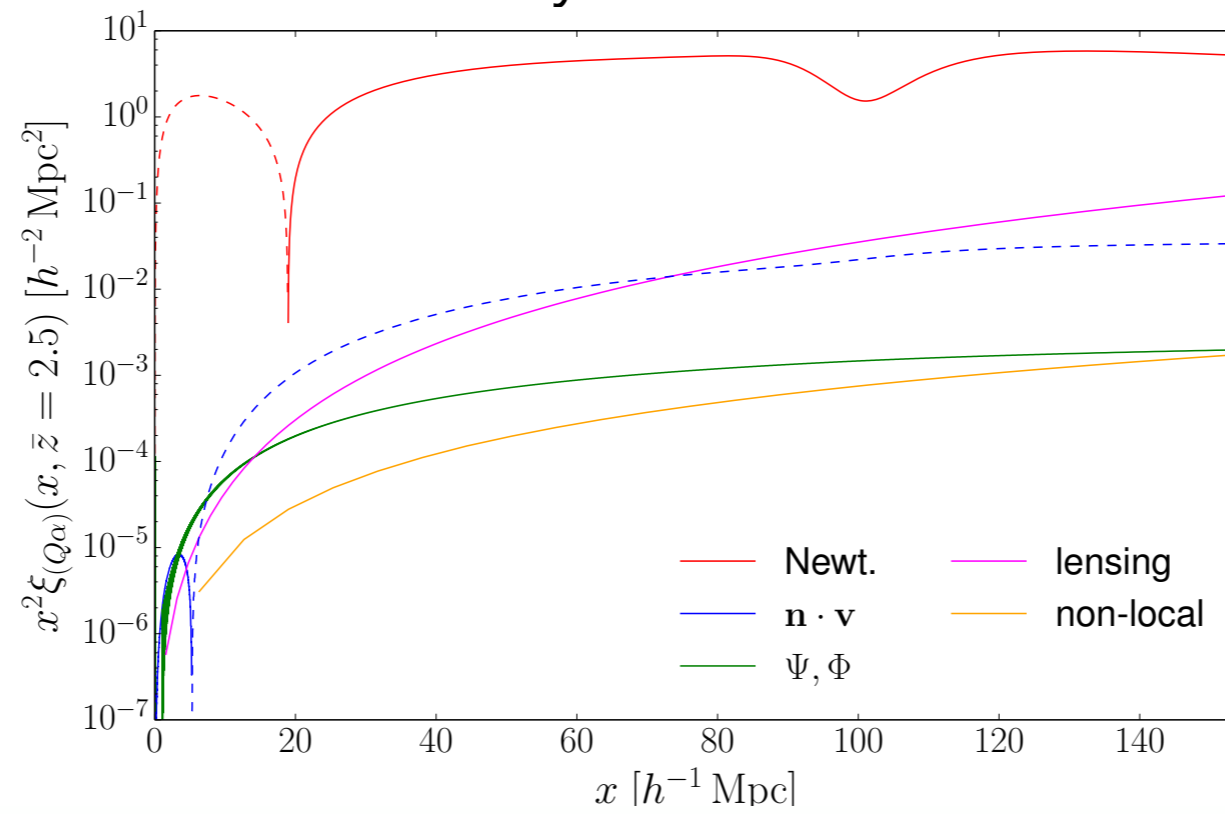
Odd spherical Bessel functions



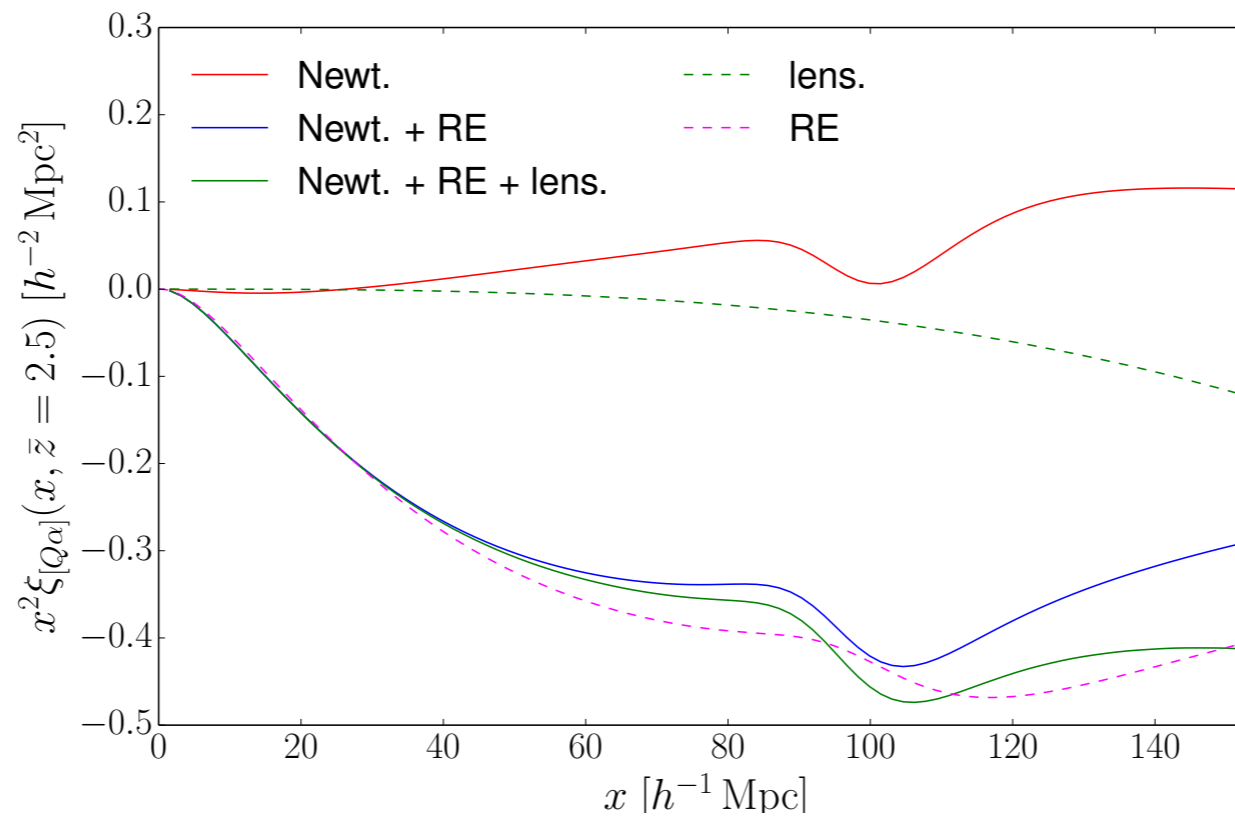
Relativistic effects on Lyman- α forest



Symmetric

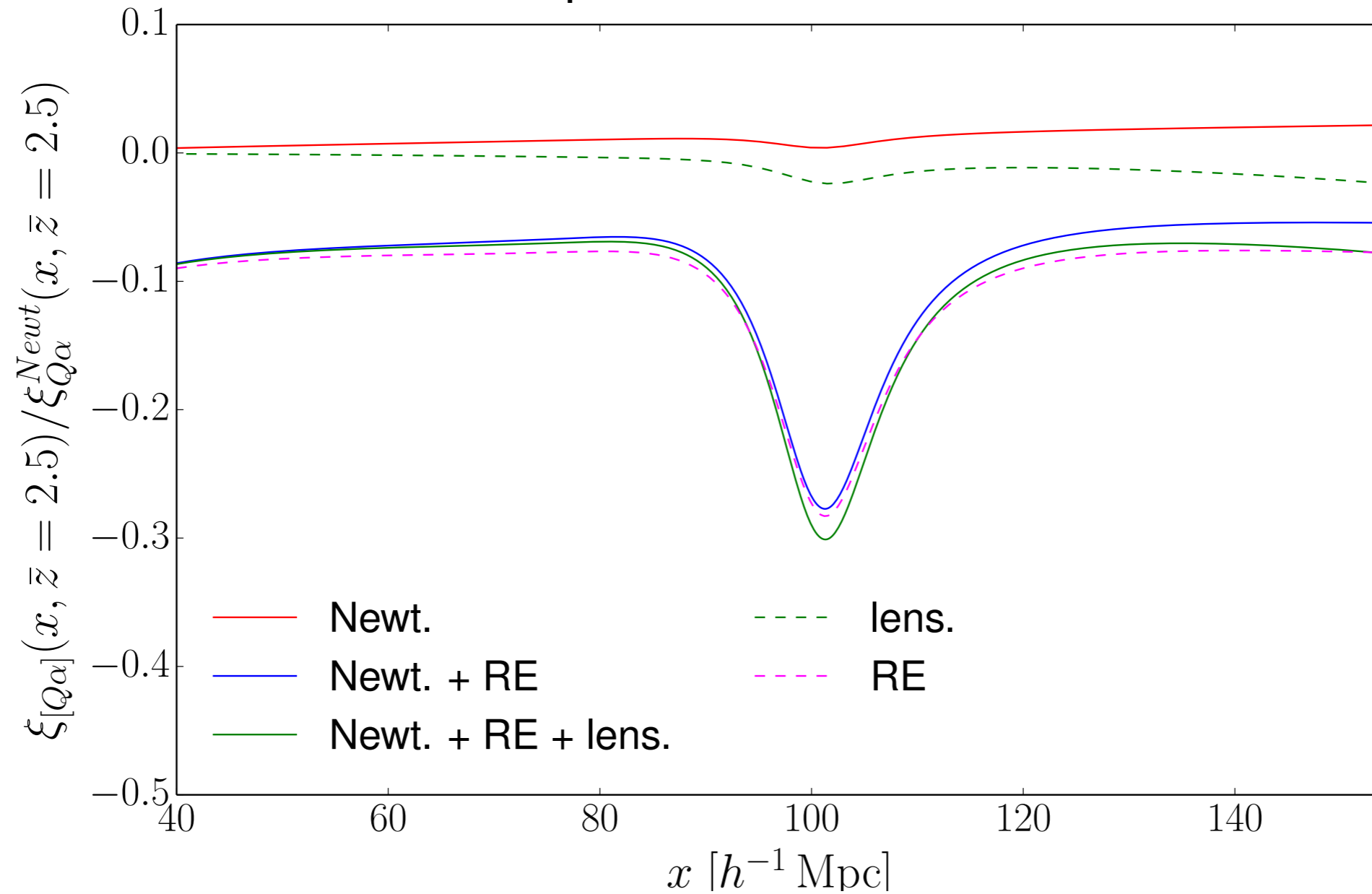


Anti-Symmetric

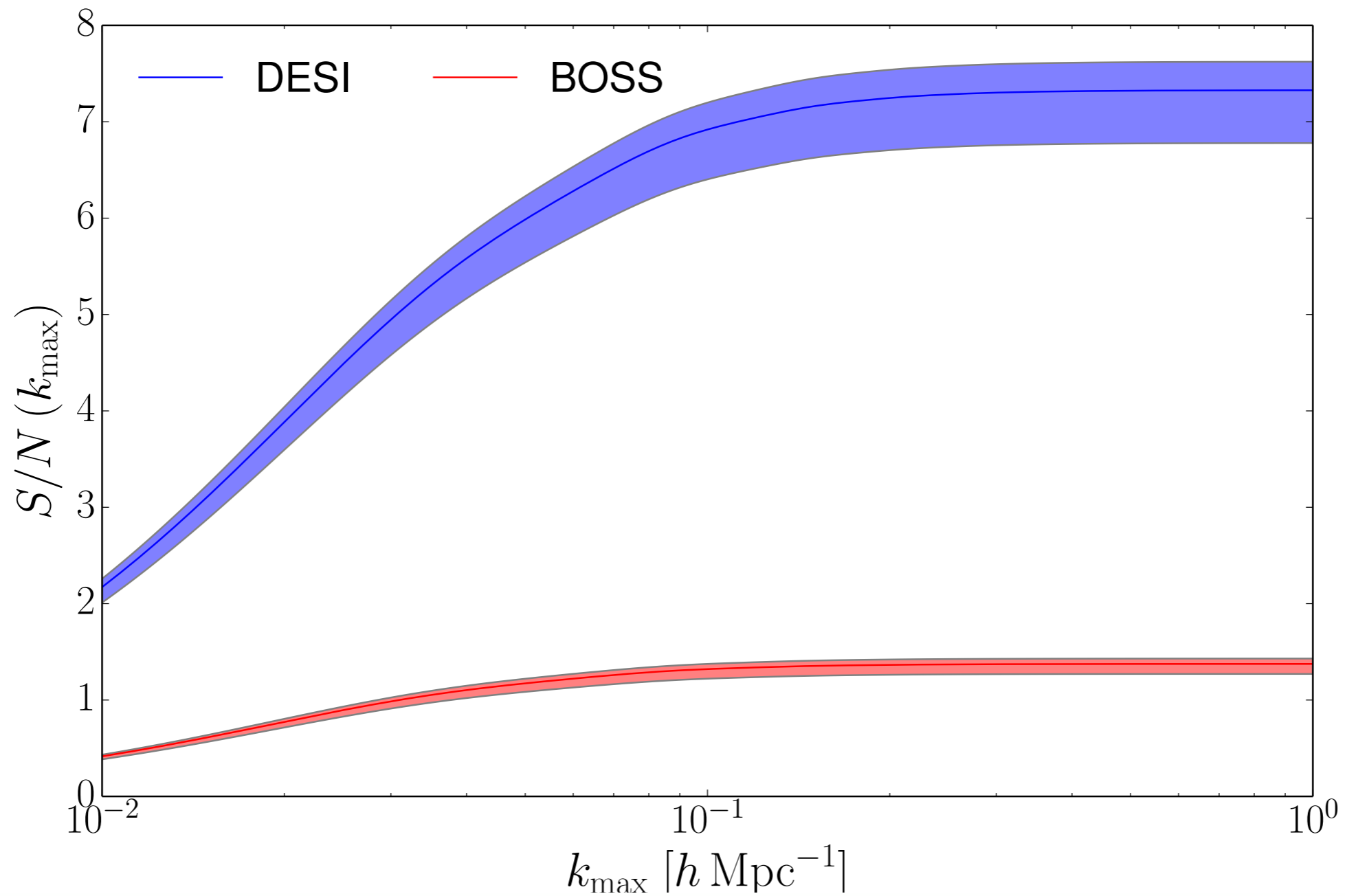


Relativistic effects on Lyman- α forest

Amplitude of effect



Relativistic effects on Lyman- α forest



Single tracer vs multi-tracers

	single	multi
Leading order correction	$\mathcal{O}(\mathcal{H}^2/k^2)$	$\mathcal{O}(\mathcal{H}/k)$
Parity	even	odd
Relevant scales	super-Hubble	all
Limited by	cosmic variance	shot noise
Forecasted detection	No	Yes

Large Scale Structures

$$\begin{aligned}
 \Delta_N(\mathbf{n}, z, m_*) &= b(z) D_{cm}(L > \bar{L}_*) + \mathcal{H}^{-1} \partial_r \mathbf{n} \cdot \mathbf{v} \\
 &\quad - \frac{2-5s}{2} \int_0^r \frac{r-r'}{rr'} \Delta_\Omega(\Psi + \Phi) dr' \\
 &\quad + \left(5s + \frac{2-5s}{\mathcal{H}r} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}}^N \right) \mathbf{n} \cdot \mathbf{v} + (f_{\text{evo}}^N - 3) \mathcal{H}v \\
 &\quad + \left(5s + \frac{2-5s}{\mathcal{H}r} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}}^N \right) \int_0^r (\dot{\Psi} + \dot{\Phi}) dr' \\
 &\quad + \left(5s + \frac{2-5s}{\mathcal{H}r} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}}^N \right) \Psi \\
 &\quad + \frac{2-5s}{r} \int_0^r (\Psi + \Phi) dr \\
 &\quad + (5s - 2) \Phi + \Psi + \mathcal{H}^{-1} \dot{\Phi}
 \end{aligned}$$

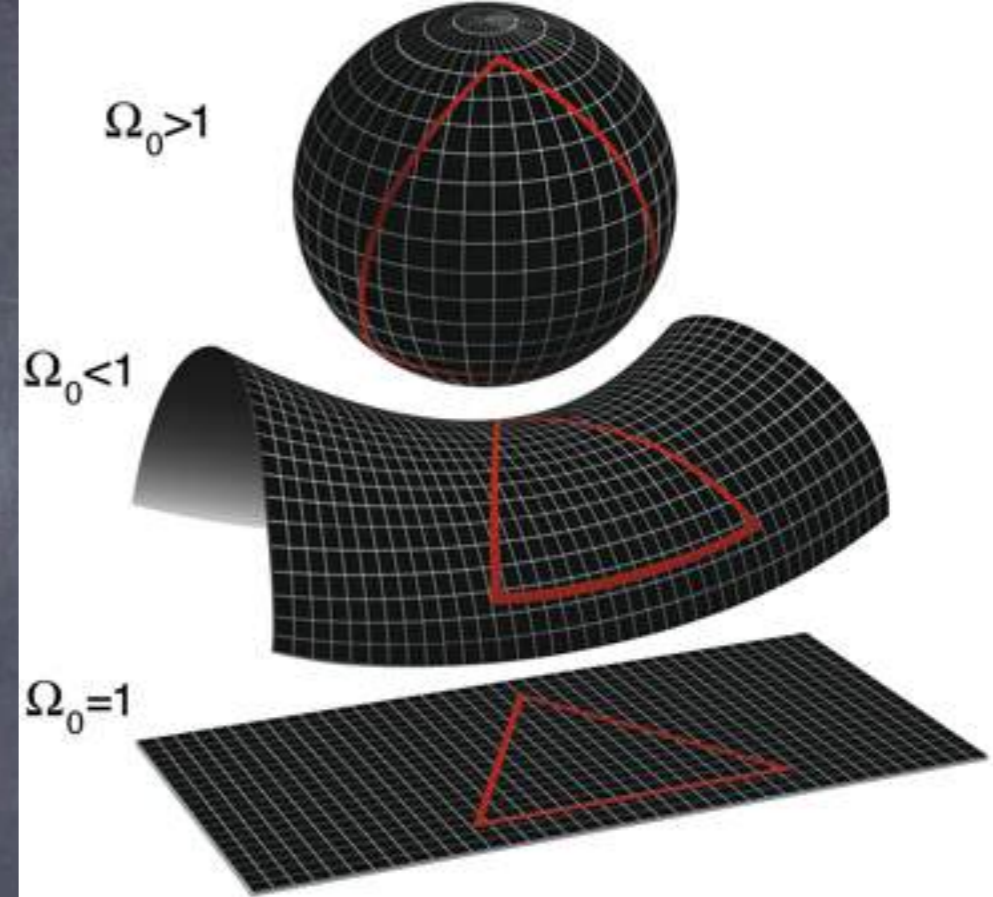
Bonvin & Durrer [arXiv:1105.5280],
 Challinor & Lewis [arXiv:1105.5292],
 Yoo [arXiv:1009.3021]

Curvature constraints from LSS

Tightest constraints $|\Omega_K| \leq 0.003$

Predicted limit of
detectability $\Omega_K \sim \mathcal{O}(10^{-4})$

Cosmic variance limit $\Omega_K \sim \mathcal{O}(10^{-5})$

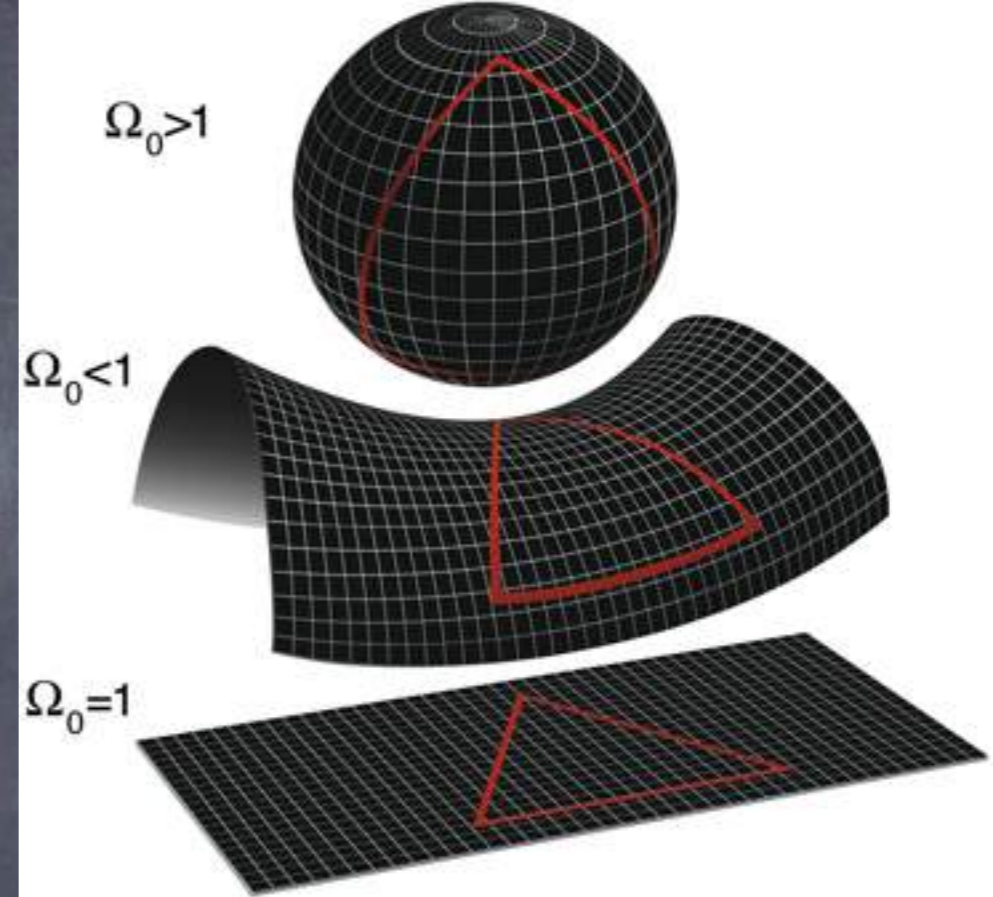


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$$d\tilde{s}^2 = a^2 \left(- (1 + 2\Psi) d\tau^2 + (1 - 2\Phi) \gamma_{ij} dx^i dx^j \right)$$

where $\gamma_{ij} dx^i dx^j = [dr^2 + \{S_K^2(r) (d\theta^2 + \sin^2 \theta d\varphi^2)\}]$

$$\text{and } S_K(r) = \begin{cases} \frac{1}{\sqrt{K}} \sin(\sqrt{K}r) & \text{for } K > 0 \\ r & \text{for } K = 0 \\ \frac{1}{\sqrt{|K|}} \sinh(\sqrt{|K|r}) & \text{for } K < 0 \end{cases}$$

Curvature constraints from LSS

$$\begin{aligned}
 \Delta(\mathbf{n}, z, m_*) &= bD_{cm} + \frac{1}{\mathcal{H}} \partial_r (\mathbf{n} \cdot \mathbf{v}) \\
 &\quad - \frac{2-5s}{2} \int_{\tau}^{\tau_0} \frac{S_K(r-\tilde{r})}{S_K(r)S_K(\tilde{r})} \Delta_{\Omega}(\Psi + \Phi) d\tilde{\tau} \\
 &\quad + \left(5s + \frac{2-5s}{\mathcal{H}} \frac{S'_K}{S_K} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} - f_{\text{evo}} \right) \\
 &\quad \times \left(\Psi + \mathbf{n} \cdot \mathbf{v} + \int_{\tau}^{\tau_0} (\dot{\Psi} + \dot{\Phi}) d\tilde{\tau} \right) \\
 &\quad + (5s-2)\Phi + \Psi + \frac{1}{\mathcal{H}} \dot{\Phi} + (f_{\text{evo}} - 3)\mathcal{H}v \\
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ED, Montanari, Raccanelli, Durrer,
Kamionkowski, Lesgourgues
[arXiv:1603.09073]

Curvature constraints from LSS

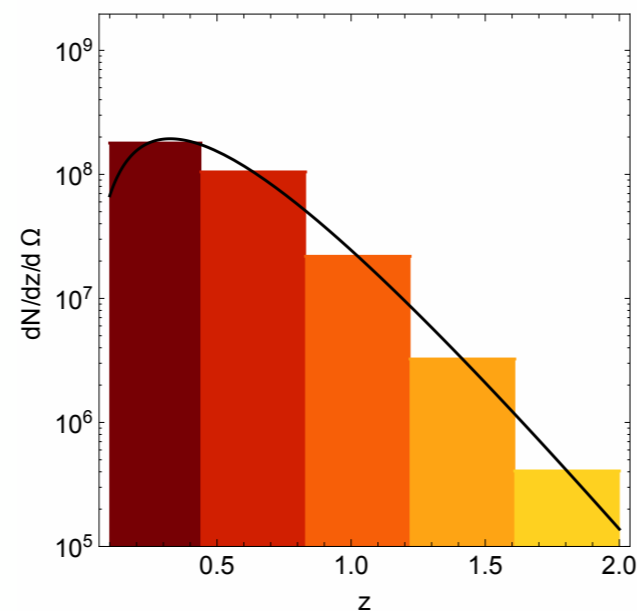
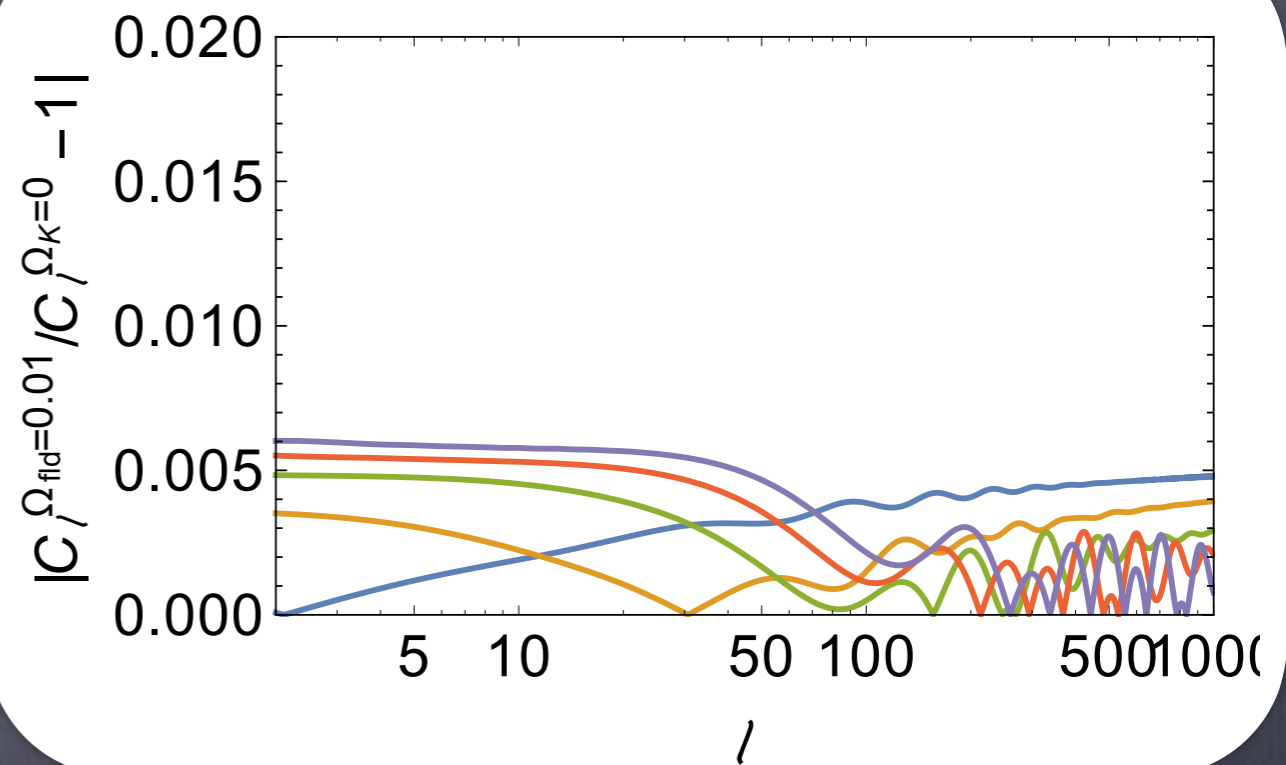
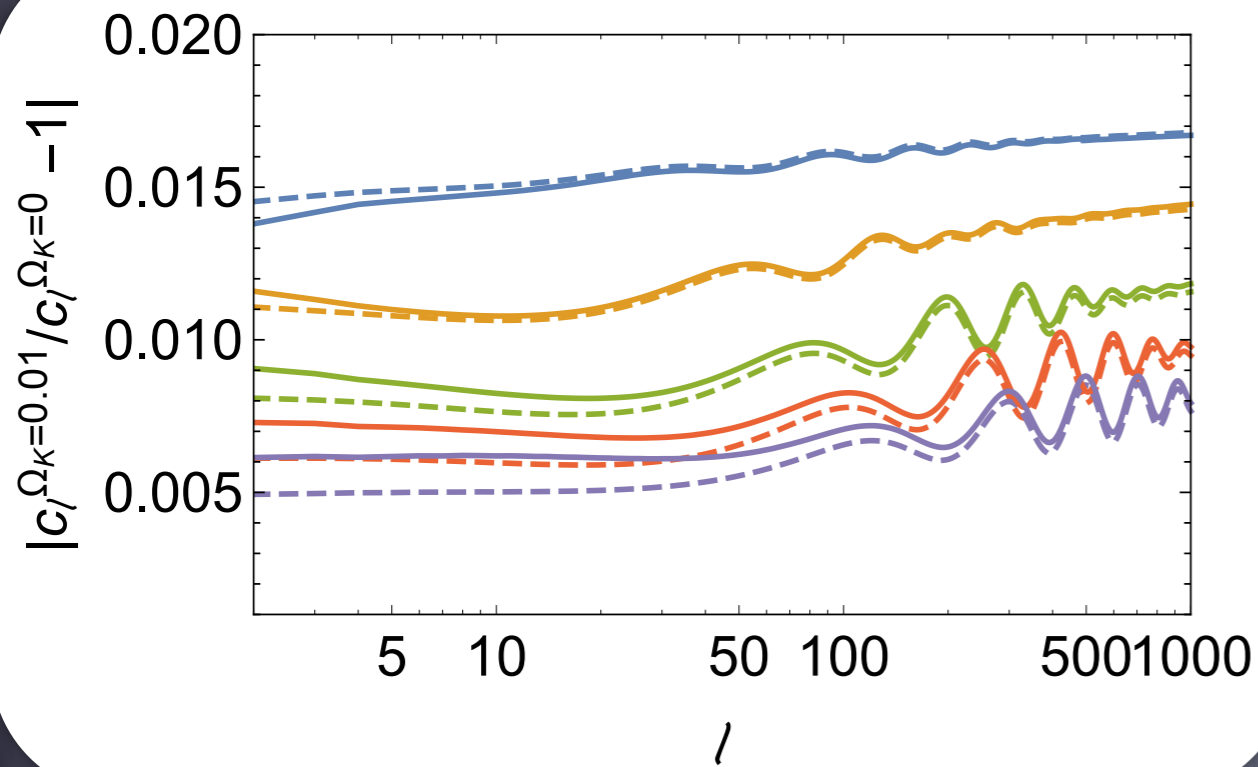
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 \end{aligned}$$

CLASS v.2.5

<http://class-code.net>

ED, Montanari, Raccanelli, Durrer,
Kamionkowski, Lesgourgues
[arXiv:1603.09073]

Curvature constraints from LSS



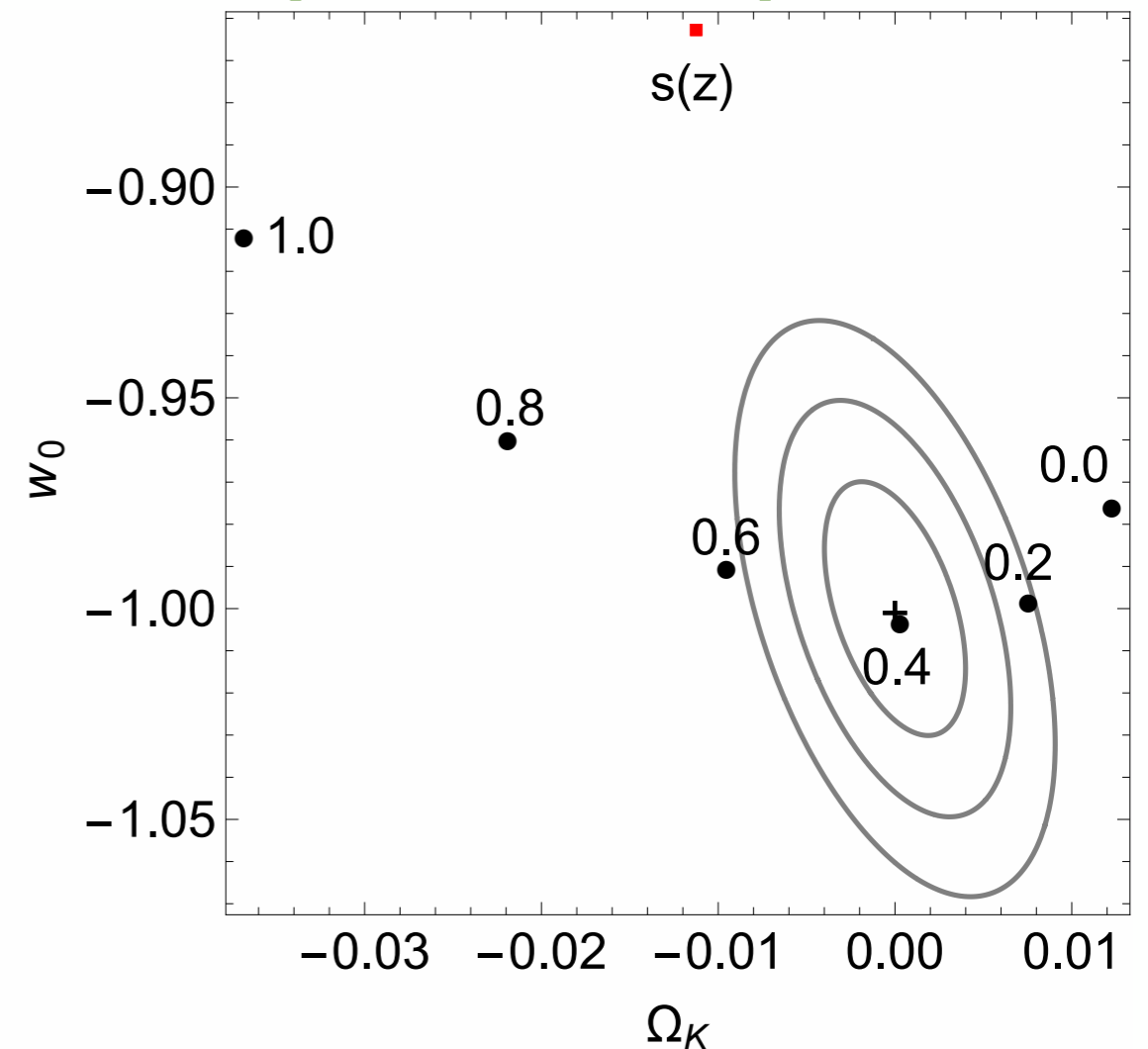
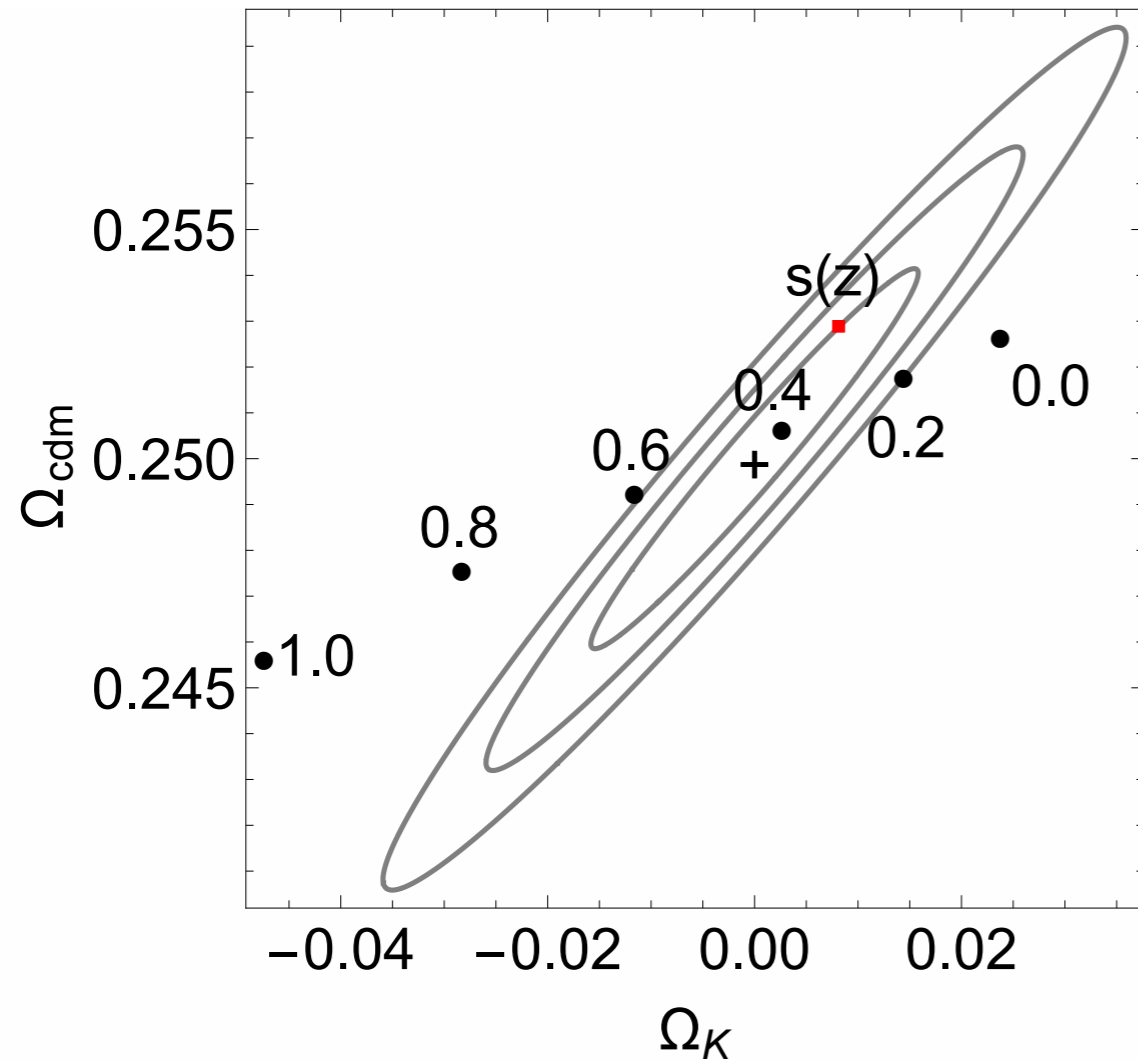
- 1-1
- 2-2
- 3-3
- 4-4
- 5-5

ED, Montanari, Raccanelli, Durrer,
Kamionkowski, Lesgourgues
[arXiv:1603.09073]

Curvature constraints from LSS

ED, Montanari, Raccanelli, Durrer,
Kamionkowski, Lesgourgues

[arXiv:1603.09073]



Spatial curvature, and other cosmological parameters, are strongly biased if relativistic effects (mainly cosmic magnification) are neglected.


Second Order

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In the weakly non-linear regime second order perturbation theory can be applied


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Born approximation fails
 Perturbed geodesics

Second Order

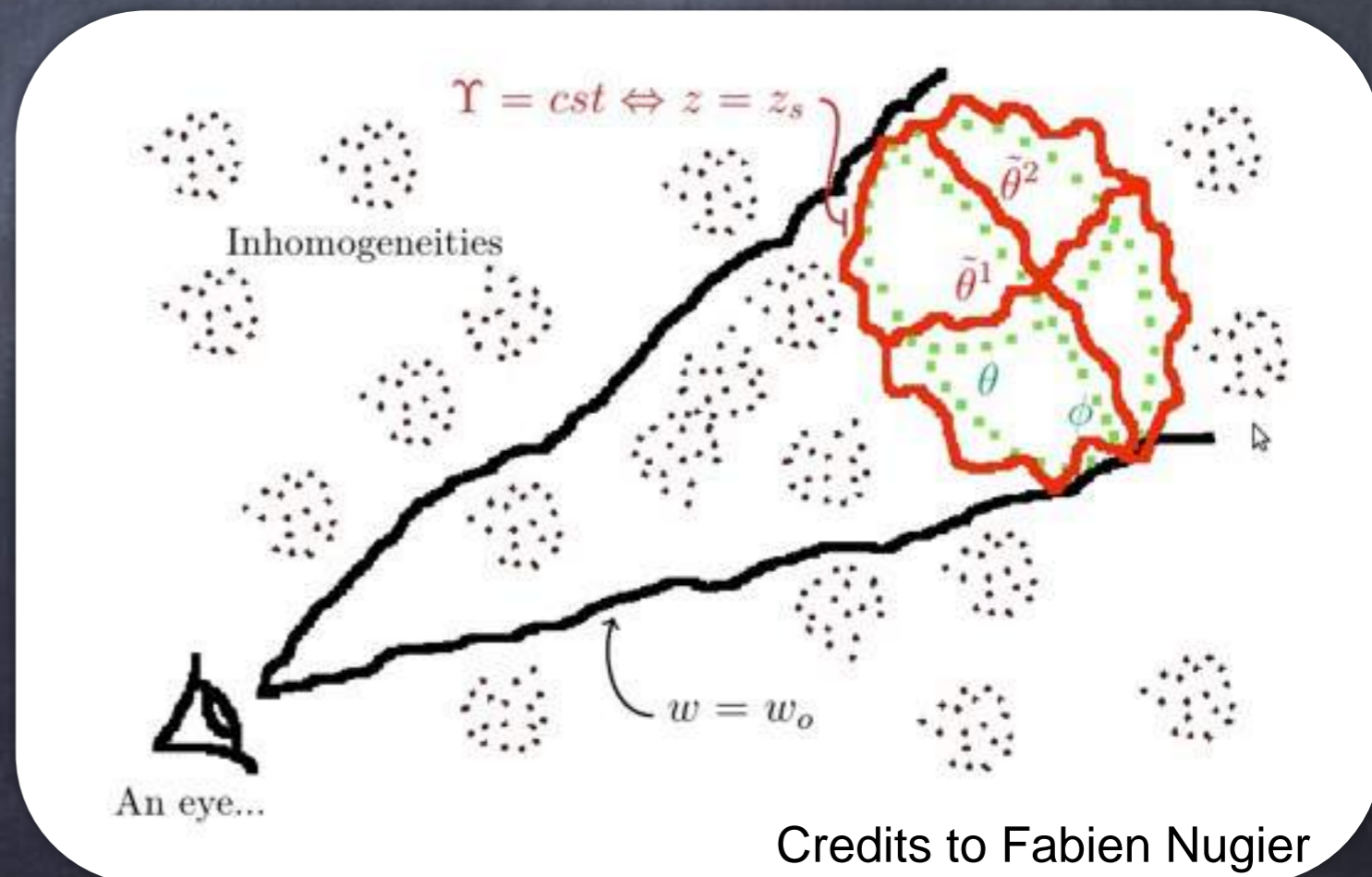
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Geodesic light-cone gauge

Gasperini, Marozzi, Nugier, Veneziano [arXiv:1104.1167]



Credits to Fabien Nugier

Second Order

$$\Delta(\mathbf{n}, z) \equiv \frac{N(\mathbf{n}, z) - \langle N \rangle (z)}{\langle N \rangle (z)}$$

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$$N(\mathbf{n}, z) = \rho(\mathbf{n}, z) V(\mathbf{n}, z)$$

Density perturbation

Volume perturbation

$$\rho(\mathbf{n}, z) = \bar{\rho} \left(1 + \delta^{(1)} + \delta^{(2)} \right) \quad V(\mathbf{n}, z) = \bar{V} \left(1 + \frac{\delta V^{(1)}}{\bar{V}} + \frac{\delta V^{(2)}}{\bar{V}} \right)$$

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$$\Delta(\mathbf{n}, z) = \left[\underbrace{\delta^{(1)} + \frac{\delta V^{(1)}}{\bar{V}}}_{1\text{-order}} + \delta^{(1)} \frac{\delta V^{(1)}}{\bar{V}} + \delta^{(2)} + \frac{\delta V^{(2)}}{\bar{V}} - \langle \delta^{(1)} \frac{\delta V^{(1)}}{\bar{V}} \rangle - \langle \delta^{(2)} \rangle - \left\langle \frac{\delta V^{(2)}}{\bar{V}} \right\rangle \right]$$

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2-order

$$\begin{aligned}
\Sigma_{IS} = & \left(-\frac{2}{\mathcal{H}_s r_s} - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \left\{ -v_{\parallel s}^{(2)} - \frac{1}{2} \phi_s^{(2)} - \frac{1}{2} \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \left[\phi^{(2)}(\eta') + \psi^{(2)}(\eta') \right] + \frac{1}{2} (v_{\parallel s})^2 \right. \\
& + \frac{1}{2} (\psi_s^I)^2 + (-v_{\parallel s} - \psi_s^I) \left(-\psi_s^I - 2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right) + \frac{1}{2} v_{\perp s}^a v_{\perp s}^a \\
& + 2a v_{\perp s}^a \partial_a \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') + 4 \int_{\eta_s}^{\eta_o} d\eta' \left[\psi^I(\eta') \partial_{\eta'} \psi^I(\eta') + \partial_{\eta'} \psi^I(\eta') \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right. \\
& \left. + \psi^I(\eta') \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''}^2 \psi^I(\eta'') - \gamma_0^{ab} \partial_a \left(\int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \partial_b \left(\int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right) \right] \\
& + 2\partial_a (v_{\parallel s} + \psi_s^I) \int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{ab} \partial_b \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \\
& + 4 \int_{\eta_s}^{\eta_o} d\eta' \partial_a \left(\partial_{\eta'} \psi^I(\eta') \right) \int_{\eta'}^{\eta_o} d\eta'' \gamma_0^{ab} \partial_b \int_{\eta''}^{\eta_o} d\eta''' \psi^I(\eta''') \left. \right\} \\
& + \left[\frac{1}{2} \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} + \frac{3}{2} \left(\frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right)^2 - \frac{1}{2} \frac{\mathcal{H}''_s}{\mathcal{H}_s^3} + \frac{1}{\mathcal{H}_s r_s} \left(1 + 3 \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} + \frac{1}{\mathcal{H}_s r_s} \right) \right] \left[(v_{\parallel s})^2 + (\psi_s^I)^2 + 2\psi_s^I v_{\parallel s} \right. \\
& + 4(v_{\parallel s} + \psi_s^I) \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') + 4 \left(\int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right)^2 \left. \right] - \psi_s^{(2)} + \frac{1}{2} \phi_s^{(2)} + \frac{1}{2\mathcal{H}_s} \partial_{\eta} \psi_s^{(2)} \\
& + \frac{1}{\mathcal{H}_s} \partial_r v_{\parallel s}^{(2)} - \frac{1}{2} \frac{1}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \left[\psi^{(2)} + \phi^{(2)} \right](\eta') + \frac{1}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \left[\psi^{(2)} + \phi^{(2)} \right](\eta') \\
& + 2 \left(1 - \frac{1}{\mathcal{H}_s r_s} \right) \left\{ -\frac{2}{\mathcal{H}_s} v_{\parallel s} \partial_r v_{\parallel s} - (v_{\parallel s})^2 - v_{\perp s}^a v_{\perp s}^a + \left[-\frac{1}{\mathcal{H}_s} \partial_{\eta} \psi_s^I - 2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right. \right. \\
& - \frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') + \frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') \left. \right] v_{\parallel s} + \left[-2\psi_s^I - 4 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right. \\
& \left. - 2\mathcal{H}_s \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right] \frac{1}{\mathcal{H}_s} \partial_r v_{\parallel s} - a v_{\perp s}^a \partial_a \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') + [\partial_r \psi_s^I + 2\partial_{\eta} \psi_s^I \\
& + 2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'}^2 \psi^I(\eta') \left. \right] \left(-2 \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right) \\
& - \left(-\psi_s^I - 2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right) \left[\frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') - \frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right. \\
& \left. - \frac{1}{\mathcal{H}_s} \partial_{\eta} \psi_s^I - \psi_s^I \right] \left. \right\} + \frac{3}{2} v_{\perp s}^a v_{\perp s}^a - \frac{2}{\mathcal{H}_s} a v_{\perp s}^a \partial_a v_{\parallel s} + \left(\frac{5}{2} + \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) (v_{\parallel s})^2 \\
& + \left(5 + 3 \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \frac{1}{\mathcal{H}_s} v_{\parallel s} \partial_r v_{\parallel s} + \frac{1}{\mathcal{H}_s^2} \left[v_{\parallel s} \partial_r^2 v_{\parallel s} + (\partial_r v_{\parallel s})^2 \right] + \left[-\frac{1}{\mathcal{H}_s^2} (\partial_r^2 \psi_s^I + \partial_{\eta}^2 \psi_s^I - \partial_{\eta} \partial_r \psi_s^I) \right. \\
& - \frac{1}{\mathcal{H}_s} \partial_r \psi_s^I - \frac{3}{\mathcal{H}_s} \left(-1 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} - \frac{2}{3} \frac{1}{\mathcal{H}_s r_s} \right) \partial_{\eta} \psi_s^I - \frac{4}{r_s} \left(-1 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \\
& + \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') - 2 \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \\
& + \frac{4}{\mathcal{H}_s r_s} \psi_s^I - \frac{2}{\mathcal{H}_s r_s^2} \int_{\eta_s}^{\eta_o} d\eta' \Delta_2 \psi^I(\eta') \left. \right] v_{\parallel s} + \left[2 \left(2 + \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} + \frac{2}{\mathcal{H}_s r_s} \right) \partial_r v_{\parallel s} \right. \\
& \left. + \frac{2}{\mathcal{H}_s} \partial_r^2 v_{\parallel s} \right] \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') + \left[\frac{2}{\mathcal{H}_s} \left(5 + 3 \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \partial_r v_{\parallel s} + \frac{2}{\mathcal{H}_s^2} \partial_r^2 v_{\parallel s} \right] \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \\
& - \frac{2}{\mathcal{H}_s} \partial_r v_{\parallel s} \frac{1}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') + \frac{2}{\mathcal{H}_s^2} \partial_{\eta} \psi_s^I \partial_r v_{\parallel s} + \frac{1}{\mathcal{H}_s} \left[\frac{1}{\mathcal{H}_s} \partial_r^2 v_{\parallel s} \right. \\
& \left. + \left(6 + 3 \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \partial_r v_{\parallel s} \right] \psi_s^I + \frac{2}{\mathcal{H}_s r_s} a v_{\perp s}^a \partial_a \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') - \frac{1}{\mathcal{H}_s} a v_{\perp s}^a \partial_a \psi_s^I \\
& - \frac{4}{\mathcal{H}_s} \gamma_0^{ab} \partial_a v_{\parallel s} \partial_b \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') + \frac{4}{r_s^2} \left(\int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right)^2 + \left\{ \left[2 \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \psi_s^I \right. \right. \\
& + 4 \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') - \frac{2}{\mathcal{H}_s} \partial_{\eta} \psi_s^I \left. \right] \frac{1}{r_s} + 2 \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'}^2 \psi^I(\eta') \\
& \left. + 2 \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \partial_{\eta} \psi_s^I + \left(-1 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \partial_r \psi_s^I - \frac{1}{\mathcal{H}_s} \partial_{\eta} \partial_r \psi_s^I \right\} \left(-2 \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right)
\end{aligned}$$

$$\begin{aligned}
& + \left[\frac{3}{\mathcal{H}_s} \left(-1 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} - \frac{2}{3} \frac{1}{\mathcal{H}_s r_s} \right) \partial_{\eta} \psi_s^I + \frac{1}{\mathcal{H}_s} \partial_r \psi_s^I - \left(2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \psi_s^I \right. \\
& \left. + \frac{1}{\mathcal{H}_s^2} (\partial_r^2 \psi_s^I + \partial_{\eta}^2 \psi_s^I - \partial_{\eta} \partial_r \psi_s^I) \right] \left(-2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right) + \left[\left(-2 - 2 \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} - \frac{2}{\mathcal{H}_s r_s} \right) \psi_s^I \right. \\
& \left. + 4 \left(-2 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} - \frac{1}{\mathcal{H}_s r_s} \right) \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') - \frac{10}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right. \\
& \left. - \frac{2}{\mathcal{H}_s} \partial_{\eta} \psi_s^I \right] \frac{1}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') + 2 \left(\frac{1}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') \right)^2 \\
& + \left(-\frac{1}{2} - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) (\psi_s^I)^2 + \frac{1}{\mathcal{H}_s^2} (\partial_{\eta} \psi_s^I)^2 + \left[-\frac{1}{\mathcal{H}_s^2} (\partial_r^2 \psi_s^I + \partial_{\eta}^2 \psi_s^I - \partial_{\eta} \partial_r \psi_s^I) \right. \\
& \left. + \frac{1}{\mathcal{H}_s} \left(4 + 3 \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} + \frac{2}{\mathcal{H}_s r_s} \right) \partial_{\eta} \psi_s^I - \frac{1}{\mathcal{H}_s} \partial_r \psi_s^I \right] \psi_s^I \\
& + \frac{4}{\mathcal{H}_s} \gamma_0^{ab} \partial_a \left(\int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right) \partial_b \left(\int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right) + 2\partial_a \left[\left(1 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) v_{\parallel s} - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \psi_s^I \right. \\
& \left. - \frac{1}{\mathcal{H}_s} (\partial_{\eta} \psi_s^I + \partial_r v_{\parallel s}) \right] \int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{ab} \partial_b \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \\
& + 4 \left(1 - \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) \int_{\eta_s}^{\eta_o} d\eta' \partial_a \left(\partial_{\eta'} \psi^I(\eta') \right) \int_{\eta'}^{\eta_o} d\eta'' \gamma_0^{ab} \partial_b \int_{\eta''}^{\eta_o} d\eta''' \psi^I(\eta''') \\
& + 2a v_{\perp s}^a \partial_a \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') + \frac{4}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \left[\psi^I(\eta') \left(-\psi^I(\eta') - 2 \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right) \right. \\
& \left. + \gamma_0^{ab} \partial_a \left(\int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \partial_b \left(\int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \right] + 4\partial_a \psi_s^I \int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{ab} \partial_b \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \\
& + \frac{4}{r_s} \left[\int_{\eta_s}^{\eta_o} d\eta' \partial_a \left(\frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') - 2\psi^I(\eta') \right) \int_{\eta'}^{\eta_o} d\eta'' \gamma_0^{ab} \partial_b \int_{\eta''}^{\eta_o} d\eta''' \psi^I(\eta''') \right] \\
& + 2\partial_a \left(\int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right) \left[-\gamma_0^{ab} \partial_b \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') - 2 \int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{ab} \partial_b \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right] \\
& + 2\partial_a \left(\int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{bc} \partial_c \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \partial_b \left(\int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{ad} \partial_d \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \\
& - 4 \left(\int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') \right) \int_{\eta_s}^{\eta_o} d\eta' \left[-\frac{1}{(\eta_o - \eta')^3} \int_{\eta'}^{\eta_o} d\eta'' \Delta_2 \psi^I(\eta'') + \frac{1}{(\eta_o - \eta')^2} \left(\frac{1}{2} \Delta_2 \psi^I(\eta') \right. \right. \\
& \left. \left. + \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} (\Delta_2 \psi^I(\eta'')) \right) \right] + \frac{2}{(\sin \theta_o)^2} \left[\frac{1}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \partial_{\theta_o} \psi^I(\eta') \right]^2 \\
& + \frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \left[\psi^I(\eta') \left(-\psi^I(\eta') - 2 \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right) \right. \\
& \left. + \gamma_0^{ab} \partial_a \left(\int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \partial_b \left(\int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \right] \\
& + 2 \int_{\eta_s}^{\eta_o} d\eta' \left\{ -2\psi^I(\eta') \frac{1}{\eta_o - \eta'} \int_{\eta'}^{\eta_o} d\eta'' \frac{\eta'' - \eta'}{\eta_o - \eta''} \Delta_2 \partial_{\eta''} \psi^I(\eta'') \right. \\
& \left. + 2\gamma_0^{ab} \partial_b \left(\int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right) \frac{1}{\eta_o - \eta'} \int_{\eta'}^{\eta_o} d\eta'' \frac{\eta'' - \eta'}{\eta_o - \eta''} \partial_a \Delta_2 \psi^I(\eta'') \right. \\
& \left. - \left(-2\psi^I(\eta') - 2 \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right) \frac{1}{(\eta_o - \eta')^2} \int_{\eta'}^{\eta_o} d\eta'' \Delta_2 \psi^I(\eta'') \right. \\
& \left. - 2\partial_a \psi^I(\eta') \int_{\eta'}^{\eta_o} d\eta'' \gamma_0^{ab} \partial_b \int_{\eta''}^{\eta_o} d\eta''' \partial_{\eta'''} \psi^I(\eta''') \right. \\
& \left. + 2\partial_a \left[\gamma_0^{ab} \partial_b \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right] \int_{\eta'}^{\eta_o} d\eta'' \partial_d \left[\gamma_0^{ac} \partial_c \int_{\eta''}^{\eta_o} d\eta''' \psi^I(\eta''') \right] \right. \\
& \left. + 2\gamma_0^{ab} \partial_a \left(\psi^I(\eta') + \int_{\eta'}^{\eta_o} d\eta'' \partial_{\eta''} \psi^I(\eta'') \right) \partial_b \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') \right\} \\
& + \left[\left(\frac{2}{\mathcal{H}_s r_s} + \frac{\mathcal{H}'_s}{\mathcal{H}_s^2} \right) (v_{\parallel s} + \psi_s^I + 2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta')) - 3v_{\parallel s} + \frac{1}{\mathcal{H}_s} \partial_r v_{\parallel s} - 4\psi_s^I \right. \\
& \left. - 6 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') + \frac{4}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') - \frac{2}{r_s} \int_{\eta_s}^{\eta_o} d\eta' \frac{\eta' - \eta_s}{\eta_o - \eta'} \Delta_2 \psi^I(\eta') + \frac{1}{\mathcal{H}_s} \partial_{\eta} \psi_s^I \right] \delta_{\rho}^{(1)} \\
& + \left[\frac{1}{\rho} \partial_{\eta} (\bar{\rho} \delta_{\rho}^{(1)}) - \partial_r \delta_{\rho}^{(1)} \right] \frac{1}{\mathcal{H}_s} \left(-v_{\parallel s} - \psi_s^I - 2 \int_{\eta_s}^{\eta_o} d\eta' \partial_{\eta'} \psi^I(\eta') \right) \\
& + 2\partial_r \delta_{\rho}^{(1)} \int_{\eta_s}^{\eta_o} d\eta' \psi^I(\eta') - 2\partial_a \delta_{\rho}^{(1)} \int_{\eta_s}^{\eta_o} d\eta' \gamma_0^{ab} \partial_b \int_{\eta'}^{\eta_o} d\eta'' \psi^I(\eta'') + \delta_{\rho}^{(2)}. \tag{4.43}
\end{aligned}$$

Second Order

$$\Delta(\mathbf{n}, z) \equiv \frac{N(\mathbf{n}, z) - \langle N \rangle(z)}{\langle N \rangle(z)}$$

Leading terms

$$\begin{aligned} \Sigma^{(2)}(\mathbf{n}, z) &= \delta^{(2)} + \mathcal{H}^{-1} \partial_r^2 v^{(2)} + \mathcal{H}^{-2} (\partial_r^2 v)^2 + \mathcal{H}^{-2} \partial_r v \partial_r^3 v + \mathcal{H}^{-1} (\partial_r v \partial_r \delta + \partial_r^2 v \delta) \\ &\quad - 2\kappa^{(2)} - 2\delta\kappa + \nabla_a \delta \nabla^a \psi + \mathcal{H}^{-1} [-2(\partial_r^2 v)\kappa + \nabla_a (\partial_r^2 v) \nabla^a \psi] + 2\kappa^2 \\ &\quad - 2\nabla_b \kappa \nabla^b \psi - \frac{1}{2r(z)} \int_0^{r(z)} dr \frac{r(z) - r}{r} \Delta_2 (\nabla^b \Psi_1 \nabla_b \Psi_1) - 2 \int_0^{r(z)} \frac{dr}{r} \nabla^a \Psi_1 \nabla_a \kappa \end{aligned}$$

Second Order

$$\Delta(\mathbf{n}, z) \equiv \frac{N(\mathbf{n}, z) - \langle N \rangle(z)}{\langle N \rangle(z)}$$

Leading terms

density-RSD

$$\begin{aligned} \Sigma^{(2)}(\mathbf{n}, z) = & \delta^{(2)} + \mathcal{H}^{-1} \partial_r^2 v^{(2)} + \mathcal{H}^{-2} (\partial_r^2 v)^2 + \mathcal{H}^{-2} \partial_r v \partial_r^3 v + \mathcal{H}^{-1} (\partial_r v \partial_r \delta + \partial_r^2 v \delta) \\ & - 2\kappa^{(2)} - 2\delta\kappa + \nabla_a \delta \nabla^a \psi + \mathcal{H}^{-1} [-2(\partial_r^2 v)\kappa + \nabla_a (\partial_r^2 v) \nabla^a \psi] + 2\kappa^2 \\ & - 2\nabla_b \kappa \nabla^b \psi - \frac{1}{2r(z)} \int_0^{r(z)} dr \frac{r(z) - r}{r} \Delta_2 (\nabla^b \Psi_1 \nabla_b \Psi_1) - 2 \int_0^{r(z)} \frac{dr}{r} \nabla^a \Psi_1 \nabla_a \kappa \end{aligned}$$

Second Order

$$\Delta(\mathbf{n}, z) \equiv \frac{N(\mathbf{n}, z) - \langle N \rangle(z)}{\langle N \rangle(z)}$$

$$Z_2(\mathbf{k}_1, \mathbf{k}_2) = \frac{b_2}{2} + b_1 F_2(\mathbf{k}_1, \mathbf{k}_2) + f\mu^2 G_2(\mathbf{k}_1, \mathbf{k}_2) + \frac{f\mu k}{2} \left[\frac{\mu_1}{k_1} (b_1 + f\mu_2^2) + \frac{\mu_2}{k_2} (b_1 + f\mu_1^2) \right]$$

density-RSD

$$\begin{aligned} \Sigma^{(2)}(\mathbf{n}, z) = & \delta^{(2)} + \mathcal{H}^{-1} \partial_r^2 v^{(2)} + \mathcal{H}^{-2} (\partial_r^2 v)^2 + \mathcal{H}^{-2} \partial_r v \partial_r^3 v + \mathcal{H}^{-1} (\partial_r v \partial_r \delta + \partial_r^2 v \delta) \\ & - 2\kappa^{(2)} - 2\delta\kappa + \nabla_a \delta \nabla^a \psi + \mathcal{H}^{-1} [-2(\partial_r^2 v)\kappa + \nabla_a (\partial_r^2 v) \nabla^a \psi] + 2\kappa^2 \\ & - 2\nabla_b \kappa \nabla^b \psi - \frac{1}{2r(z)} \int_0^{r(z)} dr \frac{r(z) - r}{r} \Delta_2 (\nabla^b \Psi_1 \nabla_b \Psi_1) - 2 \int_0^{r(z)} \frac{dr}{r} \nabla^a \Psi_1 \nabla_a \kappa \end{aligned}$$

Second Order

$$\Delta(\mathbf{n}, z) \equiv \frac{N(\mathbf{n}, z) - \langle N \rangle(z)}{\langle N \rangle(z)}$$

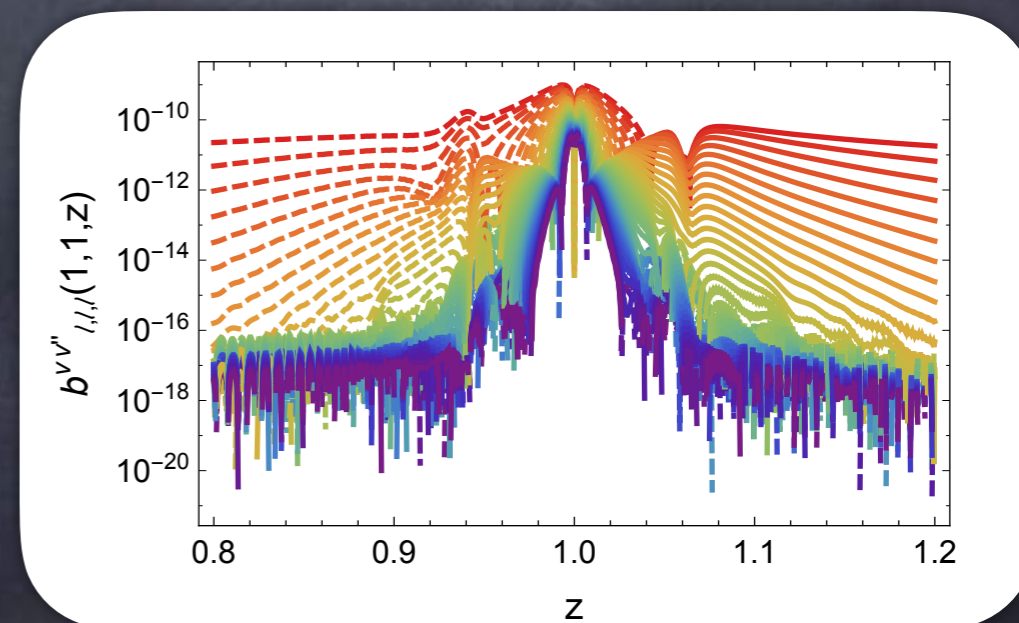
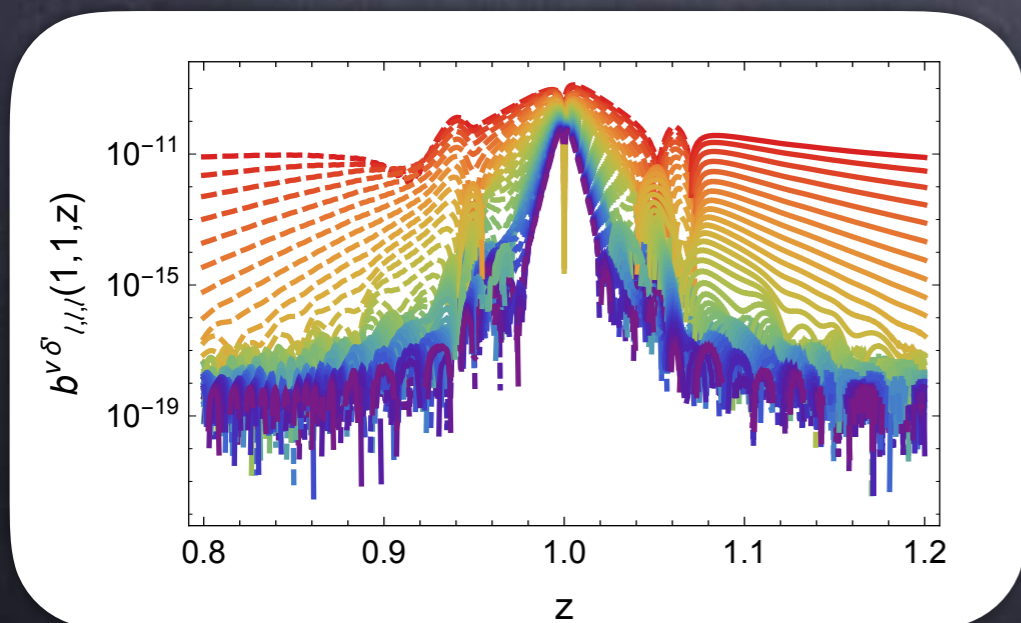
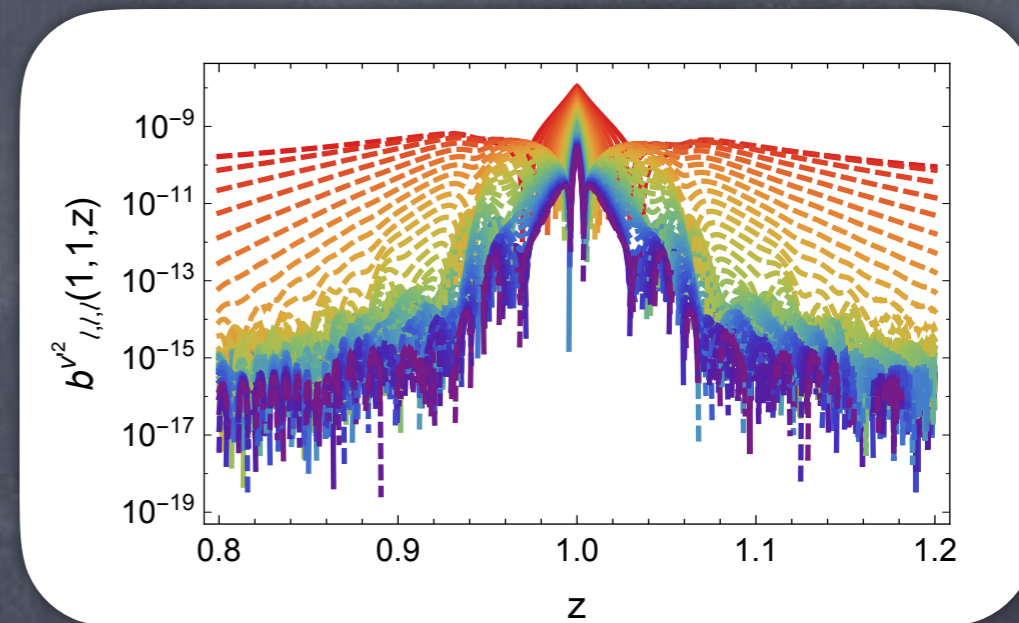
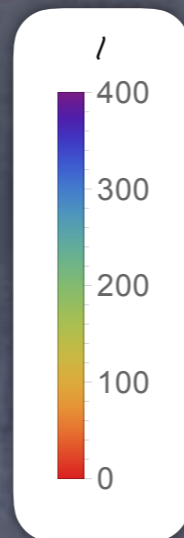
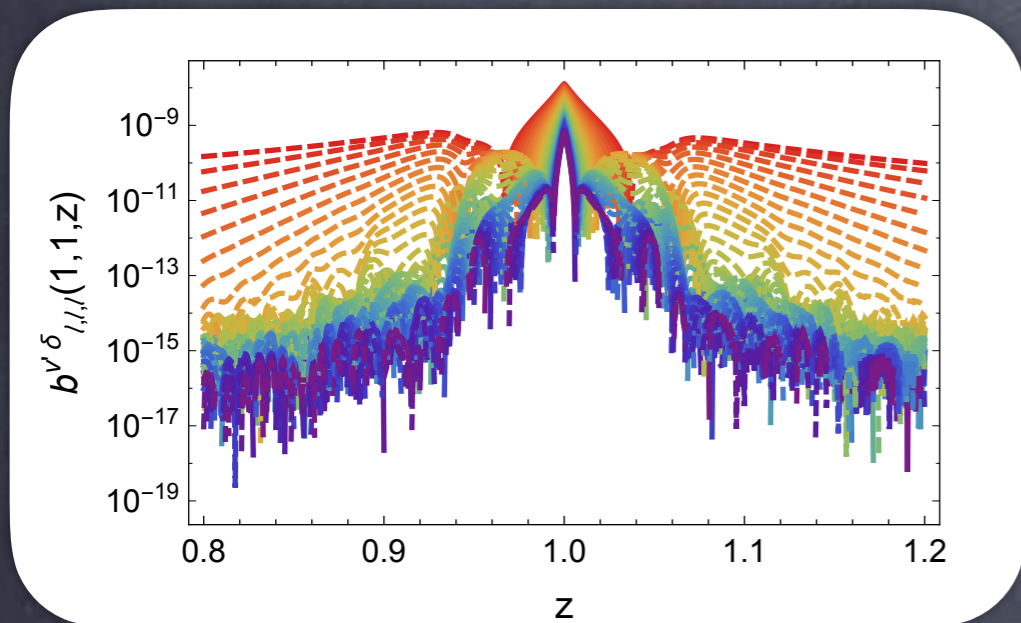
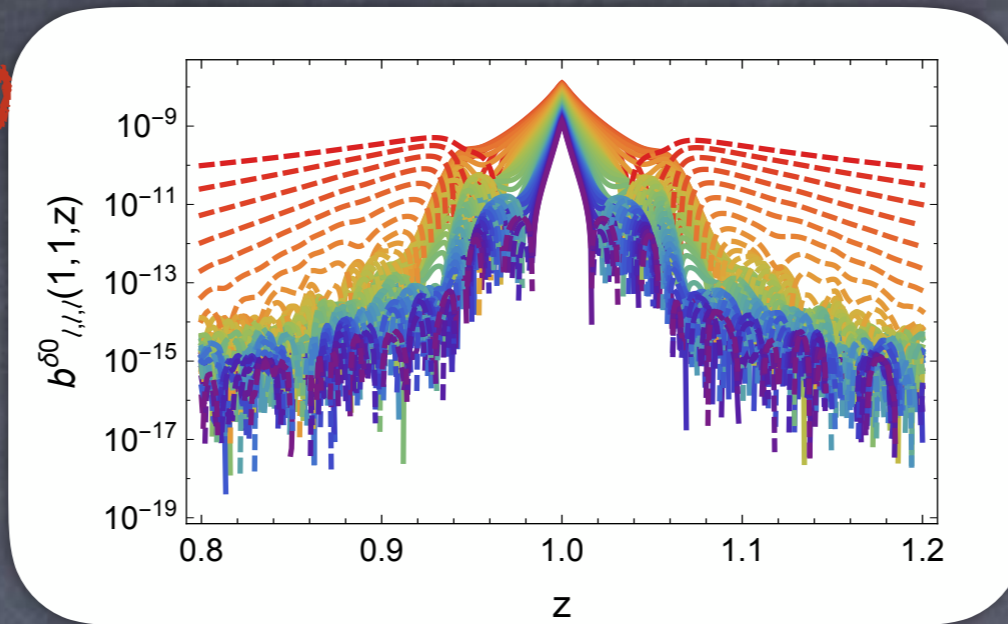
Leading terms

$$\Sigma^{(2)}(\mathbf{n}, z) = \delta^{(2)} + \mathcal{H}^{-1} \partial_r^2 v^{(2)} + \mathcal{H}^{-2} (\partial_r^2 v)^2 + \mathcal{H}^{-2} \partial_r v \partial_r^3 v + \mathcal{H}^{-1} (\partial_r v \partial_r \delta + \partial_r^2 v \delta)$$

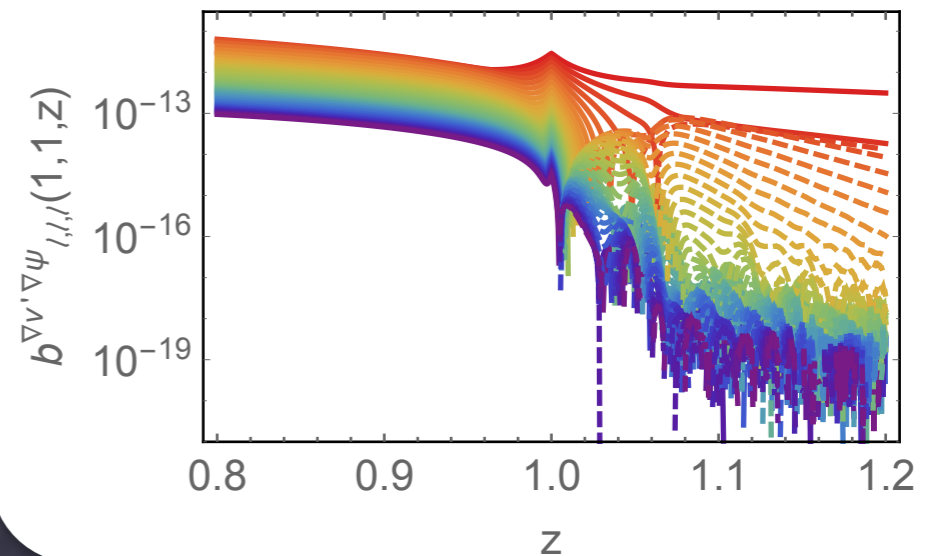
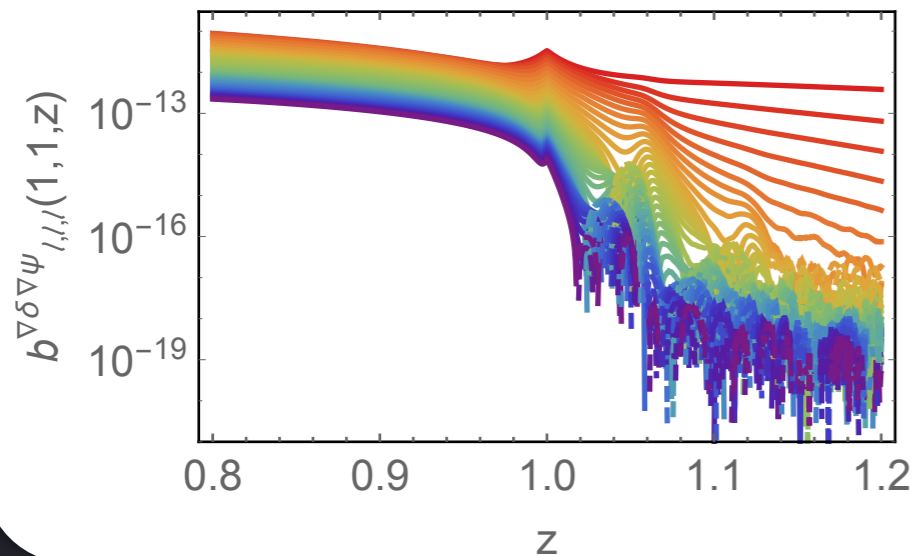
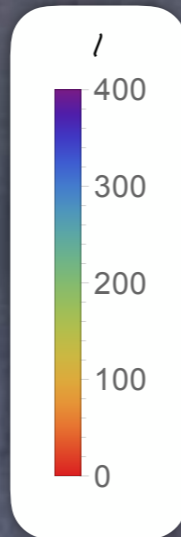
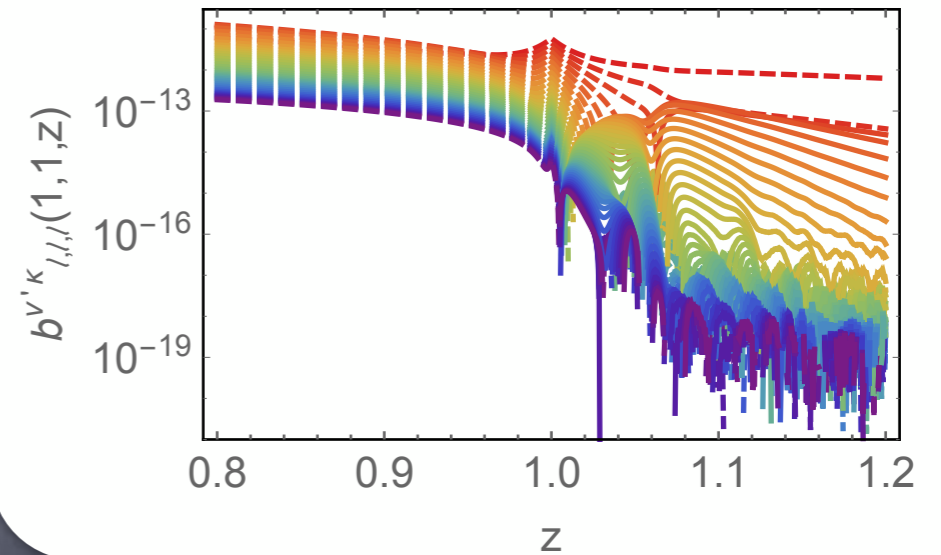
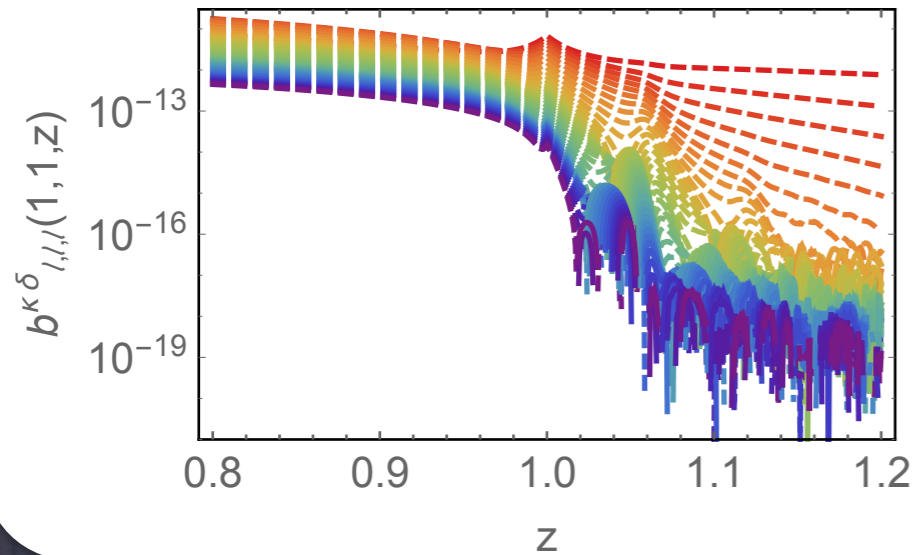
$$\begin{aligned} & - 2\kappa^{(2)} - 2\delta\kappa + \nabla_a \delta \nabla^a \psi + \mathcal{H}^{-1} [-2(\partial_r^2 v)\kappa + \nabla_a (\partial_r^2 v) \nabla^a \psi] + 2\kappa^2 \\ & - 2\nabla_b \kappa \nabla^b \psi - \frac{1}{2r(z)} \int_0^{r(z)} dr \frac{r(z) - r}{r} \Delta_2 (\nabla^b \Psi_1 \nabla_b \Psi_1) - 2 \int_0^{r(z)} \frac{dr}{r} \nabla^a \Psi_1 \nabla_a \kappa \end{aligned}$$

like-lensing terms

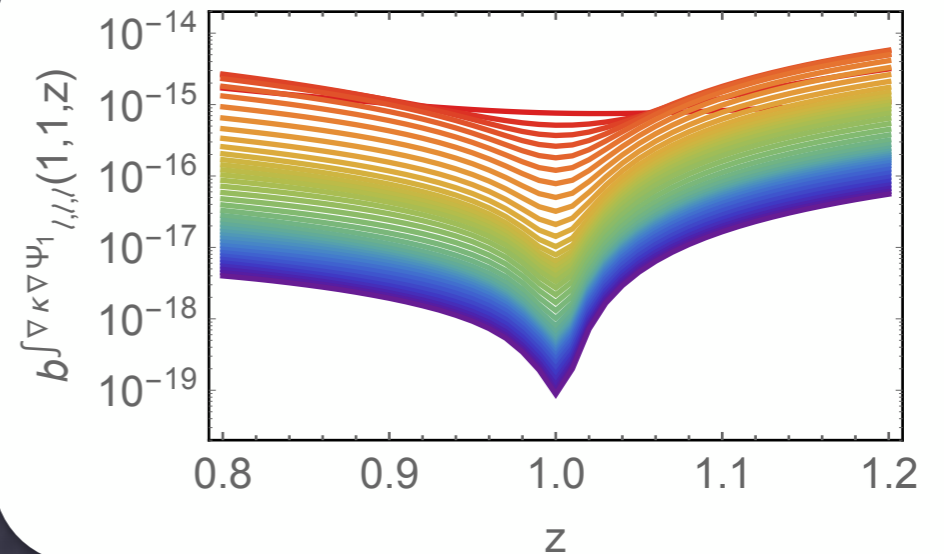
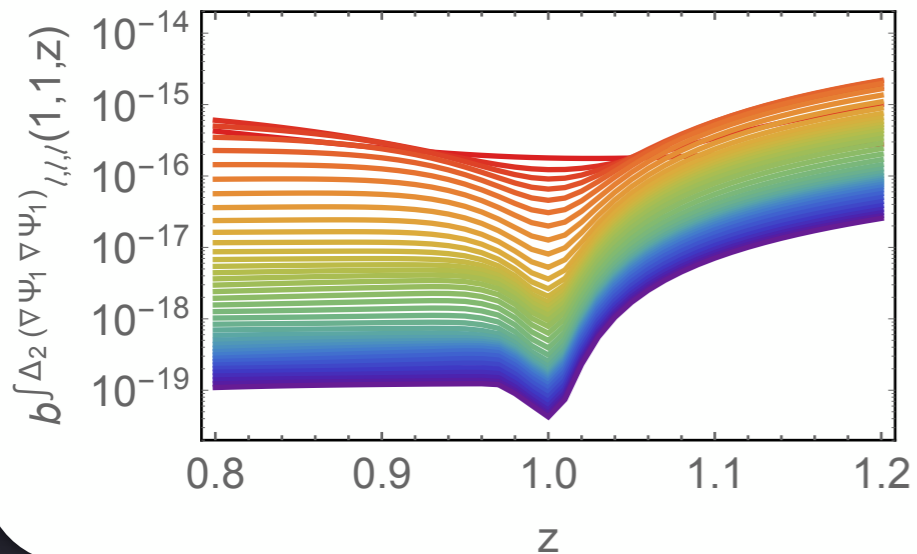
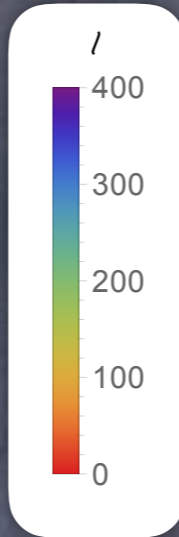
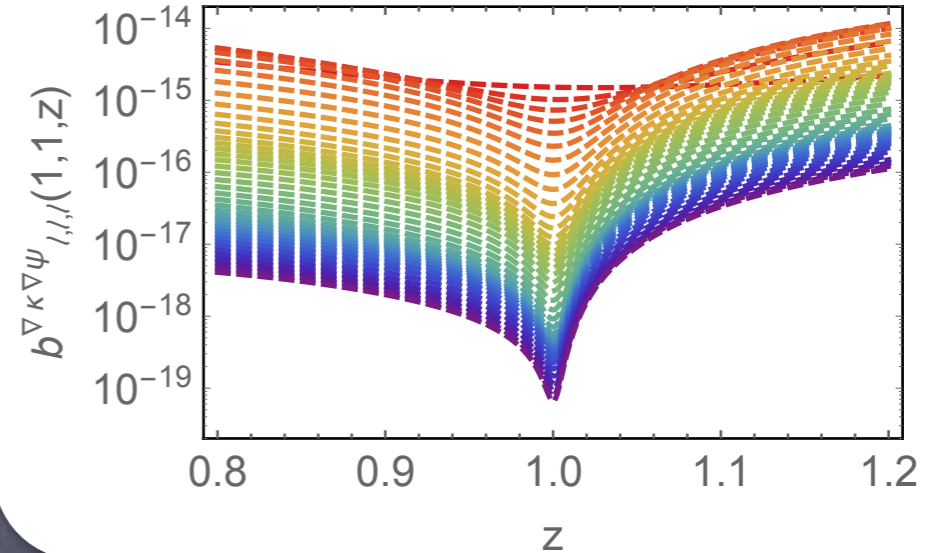
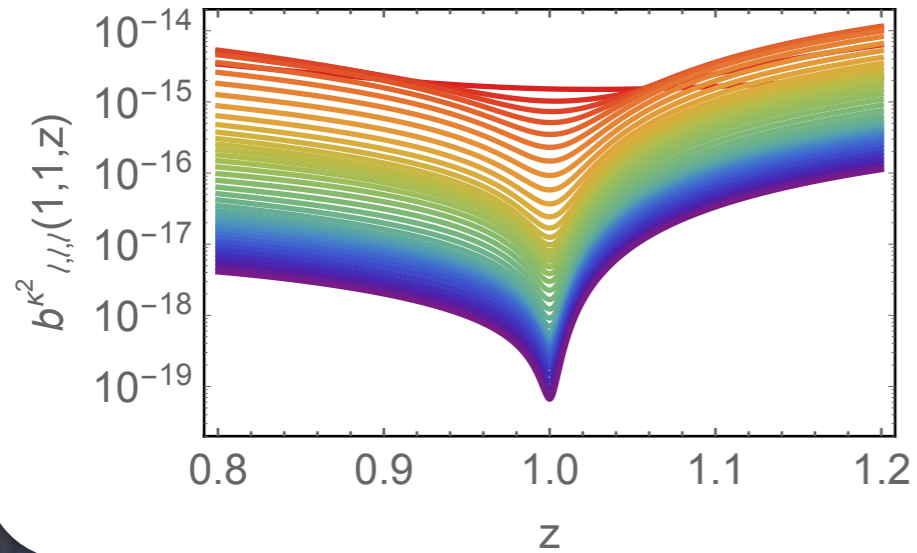
density-RSD



lensing x newtonian

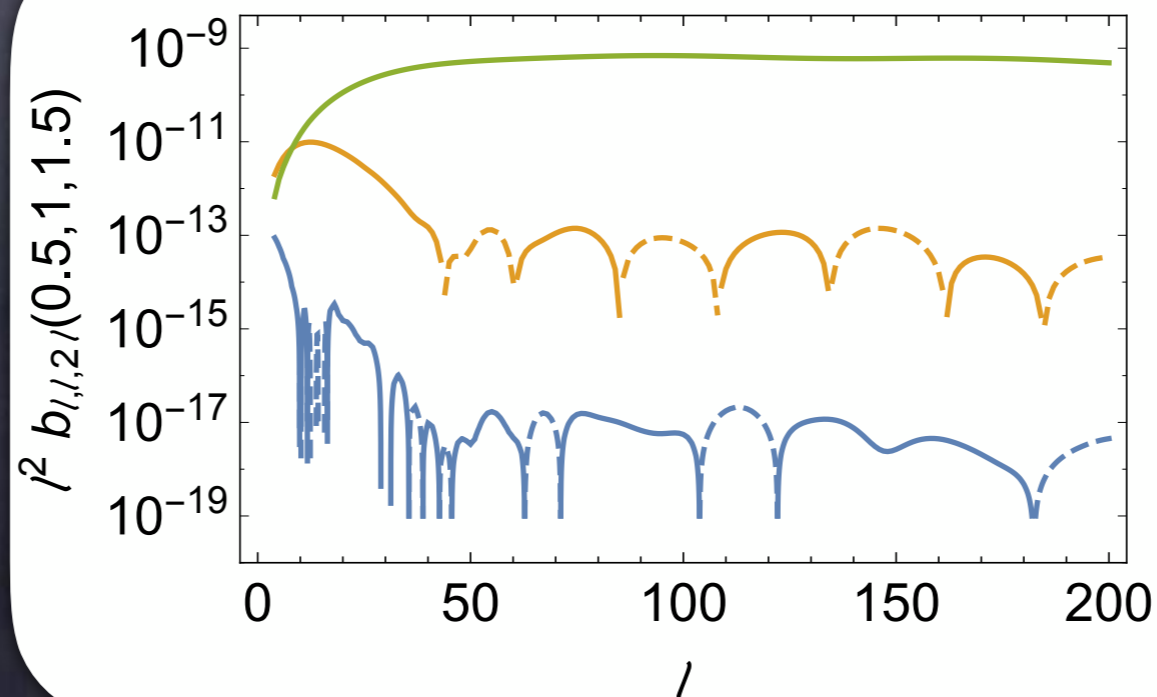
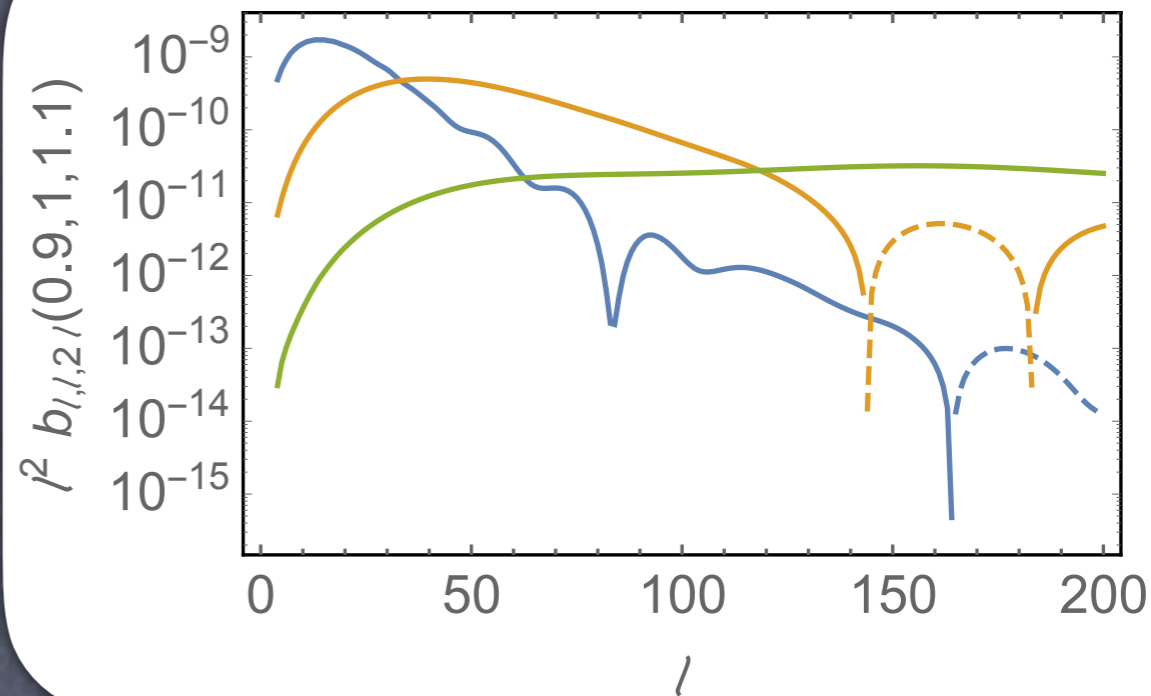
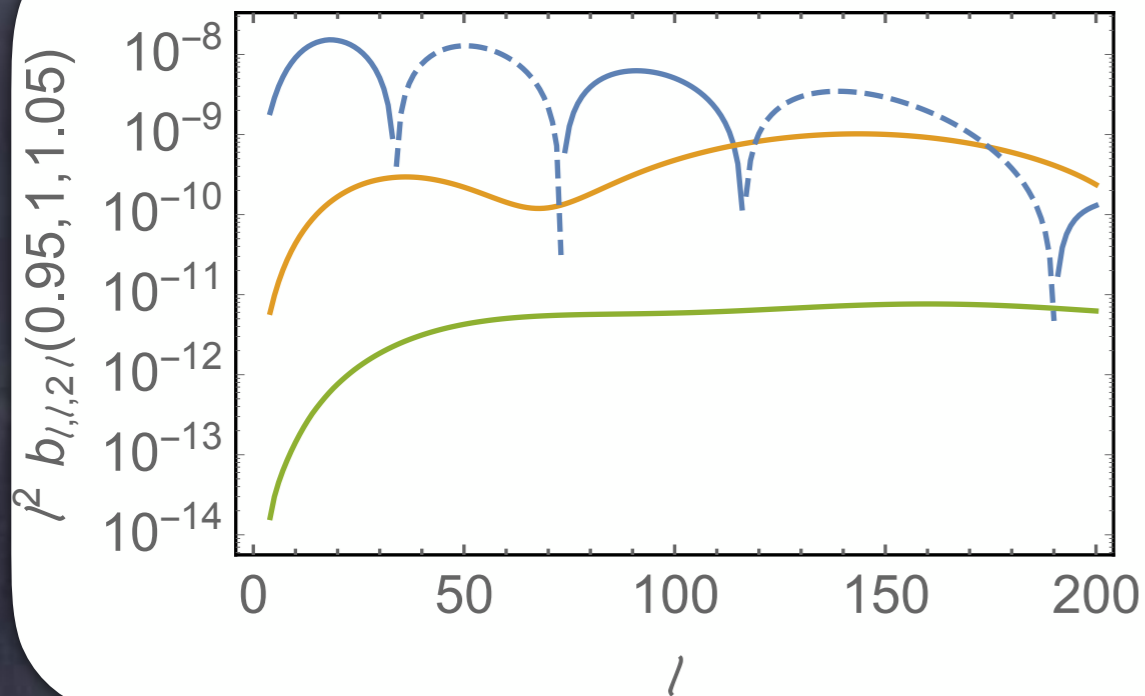


Lensing



Beyond
Born-Approximation

Bispectrum



- Newtonian
- Newtonian x Lensing
- Lensing

Conclusions

- The observable quantity includes several relativistic corrections
 - which encode useful information
 - and may bias the parameter estimation, if neglected
- Multi-tracers technique allows to isolate some relativistic correction. Future surveys should be able to measure them
- relativistic number counts can be generalized to non-flat universes
 - neglecting relativistic effects may bias the curvature parameter as well
- Many new contributions to the bispectrum
 - how do they affect f_{NL} ?