

# Thermal loophole in the Higgs Portal

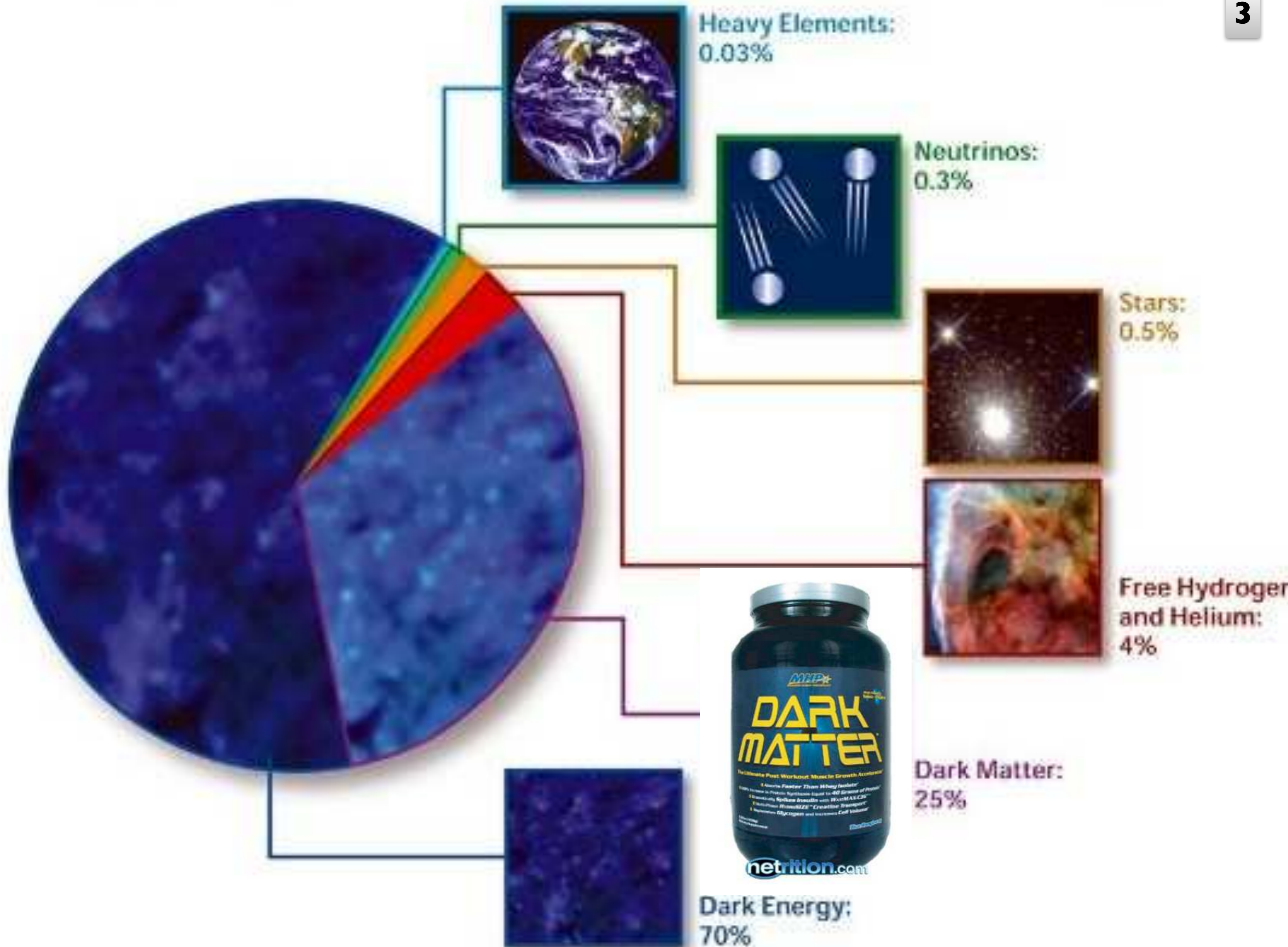
**Andi Hektor**  
**NICPB, Tallinn**

**in collaboration with**  
**Kristjan Kannike & Ville Vaskonen**

**Helsinki ✧ 7/9/2017**

# Motivation & outlines

- Higgs Portal, rise and fall
- Escapes from the direct detection constraints:
  - ‘Pseudoscalar portal’
  - ‘Resonance portal’
- **‘Thermal portal’**
- Thermal effects of primordial plasma at the freeze-out
- A minimal model
- Conclusions



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## Trump searches for the “dark matter” of politics

Details Last Updated on Tuesday, 14 June 2016 11:12  
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It exists mostly in theory and the realm of complicated mathematical equations. But, scientists know it's there. It's called "dark matter" and it makes up a substantial part of our universe. Politics, of course, isn't physics, but in the theater of American elections there is a sort of dark matter as well. We can't always trace its source, sometimes it's completely hidden, but on occasion, it's more evident than we would like to believe. And this year, the dark matter of American politics isn't hard to find at all. It's that undercurrent of hostility that stays just below the surface. Most of us admit it. come on. have thought some of the outlandish things Donald



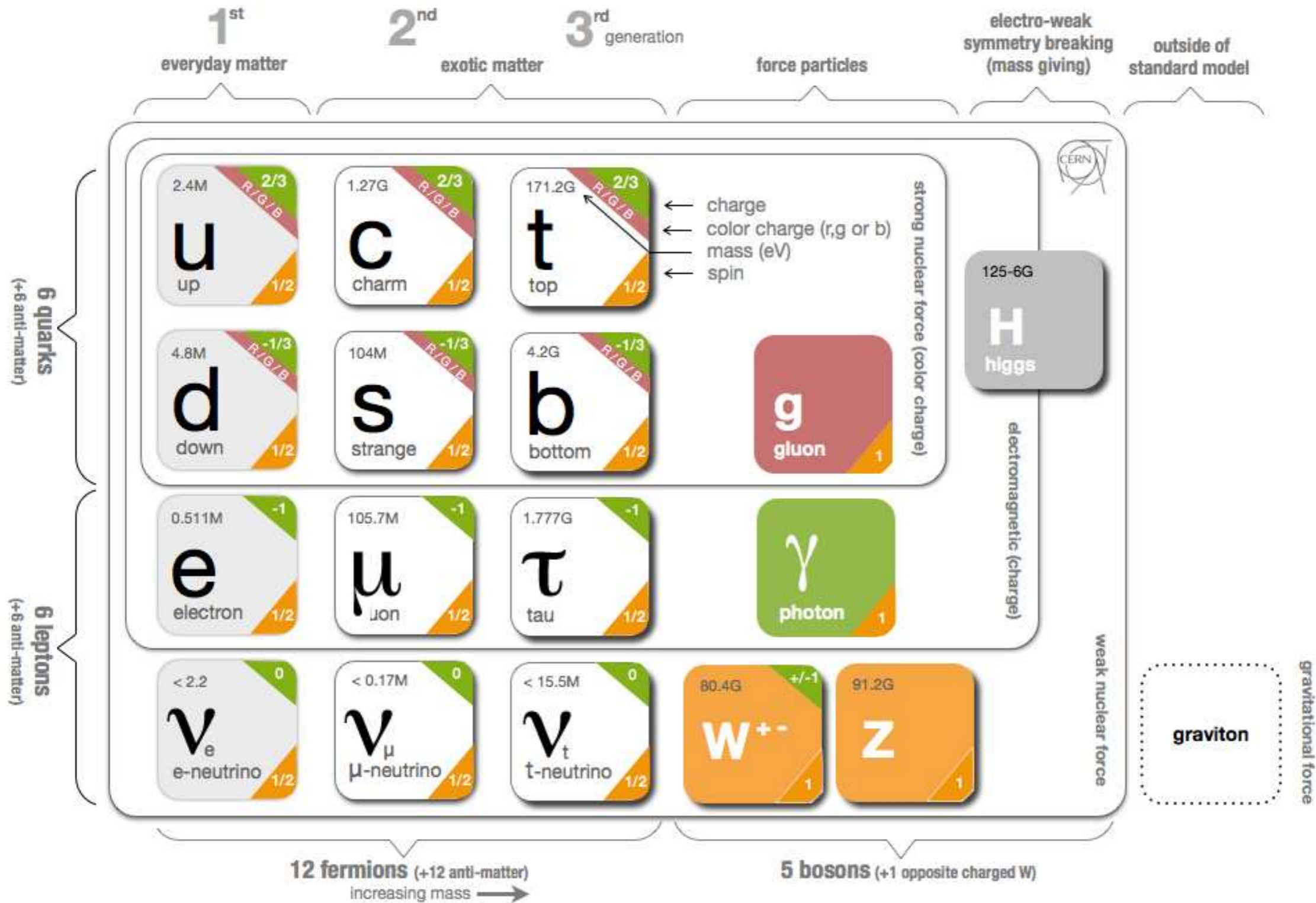
Looking for a Way to Save  
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is the Solution!

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sell ads and get

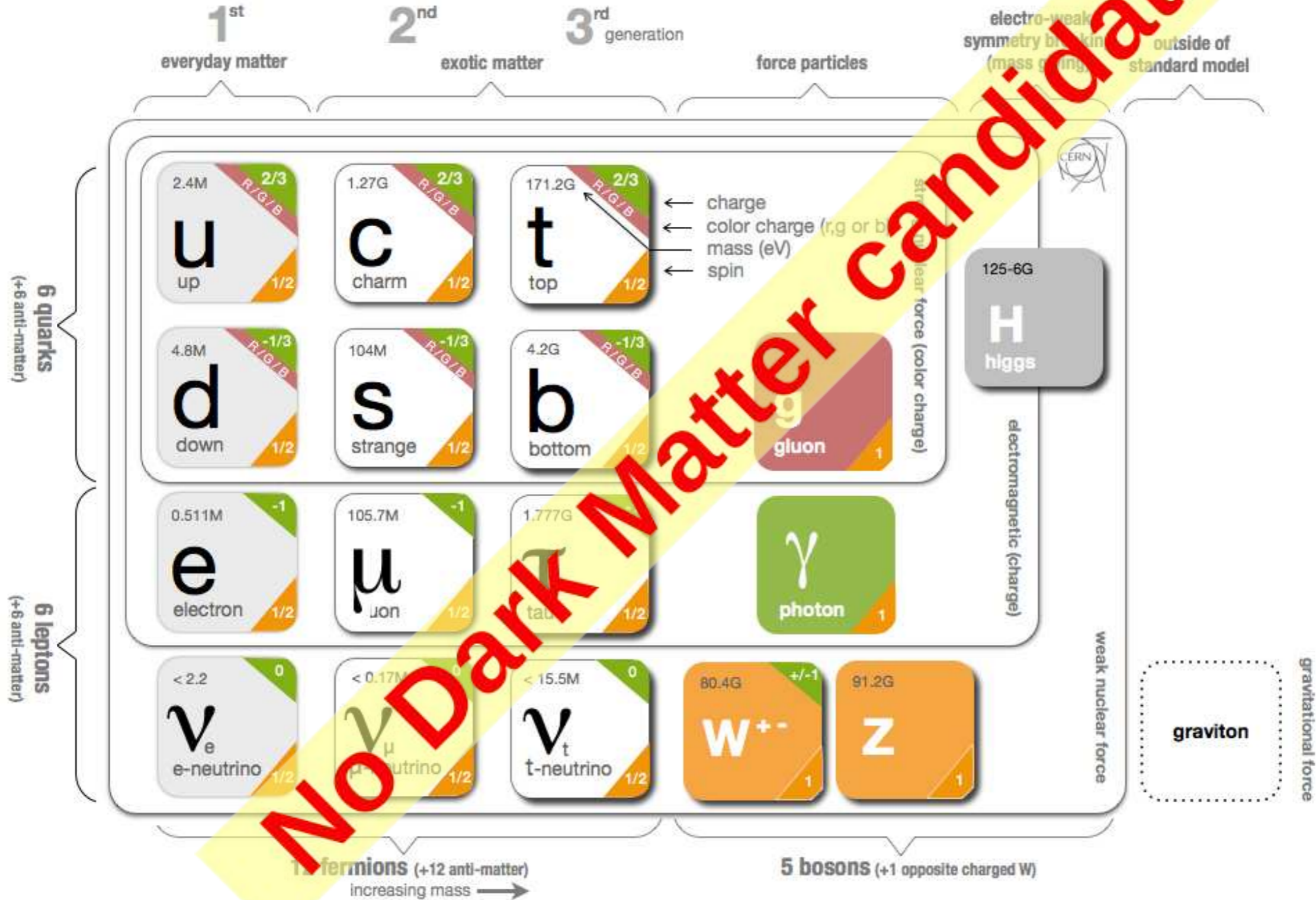


# DM in the Standard Model?



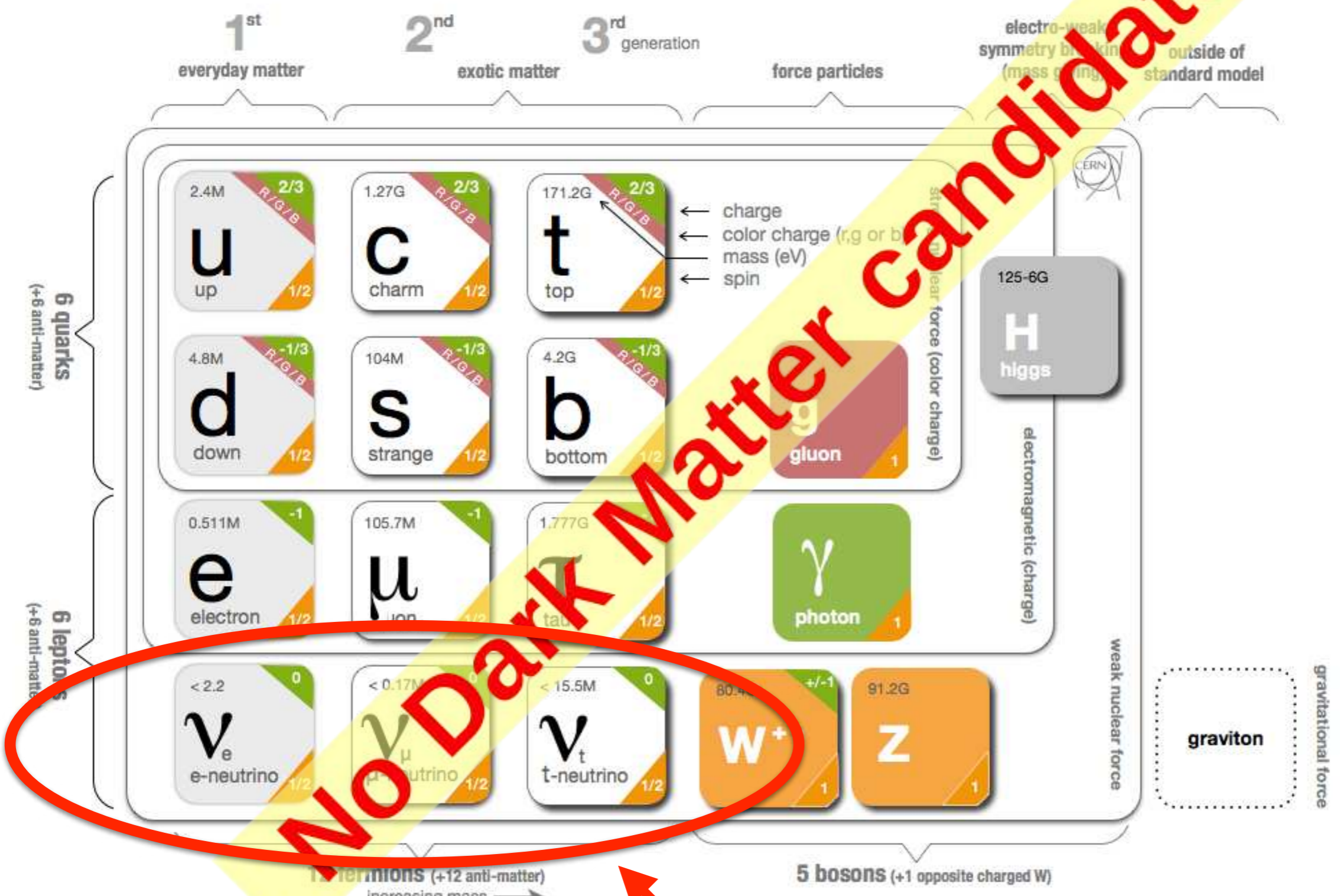


# DM in the Standard Model?



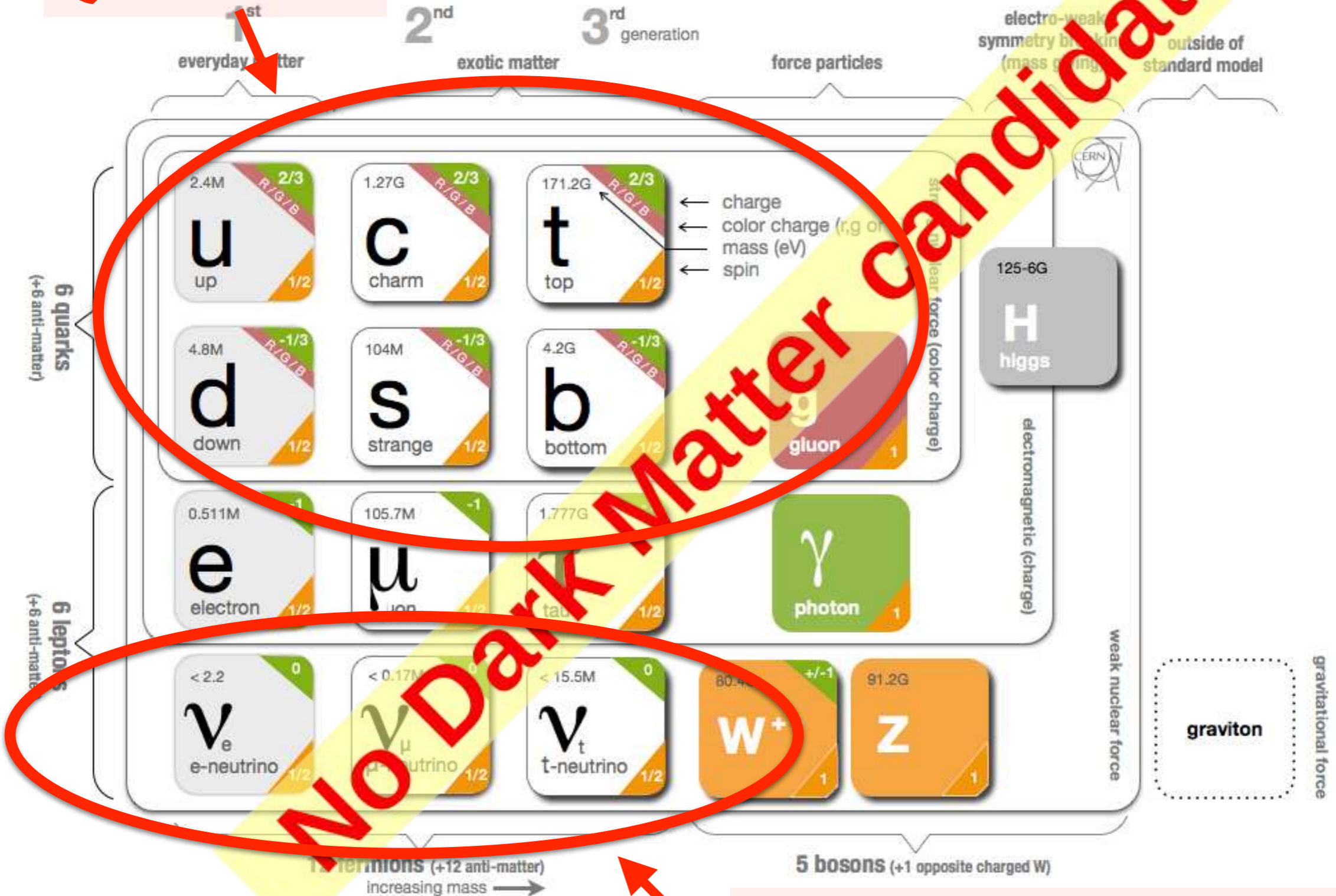


# DM in the Standard Model?





# Stop! DM in the Standard Model? QCD axion?

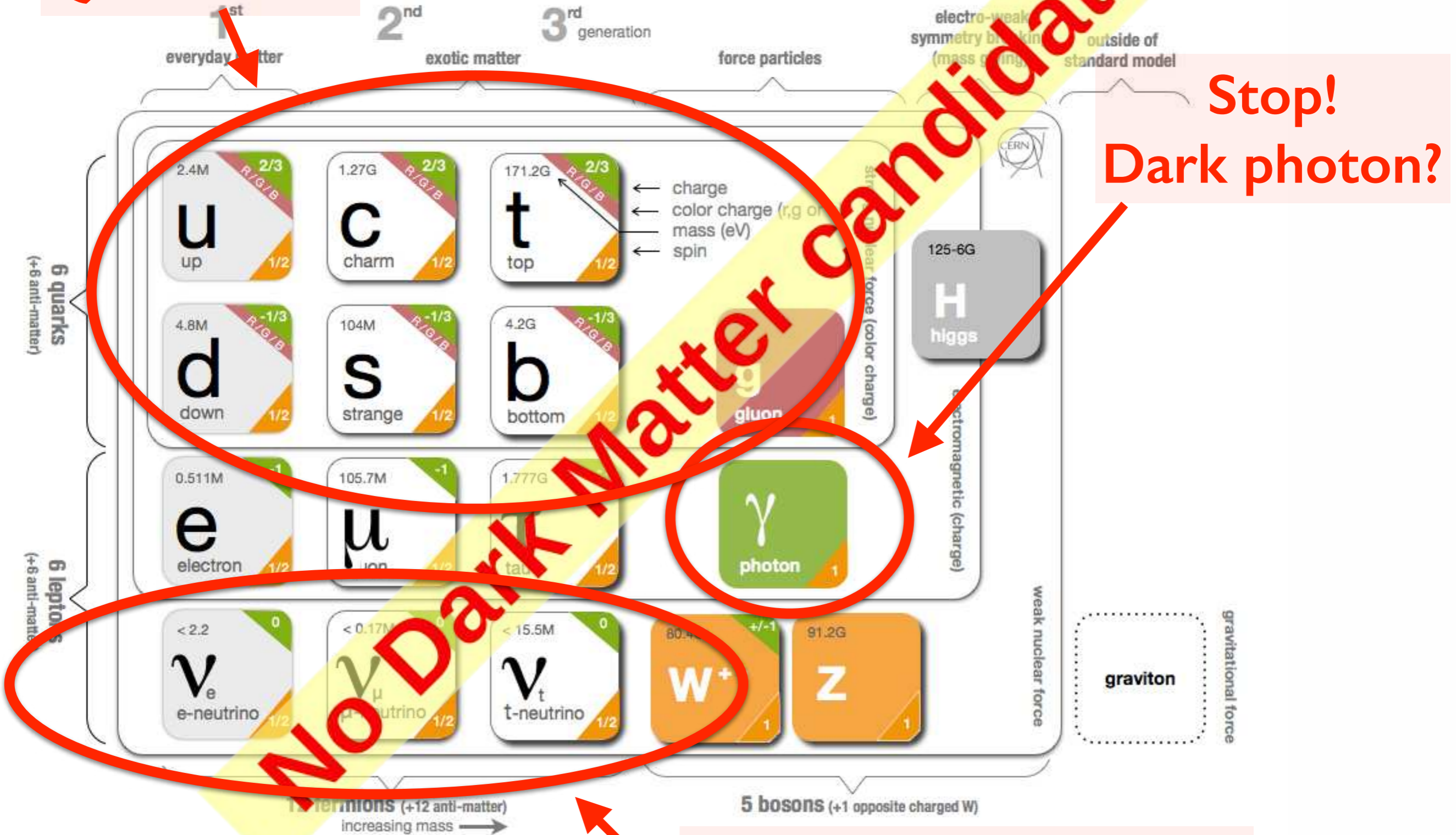


Stop! Sterile neutrino?





# Stop! DM in the Standard Model? QCD axion?



Stop!  
Dark photon?

Stop! Sterile neutrino?



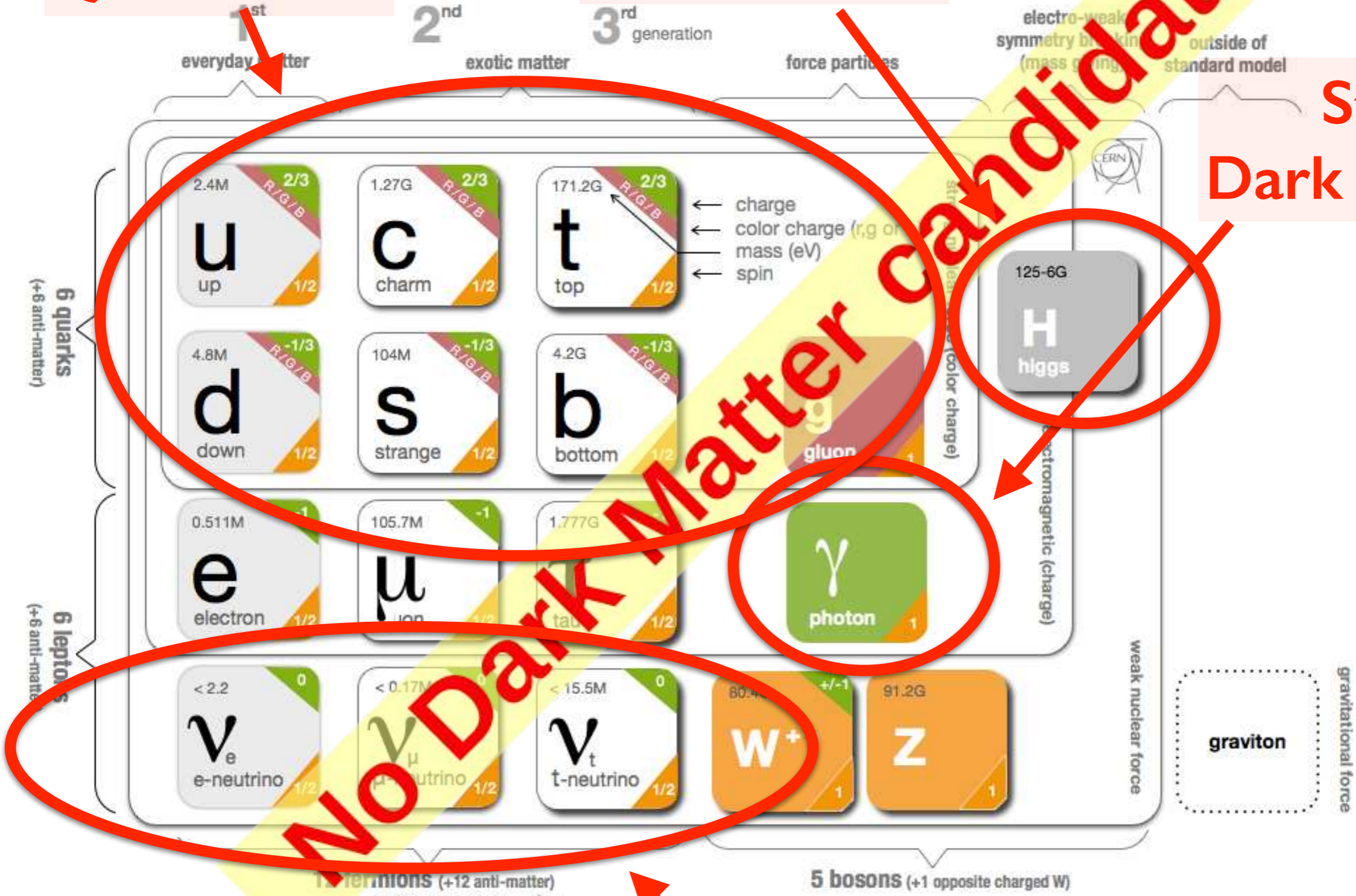
# Stop! DM in the Standard Model?

QCD axion?

Higgs portal?

Stop! Dark photon?

Stop! Sterile neutrino?



NO Dark Matter candidate!



# Stop! DM in the Standard Model?

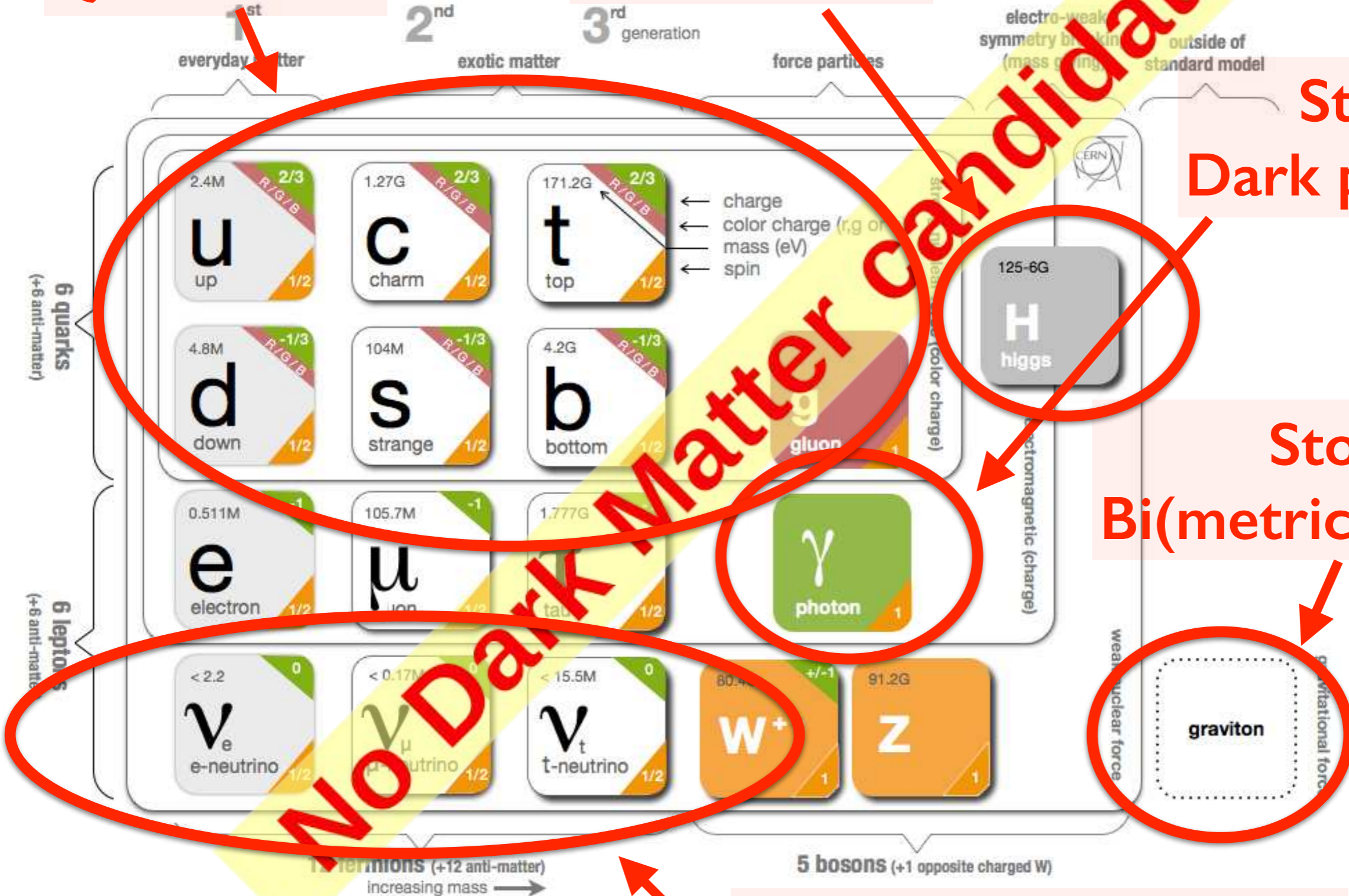
QCD axion?

Higgs portal?

Stop! Dark photon?

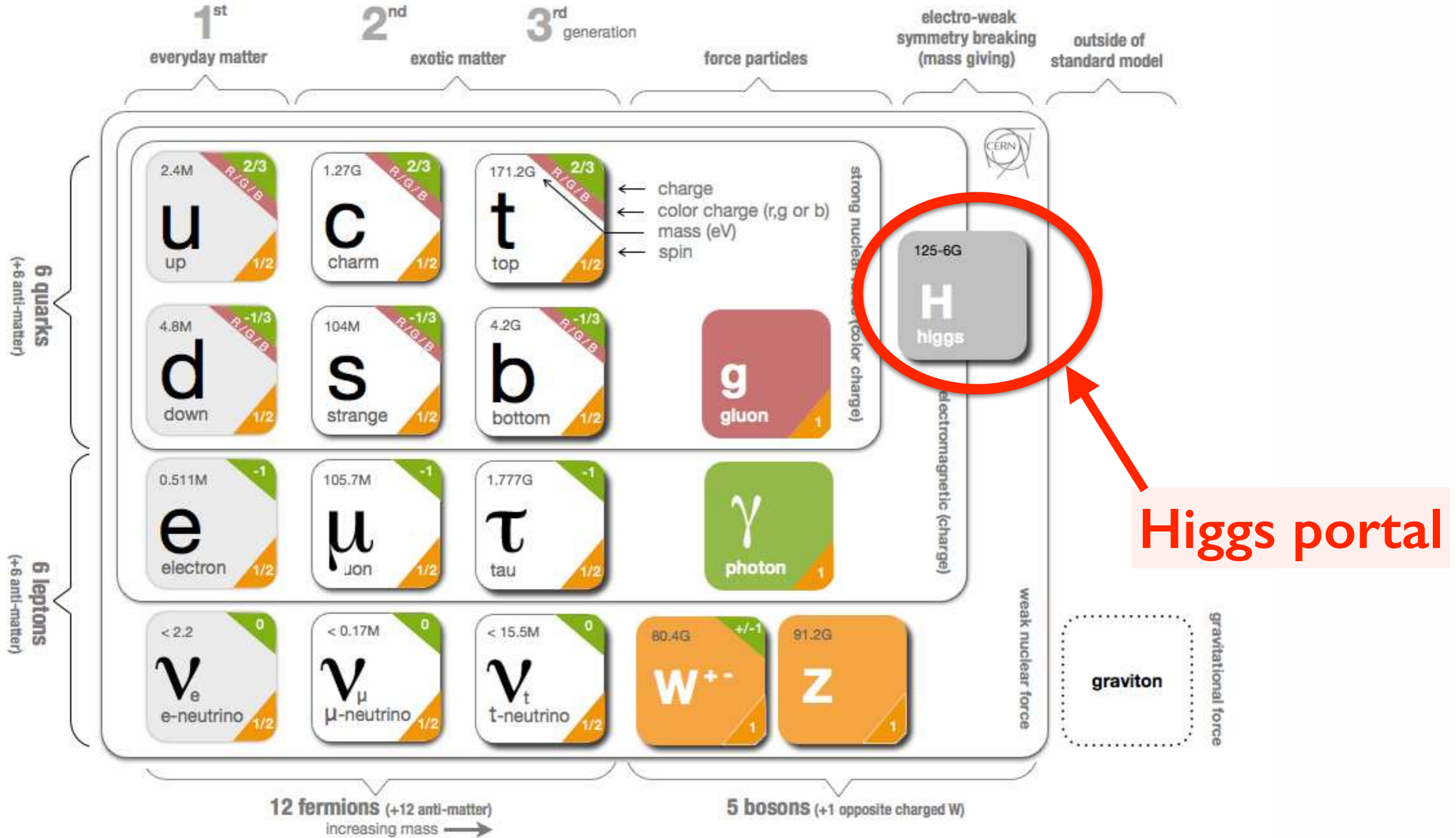
Stop! Bi(metric)gravity?

Stop! Sterile neutrino?



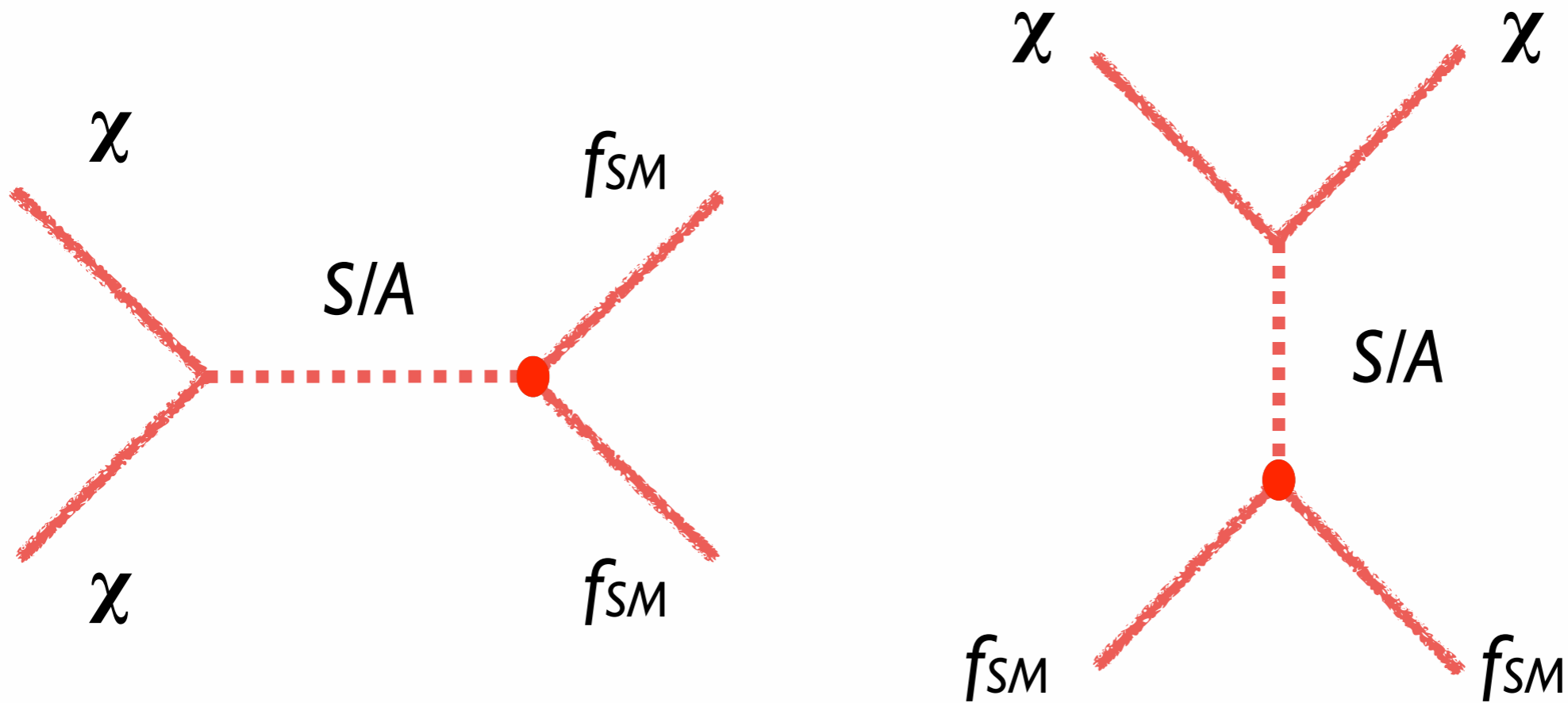


# In this talk, the favourite is...





# The (simplified) portal DM



$$\mathcal{L}_{0_S S^{\frac{1}{2}}} = \frac{1}{2}(\partial_\mu S)^2 - \frac{1}{2}m_S^2 S^2 + \bar{\chi}(i\partial - m_\chi)\chi - g_\chi S \bar{\chi}\chi - g_{SM} S \sum_f \frac{y_f}{\sqrt{2}} \bar{f}f$$

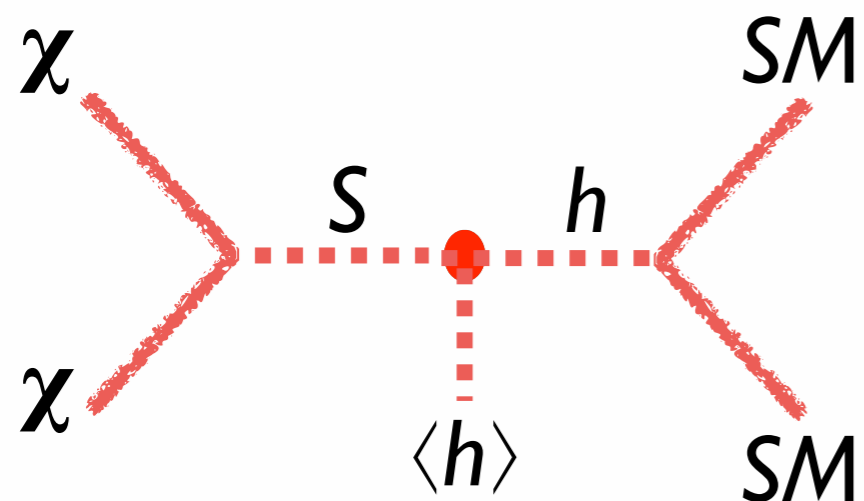
$$\mathcal{L}_{0_A S^{\frac{1}{2}}} = \frac{1}{2}(\partial_\mu A)^2 - \frac{1}{2}m_A^2 A^2 + \bar{\chi}(i\partial - m_\chi)\chi - ig_\chi A \bar{\chi}\gamma^5\chi - ig_{SM} A \sum_f \frac{y_f}{\sqrt{2}} \bar{f}\gamma^5 f$$

The portal provides  
a good old WIMPy scenario!

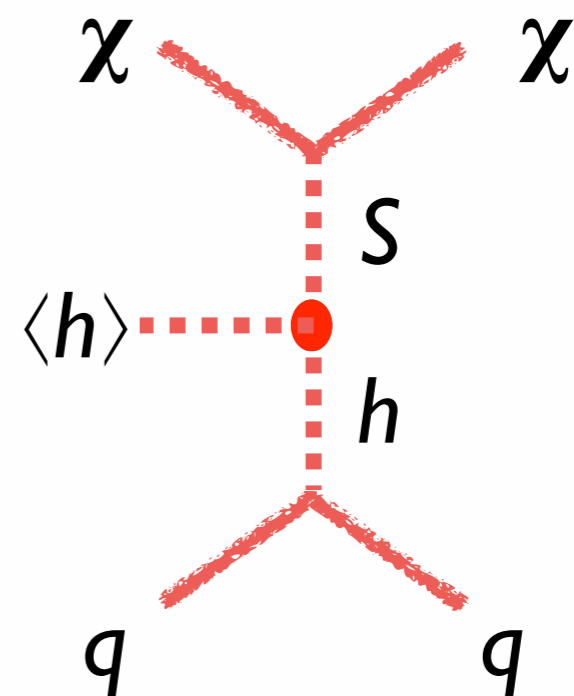


# Higgs portal, a realization

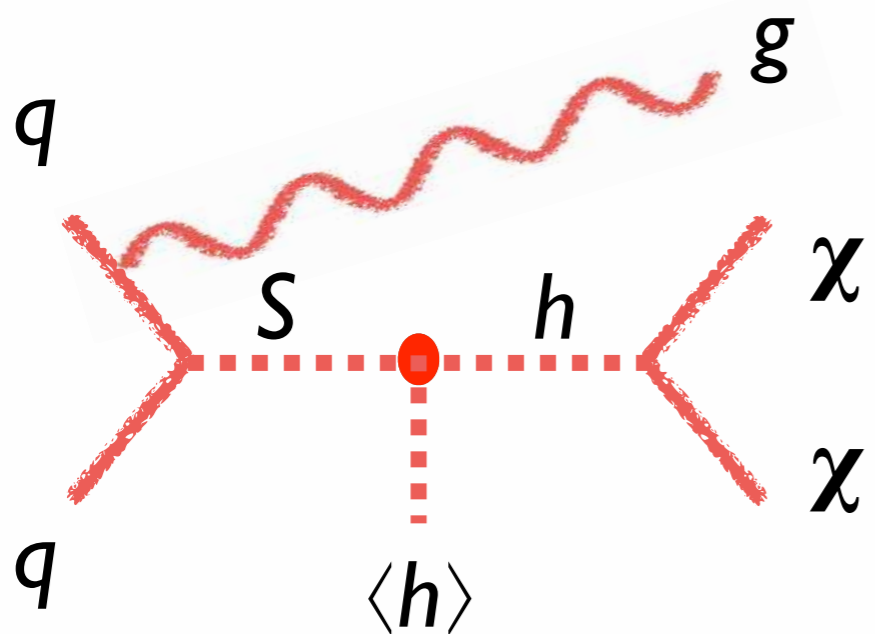
Freeze-out



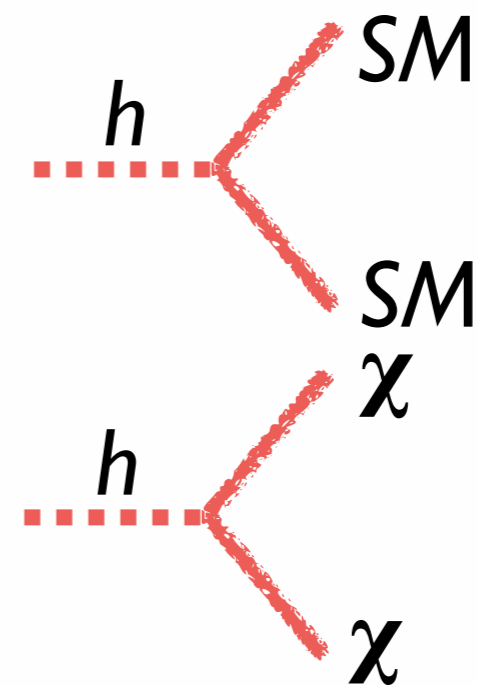
Direct detection



Accelerator: mono-jet

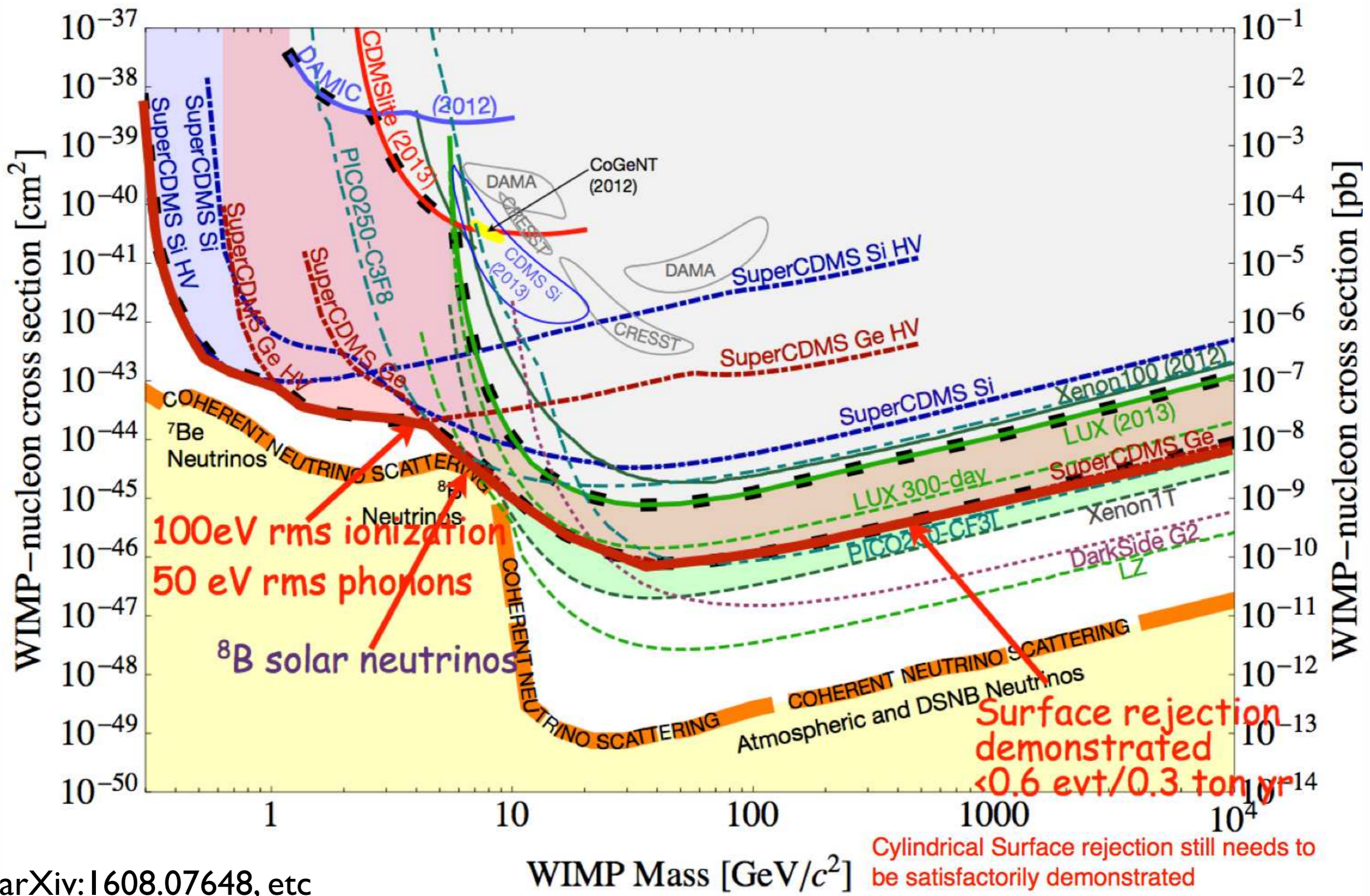


Accelerator: the Higgs width





# The model killer, Direct Detection

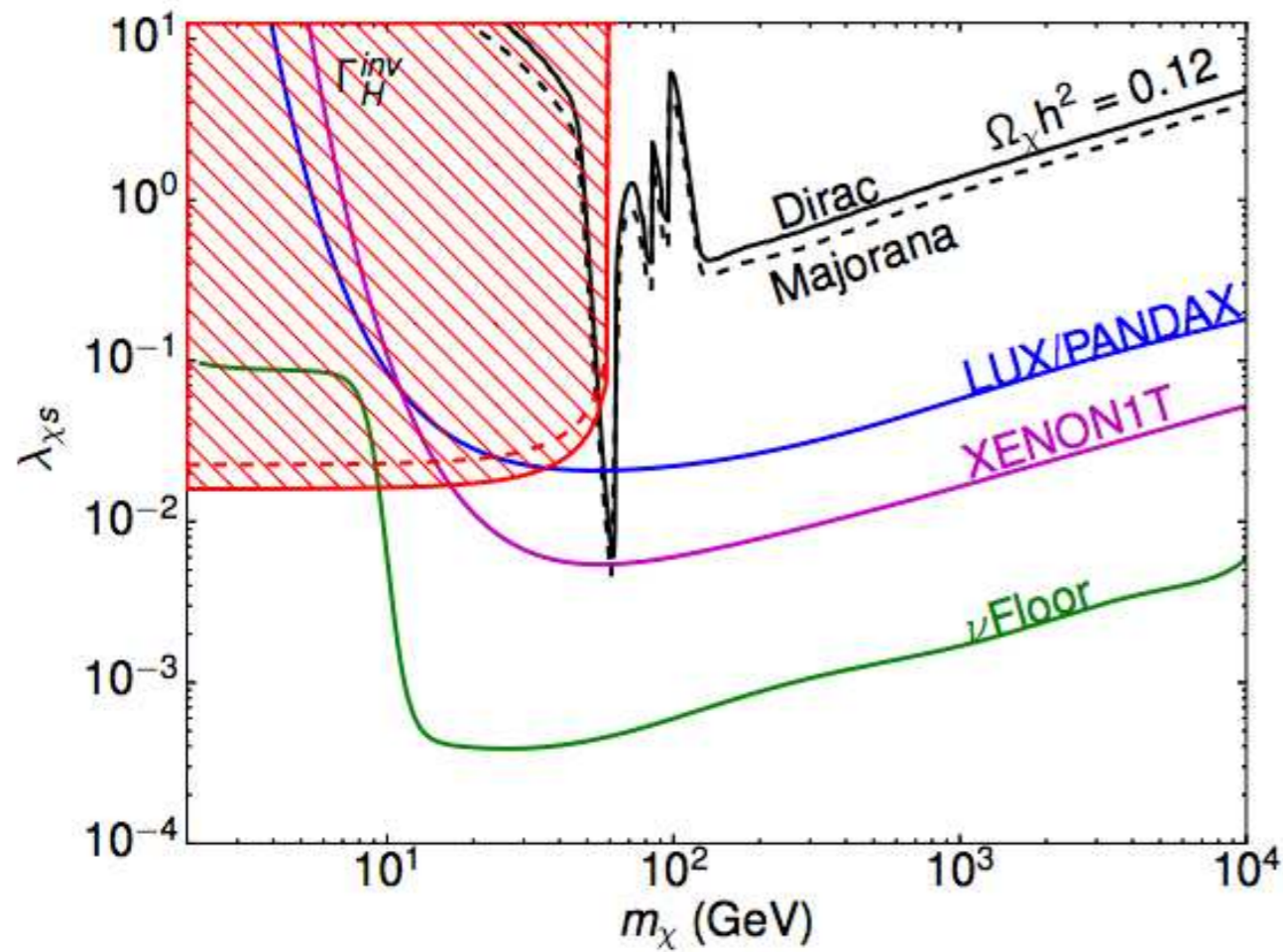


See arXiv:1608.07648, etc

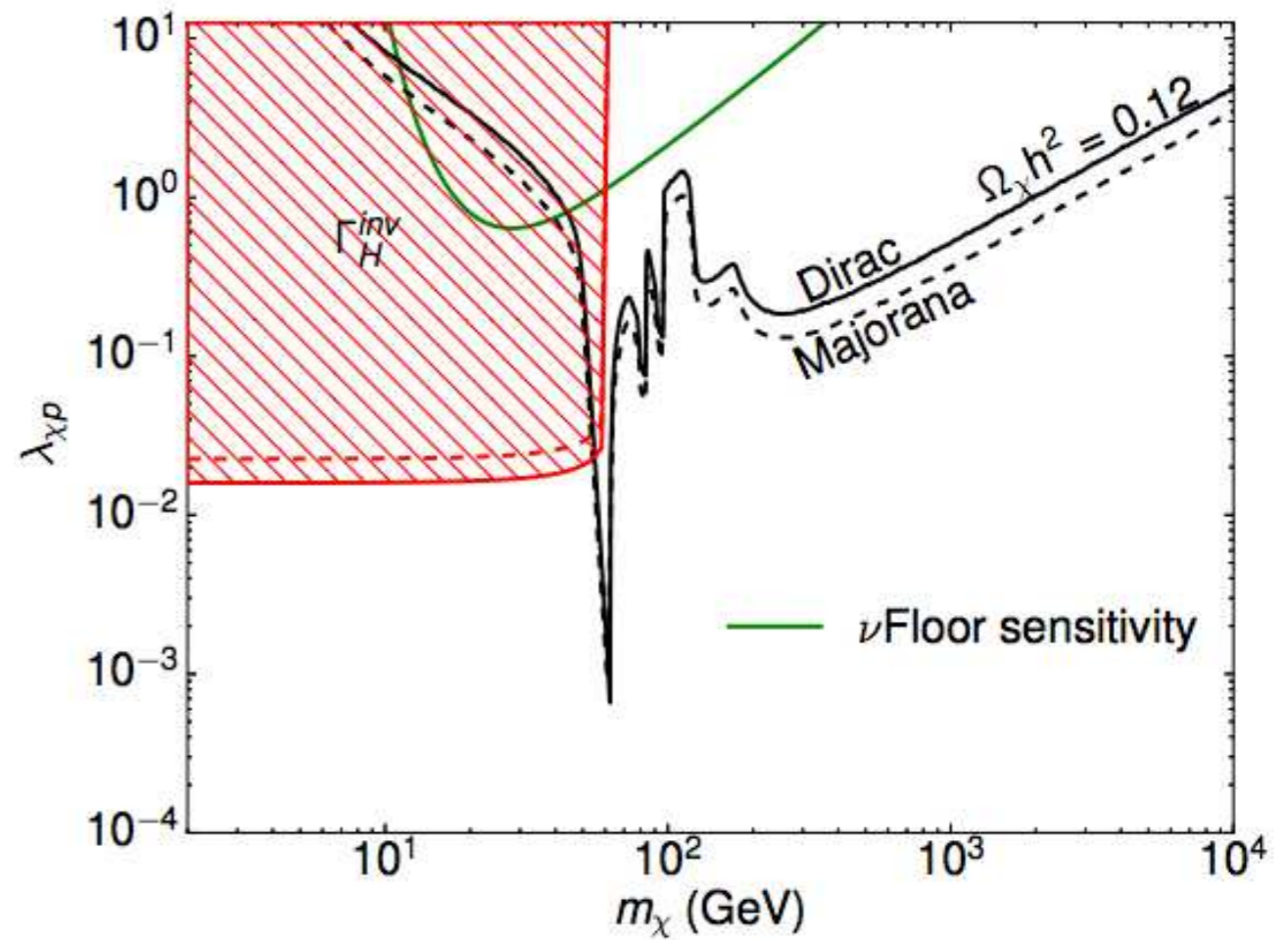


So, after LUX/Pandax there is just one problem —  
the model is ruled out?

Scalar



Pseudoscalar



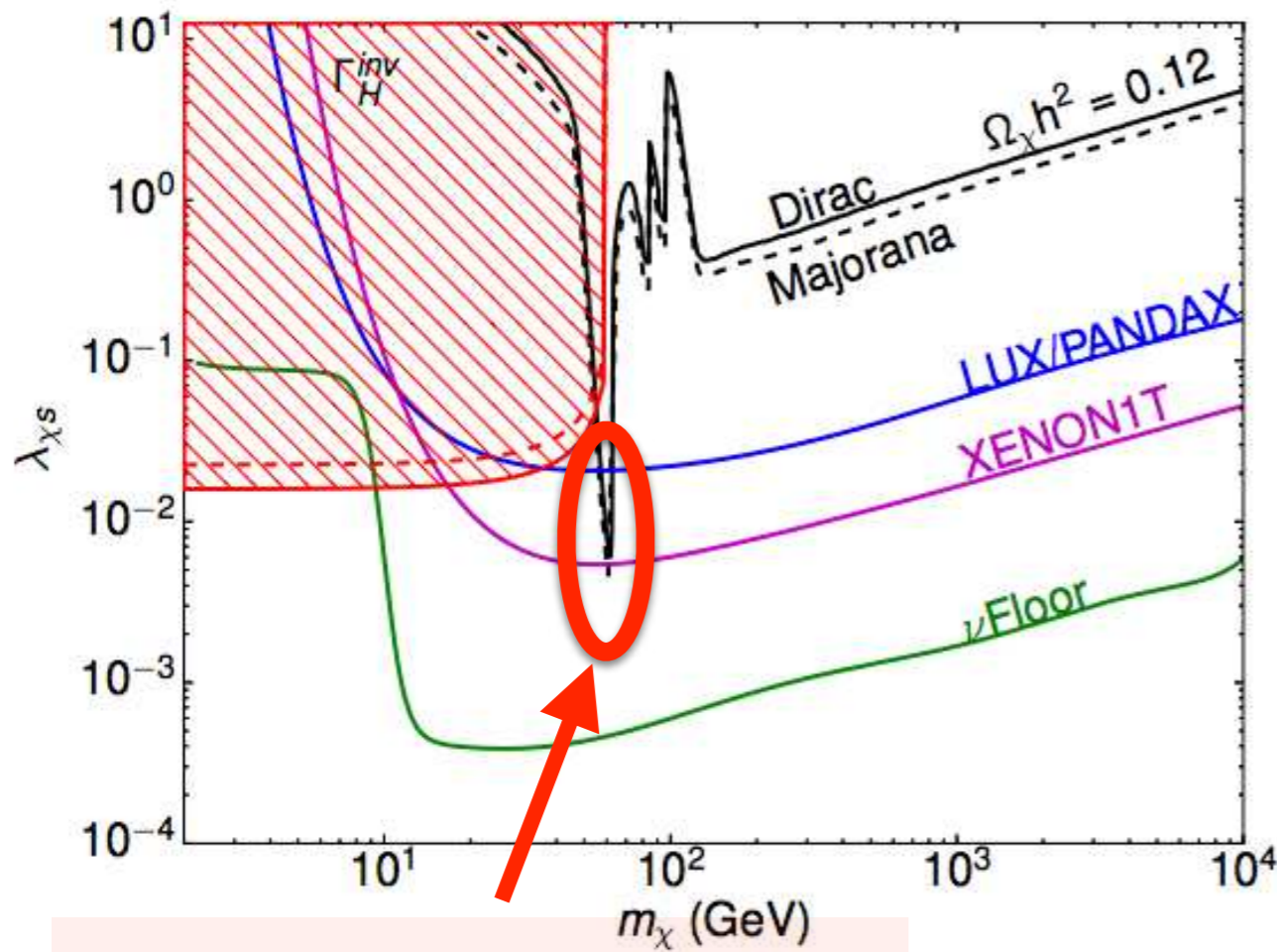
Escudero et al, arXiv:1609.09079





So, after LUX/Pandax there is just one problem — the model is ruled out?

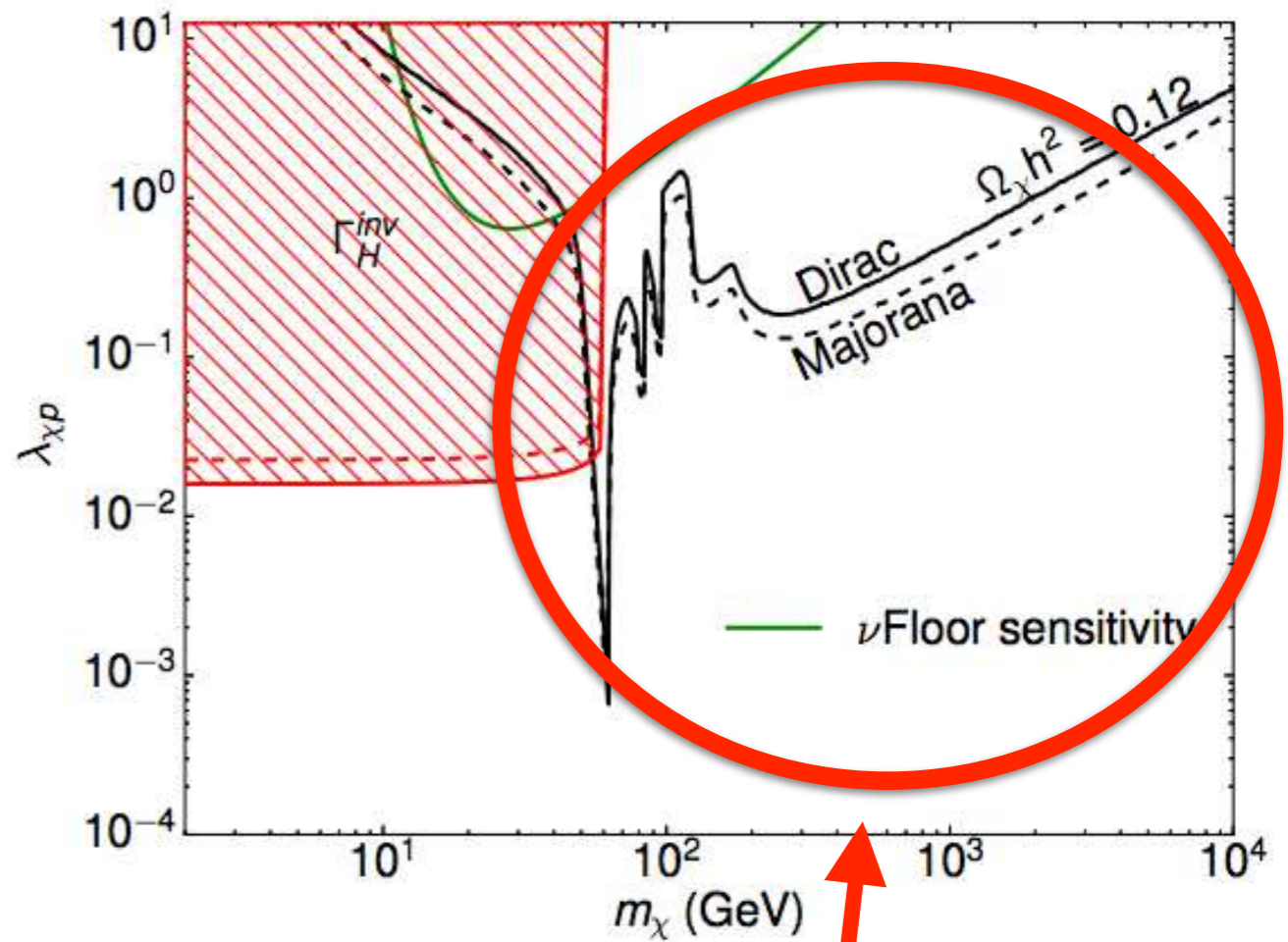
### Scalar



**‘Resonance portal’**

Escudero et al, arXiv:1609.09079

### Pseudoscalar



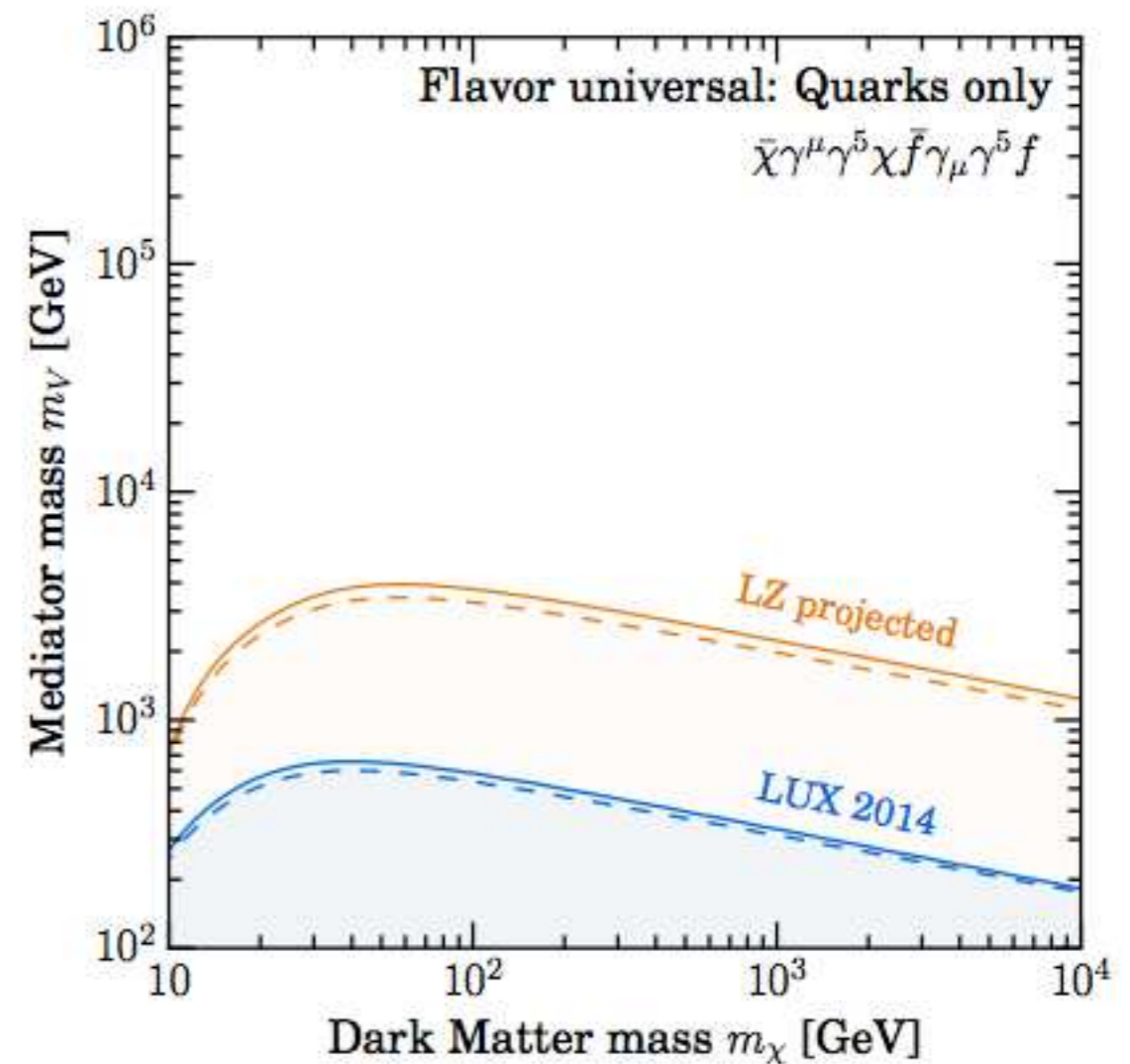
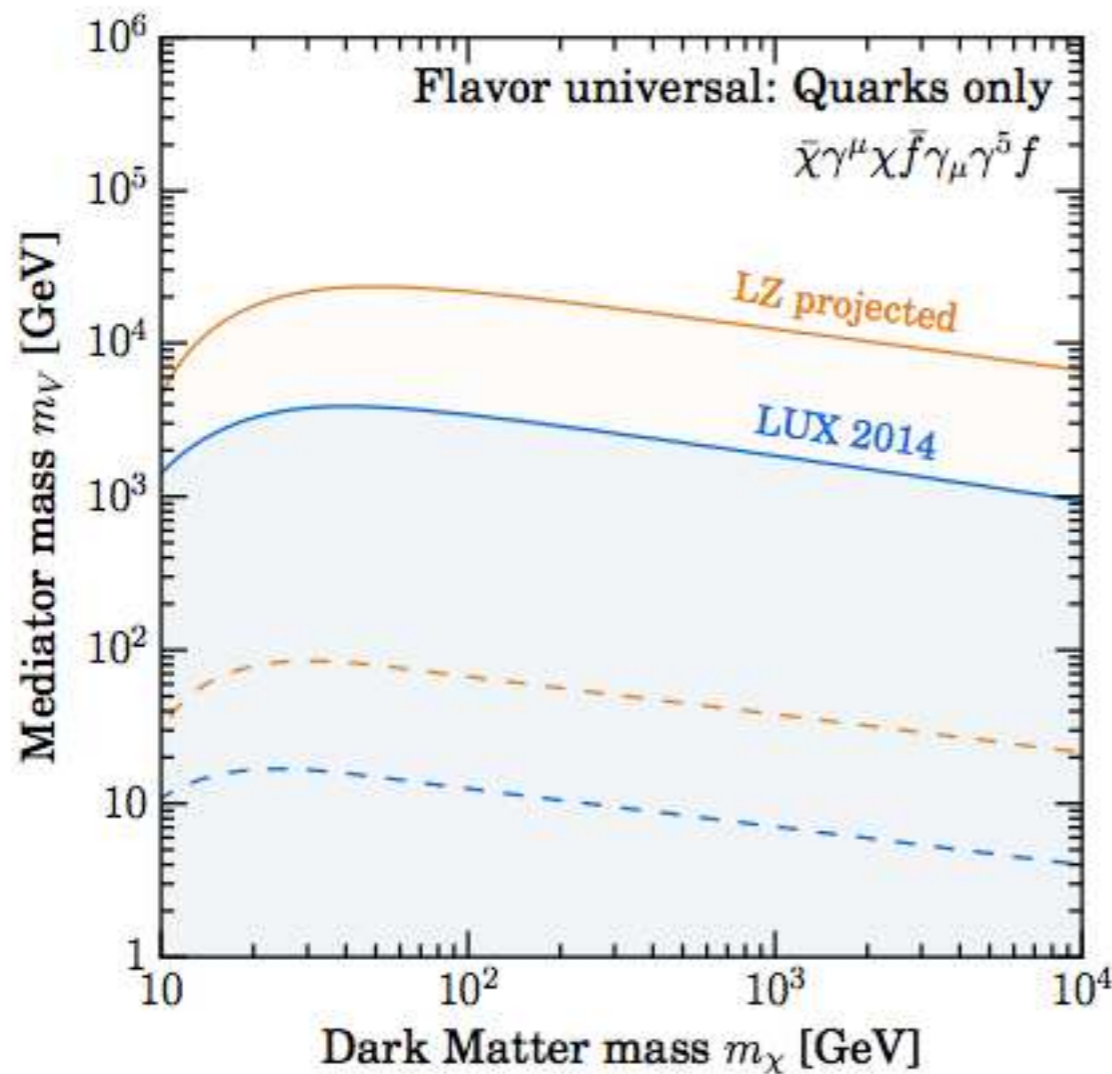
**‘Pseudoscalar portal’**

Escudero et al, arXiv:1609.09079



# A note added

One has to be careful comparing very different energy scales — running!

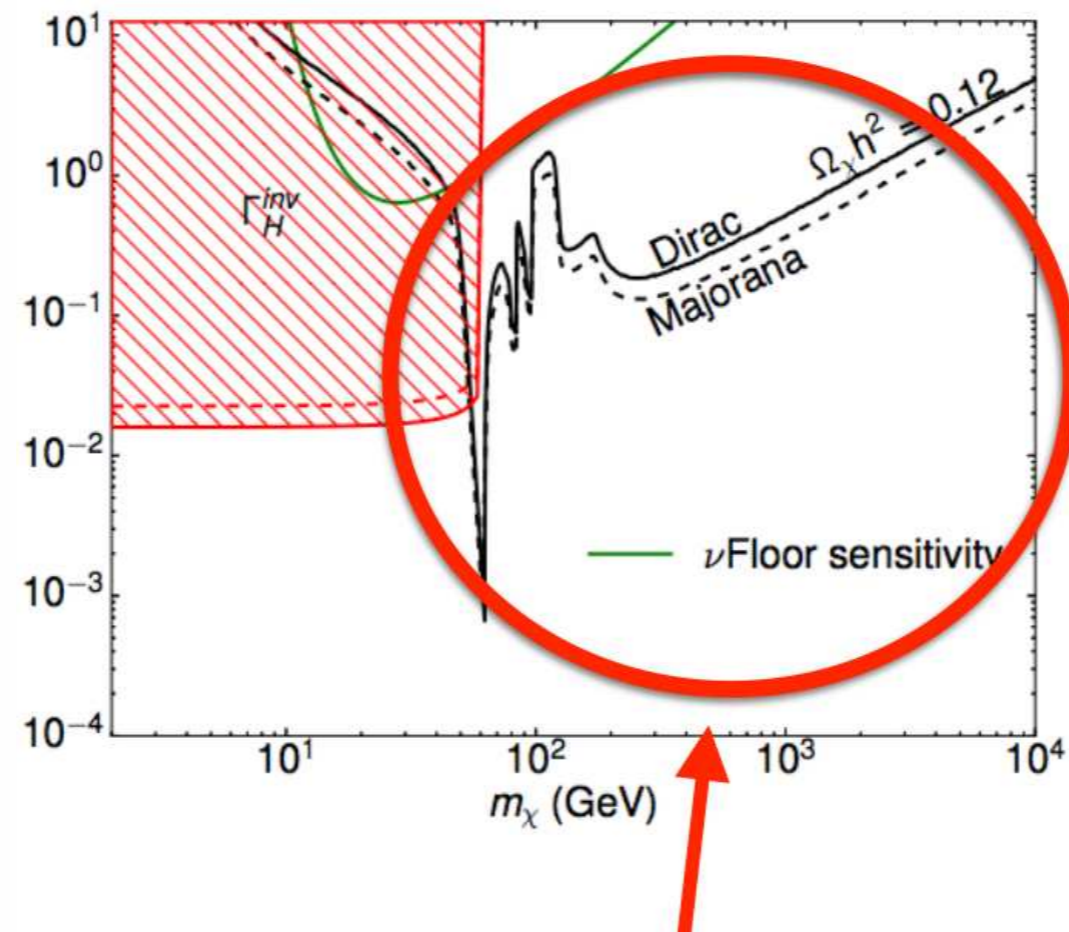


D'Eramo, Kavanagh & Panci, arXiv:1609.09079



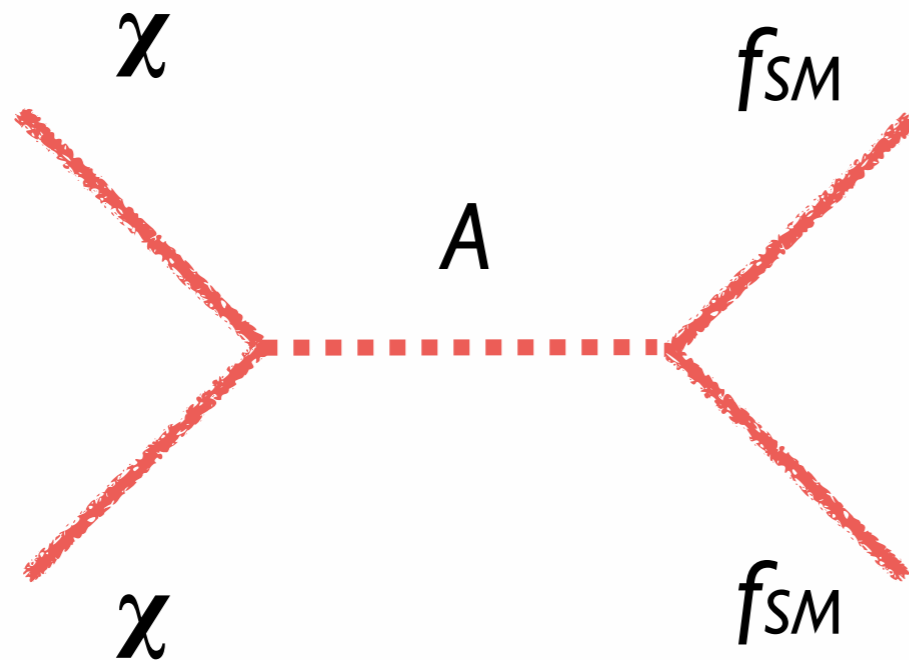
# 'Pseudoscalar portal'

## Pseudoscalar





# Pseudoscalar portal



$$\sigma_{SI} \propto v^2$$

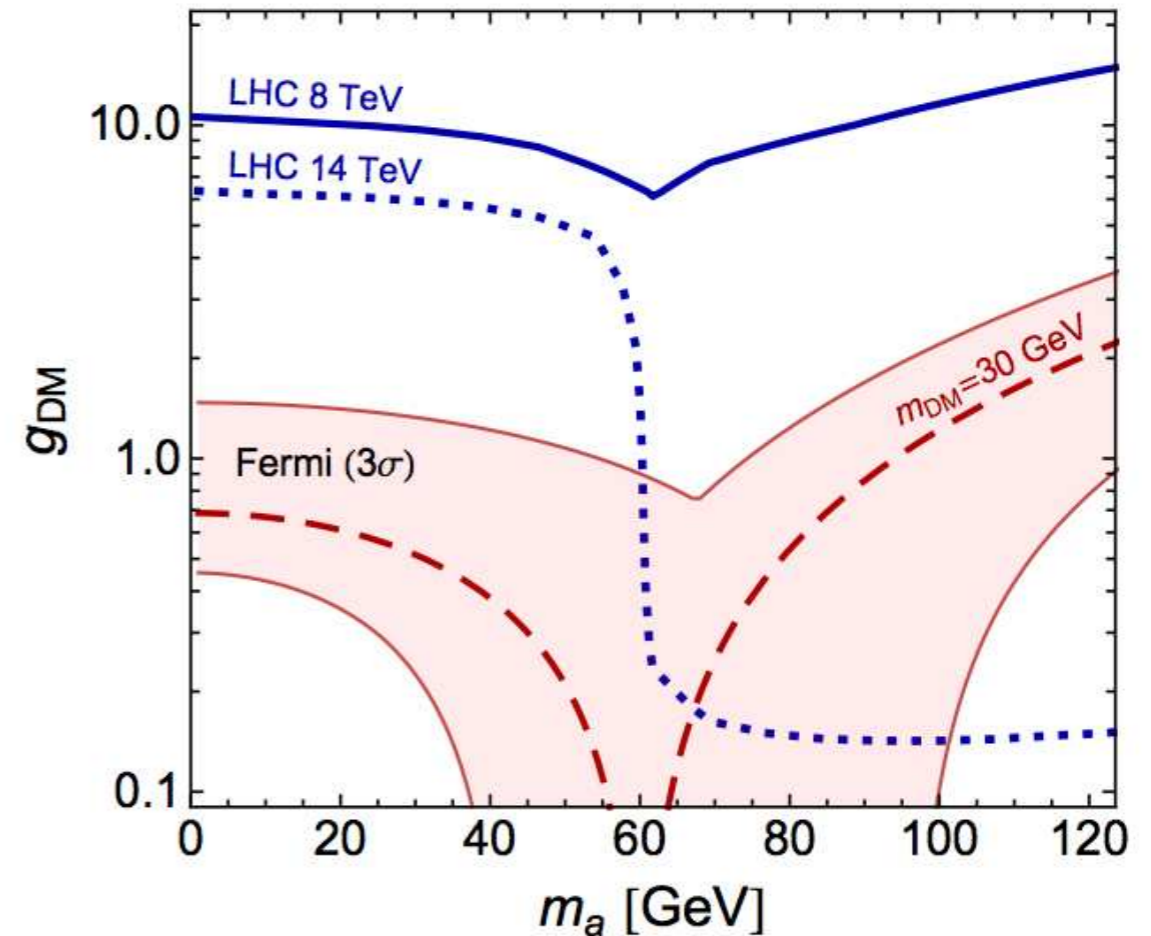
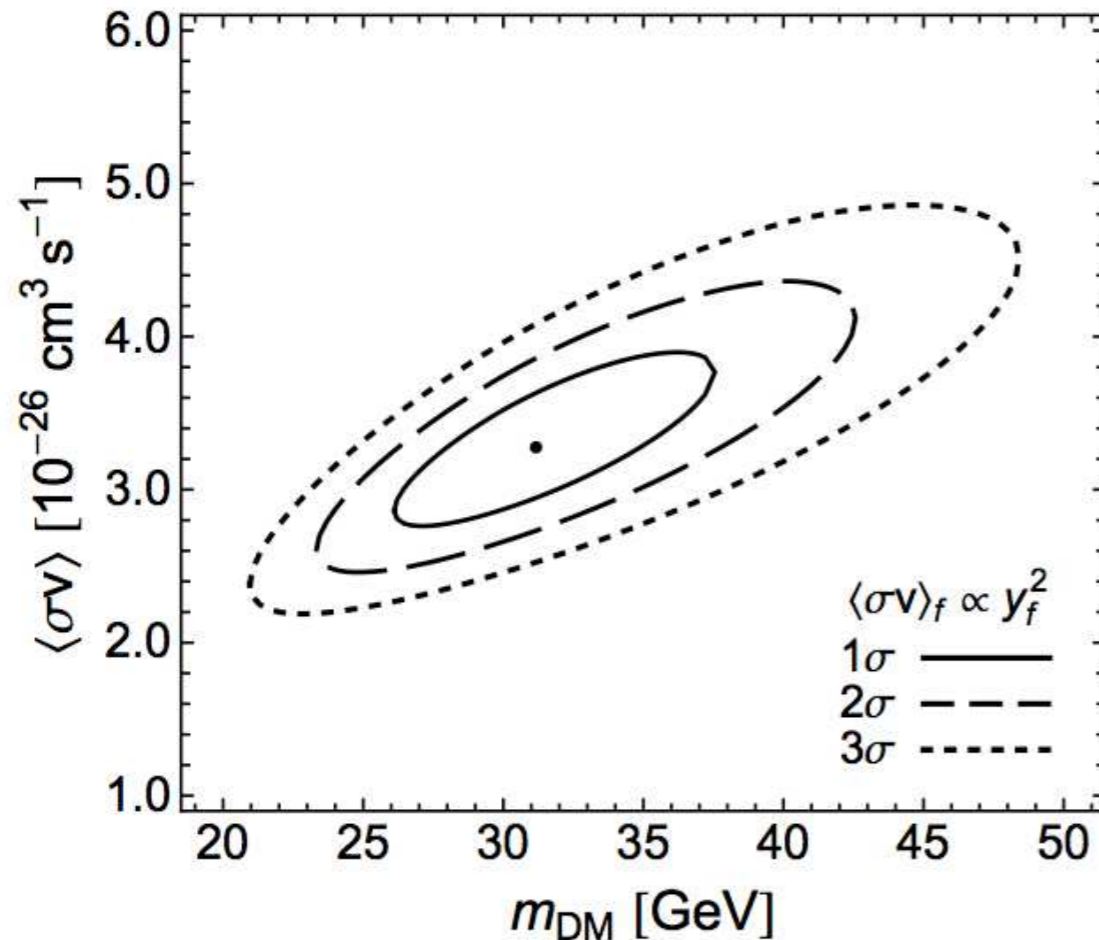
$$\mathcal{L}_{0As\frac{1}{2}} = \frac{1}{2}(\partial_\mu A)^2 - \frac{1}{2}m_A^2 A^2 + \bar{\chi}(i\not{\partial} - m_\chi)\chi - ig_\chi A \bar{\chi}\gamma^5\chi - ig_{SM}A \sum_f \frac{y_f}{\sqrt{2}} \bar{f}\gamma^5 f$$

Galactic Centre Excess — Bohm et al, 1401.6458; Berlin et al, 1404.0022;  
Constraints from flavour physics — Dolan et al, 1412.5174



# Pseudoscalar portal

- Fits the 2-5 GeV gamma-ray excess at GC

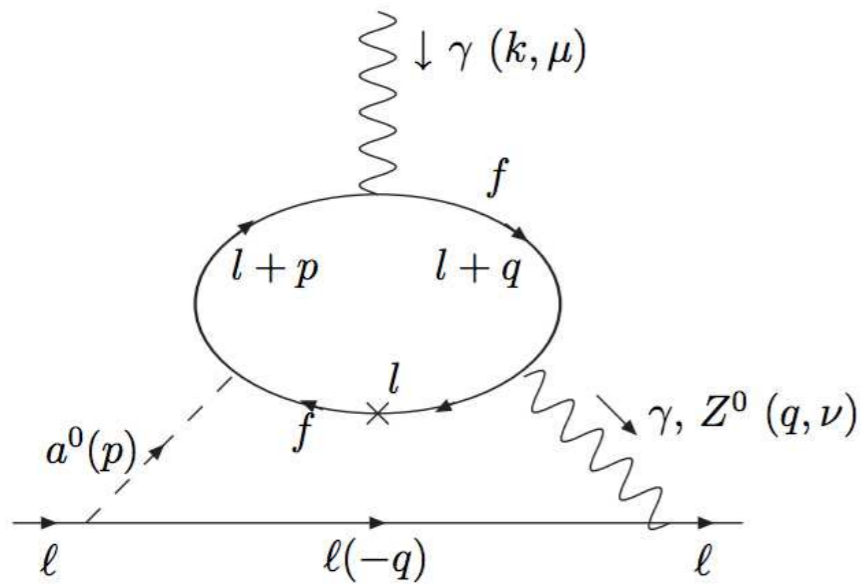


Bohm et al, 1401.6458

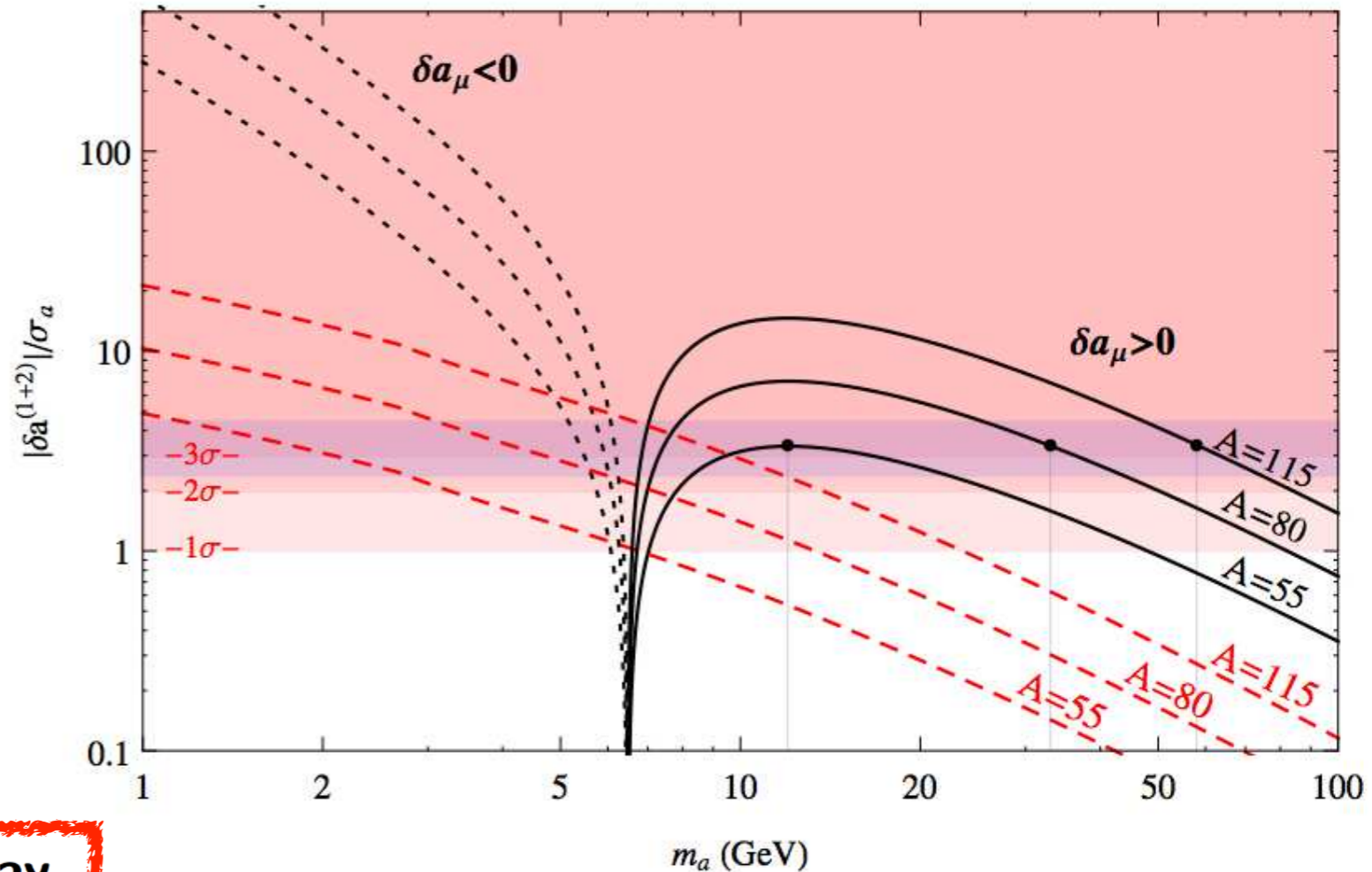


# Pseudoscalar portal & muon $g-2$

'Leptophilic' Coy DM,  
Hektor & Marzola, arXiv:1403.3401



[Chang et al, hep-ph/0009292]



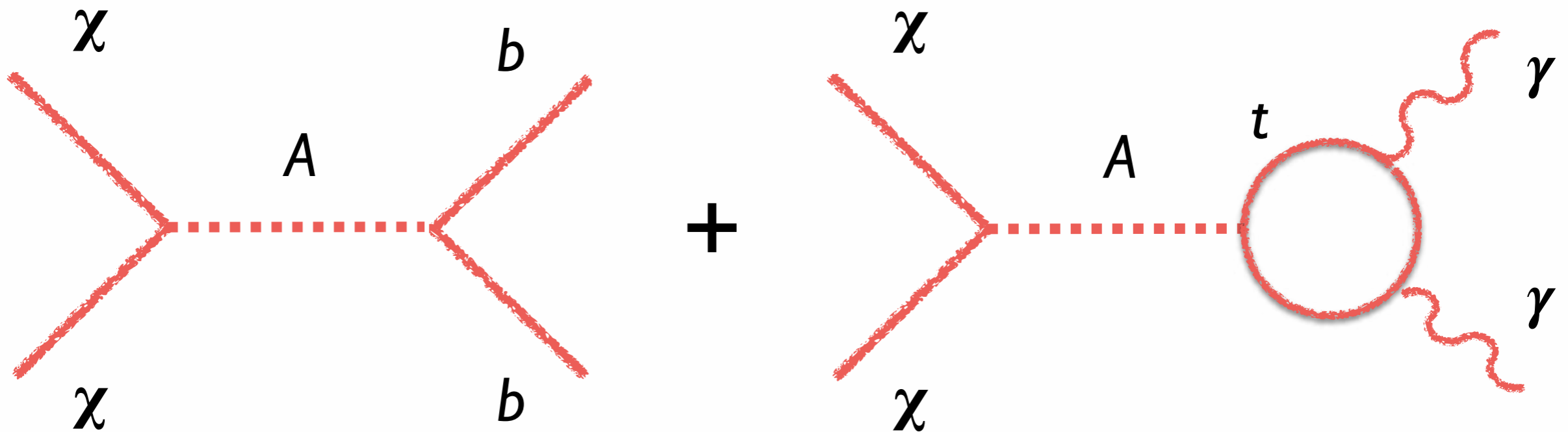
Muon  $g - 2$  and Galactic Centre  $\gamma$ -ray  
excess in a scalar extension of the  
2HDM type-X

arXiv:1507.05096

Andi Hektor,<sup>a,1</sup> Kristjan Kannike<sup>a</sup> and Luca Marzola<sup>a,b</sup>



# Pseudoscalar portal & $\gamma$ -ray line

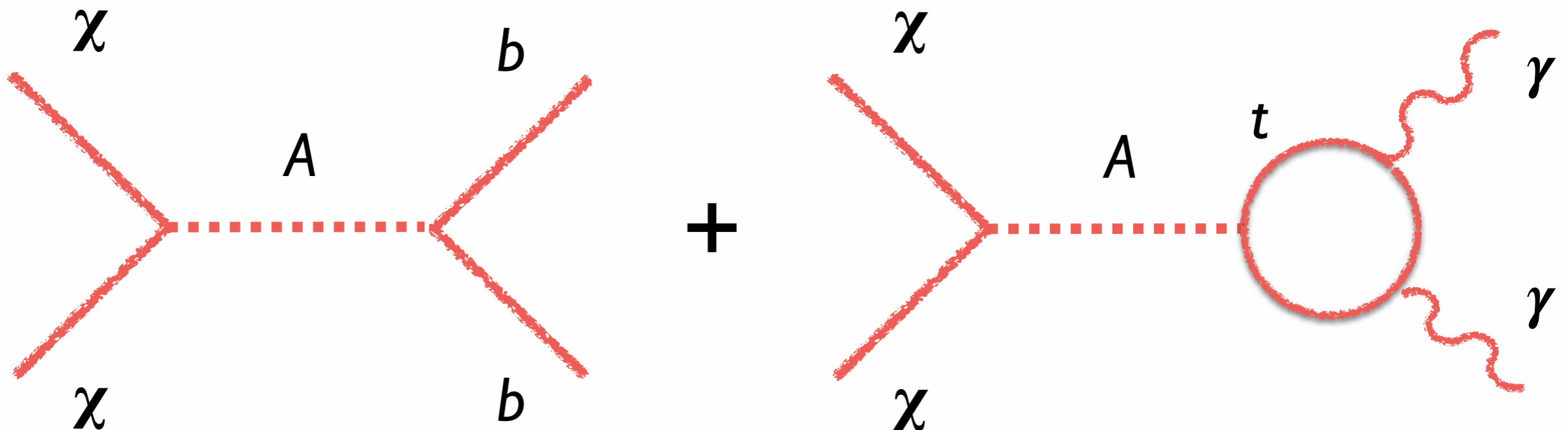


$$m_\chi < m_t$$

$$\mathcal{L}_{0As\frac{1}{2}} = \frac{1}{2}(\partial_\mu A)^2 - \frac{1}{2}m_A^2 A^2 + \bar{\chi}(i\partial - m_\chi)\chi - ig_\chi A \bar{\chi}\gamma^5\chi - ig_{\text{SMA}} A \sum_f \frac{y_f}{\sqrt{2}} \bar{f}\gamma^5 f.$$



# Pseudoscalar portal & $\gamma$ -ray line

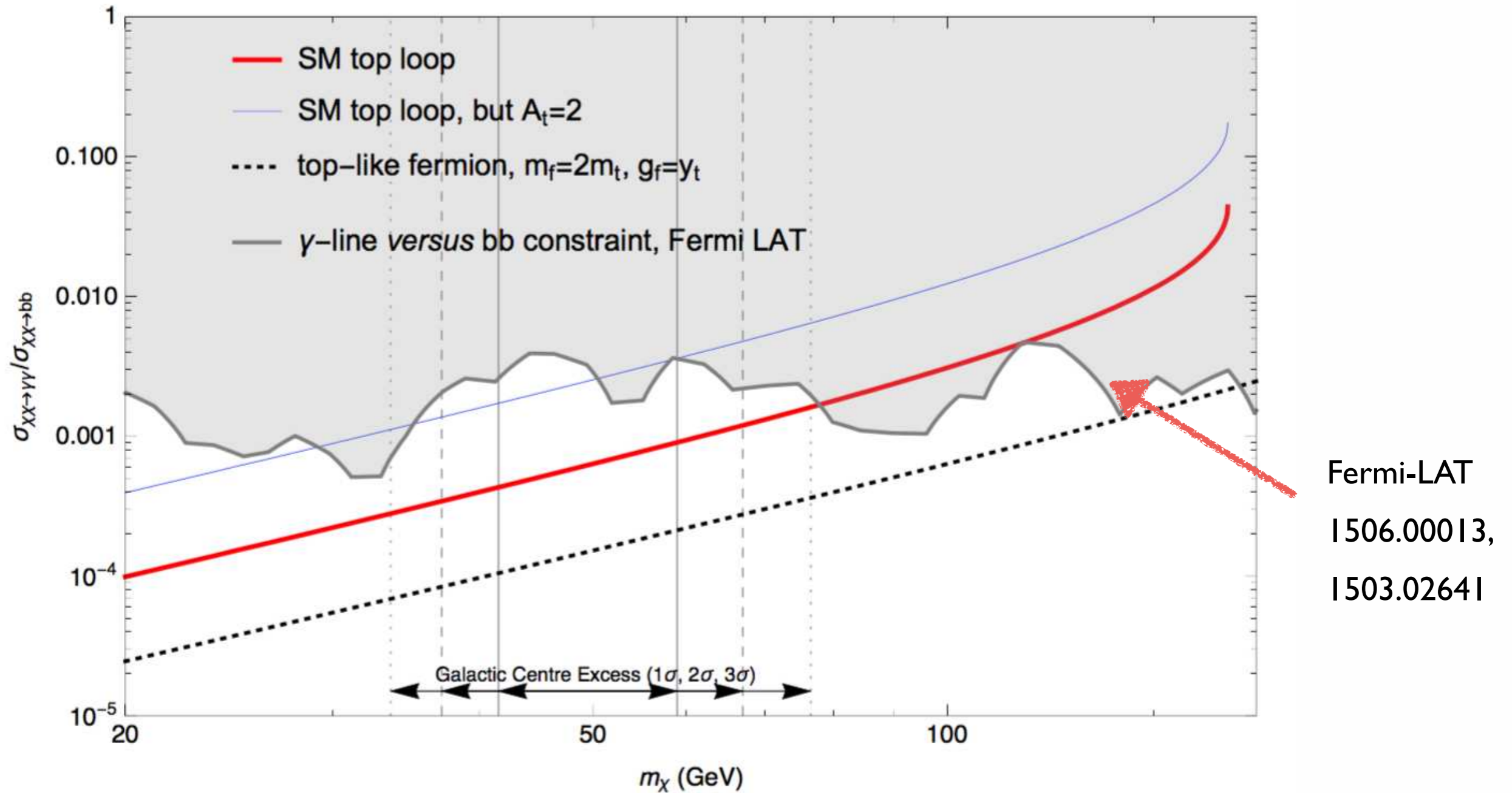


- Loop suppression,  $\sim 10^{-2} \dots 10^{-3}$
- Sensitivity of line/broader distribution search,  $\sim 10^2$





# Ratio of the $\gamma$ -ray line and $bb$ signal

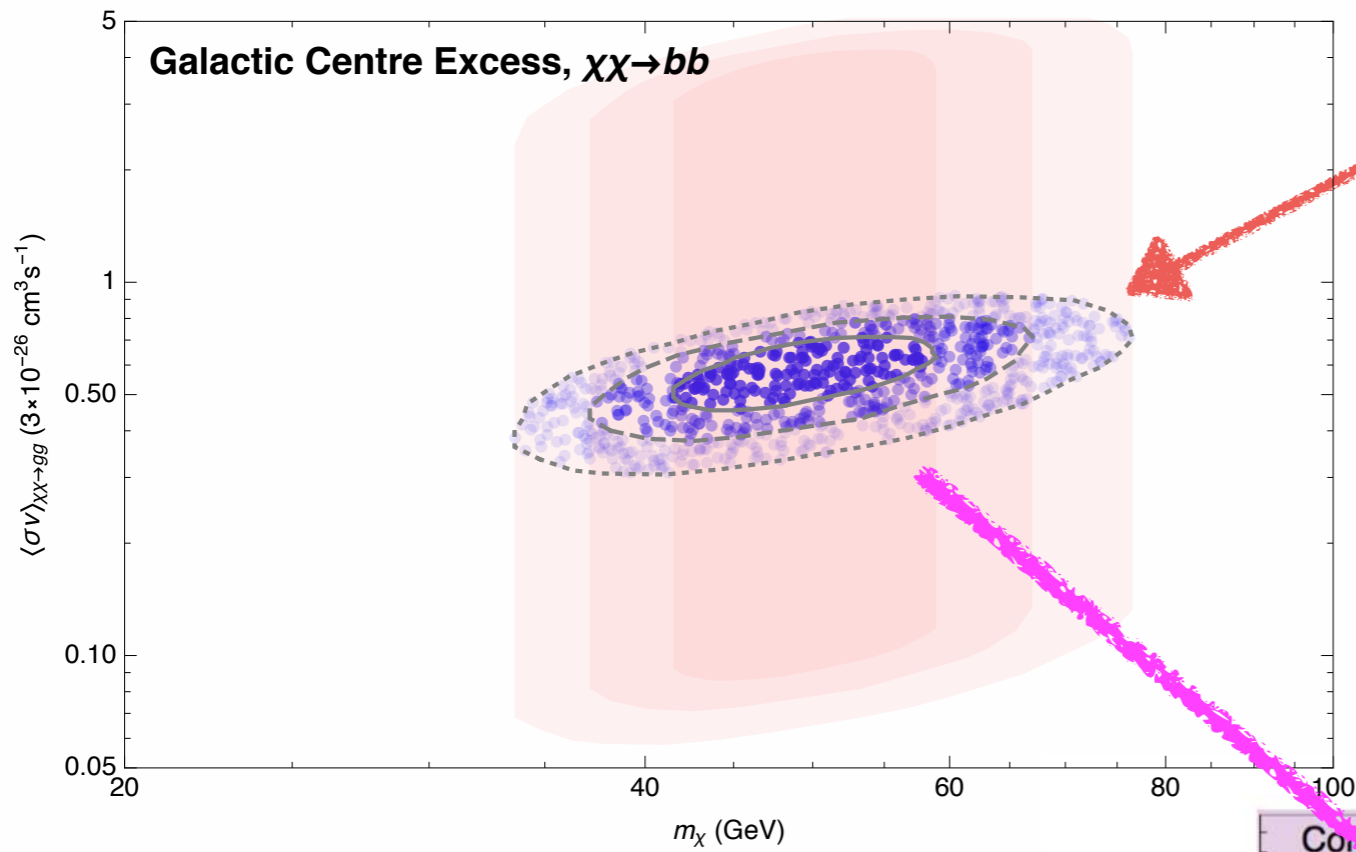


$$R = \frac{(N_c^t \alpha Q_t y_t)^2}{2\pi^2 N_c^b y_b^2} \frac{m_t^2}{m_\chi^2 \sqrt{1 - \frac{m_t^2}{m_\chi^2}}} \left| \arcsin^2\left(\frac{m_\chi}{m_t}\right) \right|^2$$

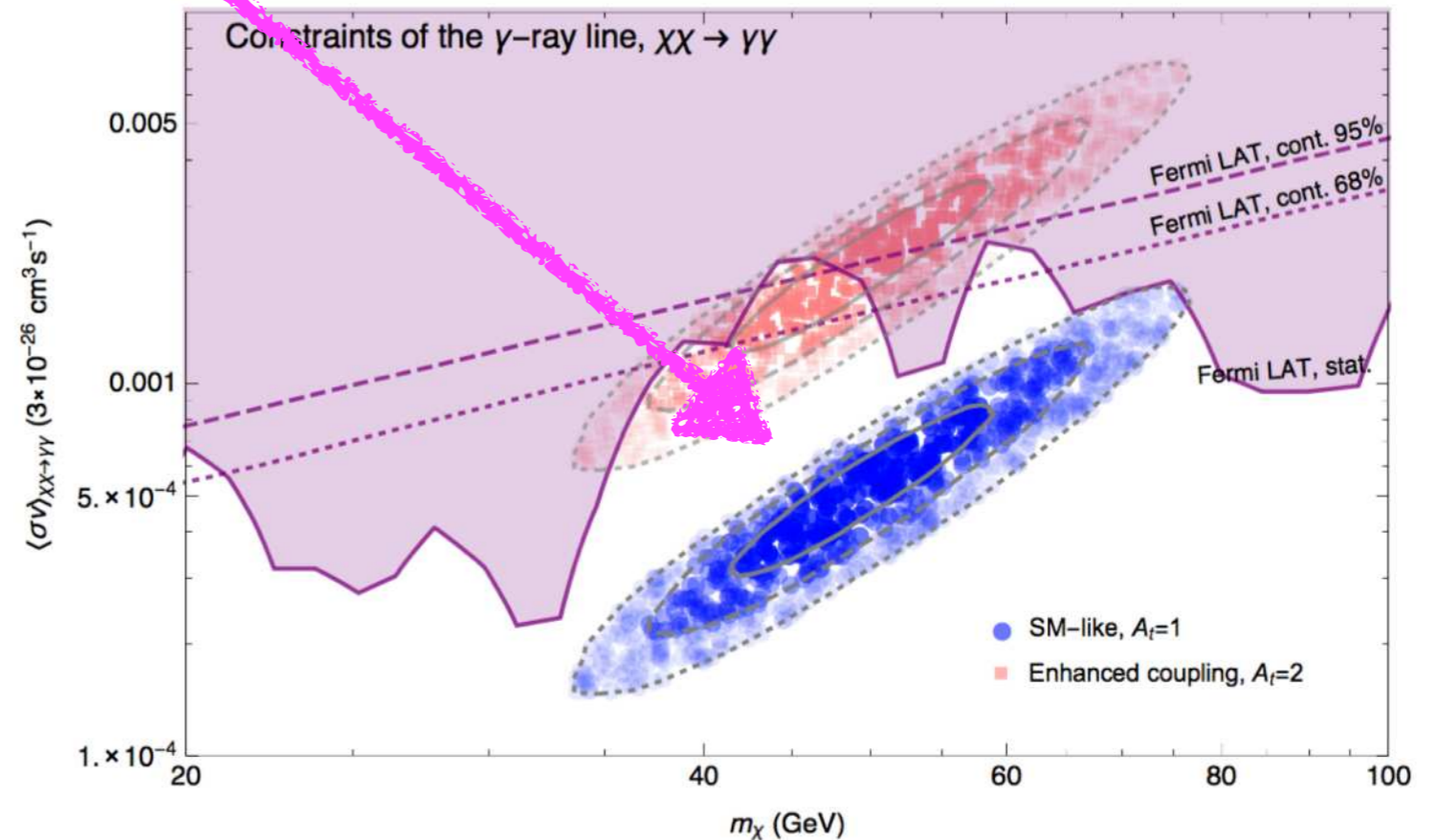


# The $\gamma$ -ray line and the Galactic Centre Excess

Calore et al,  
1411.4647

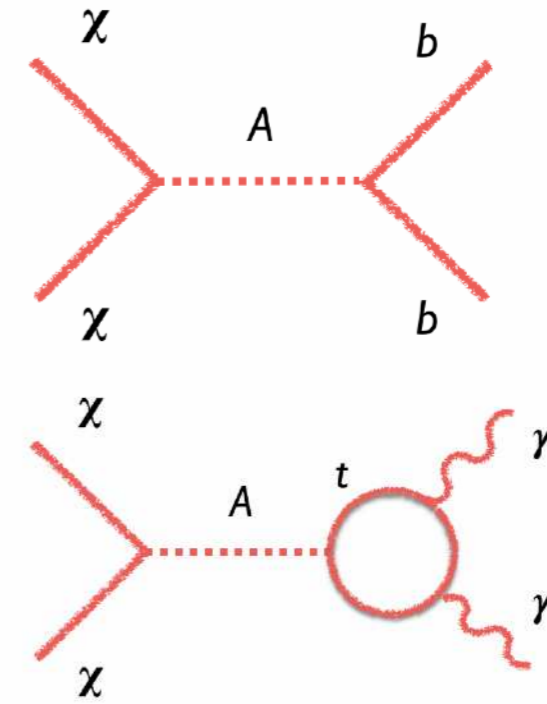
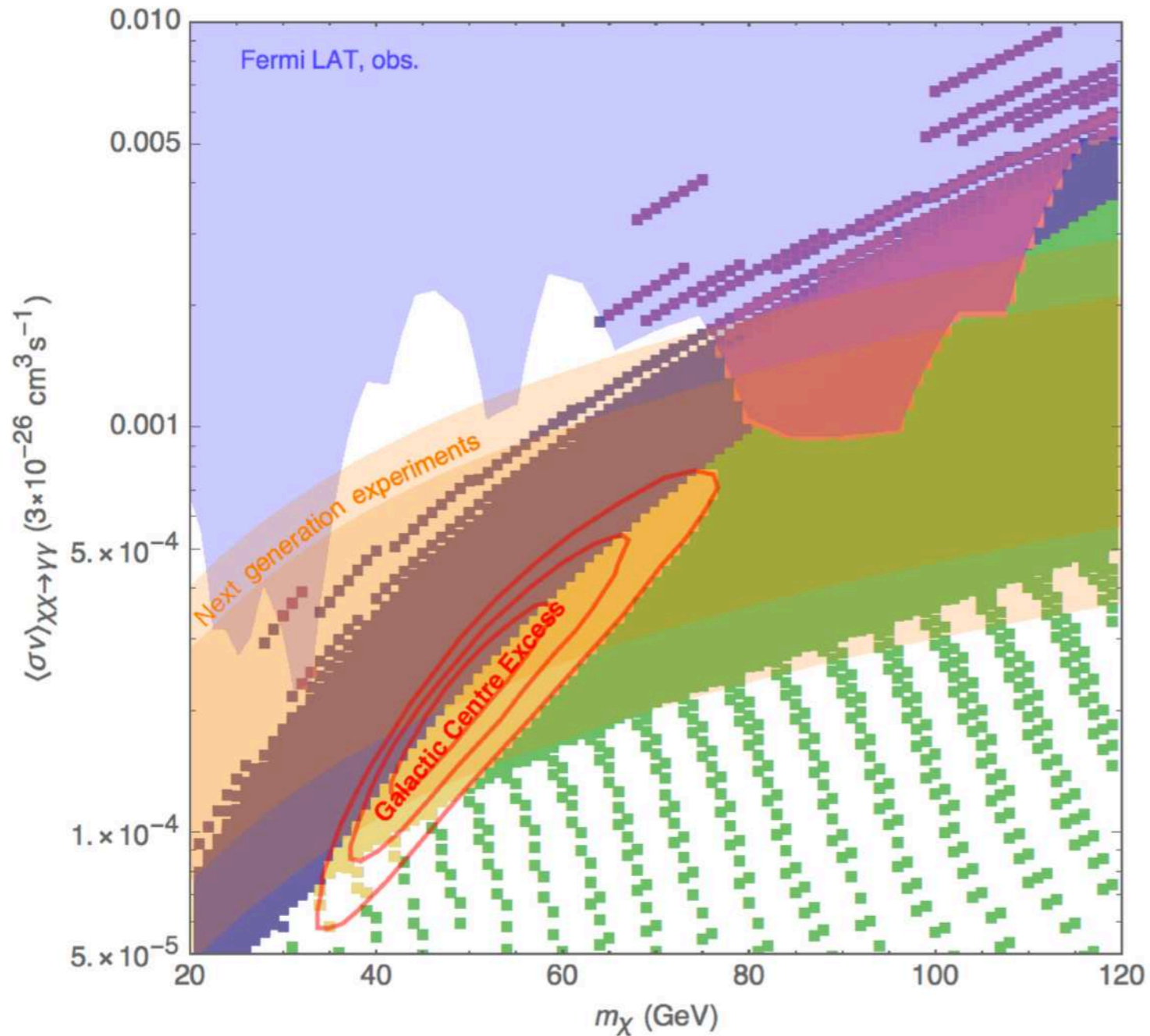


$$R = \frac{(N_c^t \alpha Q_t y_t)^2}{2\pi^2 N_c^b y_b^2} \frac{m_t^2}{m_\chi^2 \sqrt{1 - \frac{m_t^2}{m_\chi^2}}} \left| \arcsin^2\left(\frac{m_\chi}{m_t}\right) \right|^2$$





# The $\gamma$ -ray line and the Galactic Centre Excess



Gamma-ray line constraints on coy dark matter

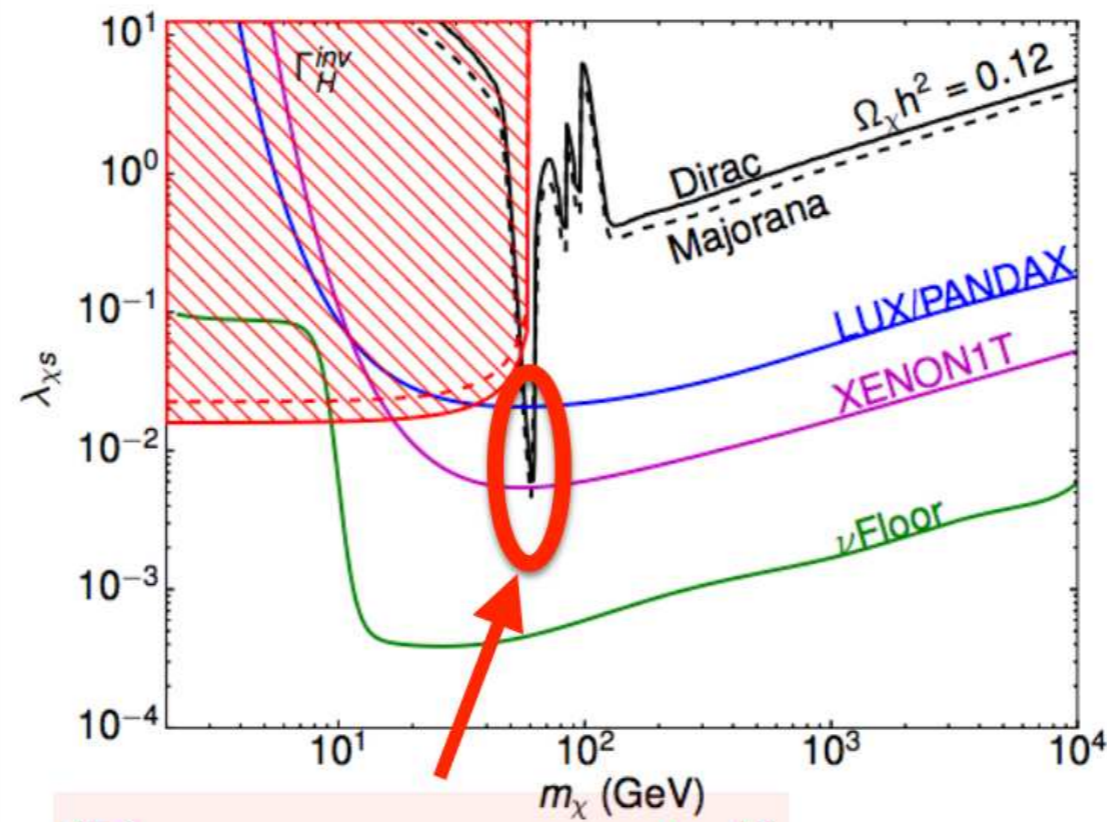
Andi Hektor, Luca Marzola, and Taavi Tuvi

Phys. Rev. D **95**, 121301(R) – Published 16 June 2017



# 'Resonance portal'

Scalar

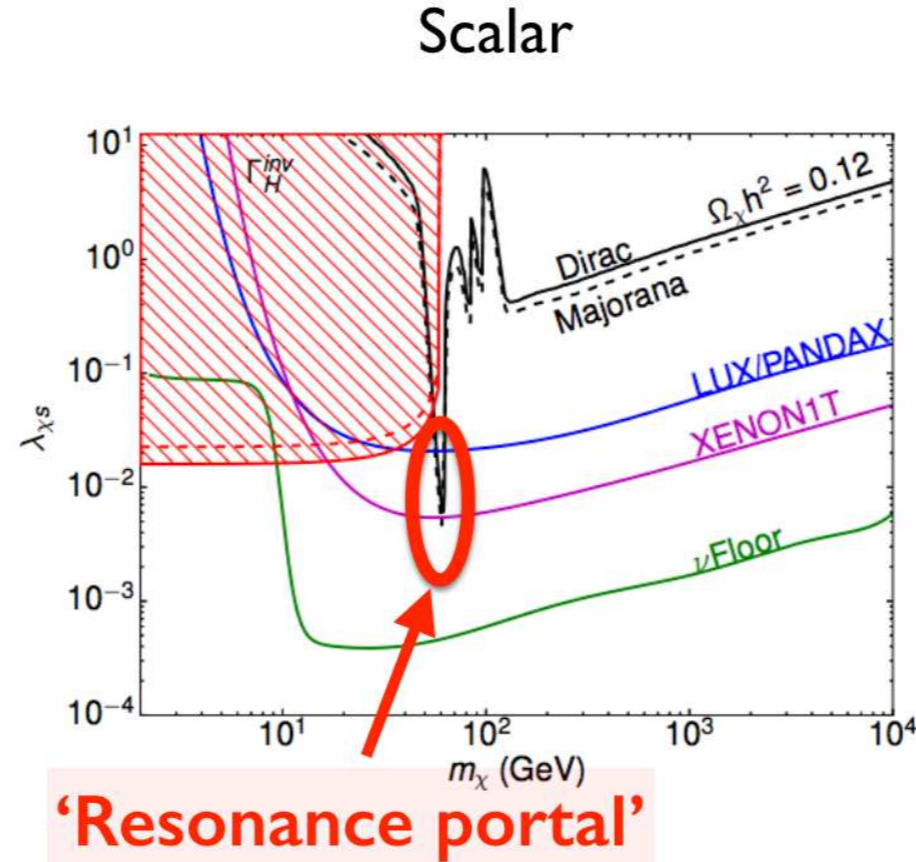


'Resonance portal'



# Is 'Resonance Portal' very tuned?

1. Yes
2. No. There can be theoretical motivation, e.g. hep-ph/9207234, hep-ph/9704403, hep-ph/9804231





# Interesting phenomenology for indirect section!

$$\sigma_{XX \rightarrow S \rightarrow ff} = \frac{N_c (g_\chi A y_f)^2}{16 \pi} \frac{m^2 \left(1 - \frac{m_f^2}{m^2}\right)^{3/2}}{(M^2 - s[m^2, v_{rel}])^2 + M^2 \Gamma^2} v_{rel}$$

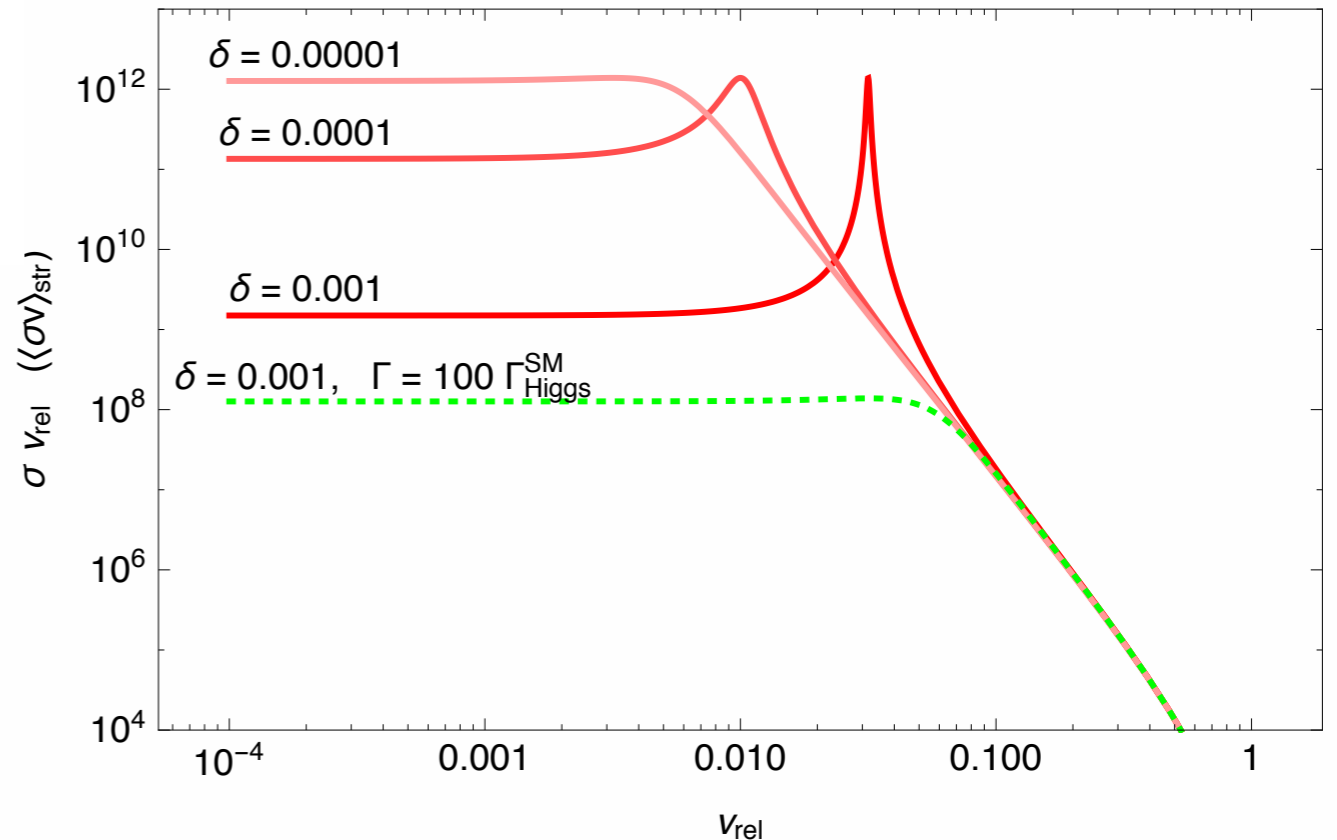
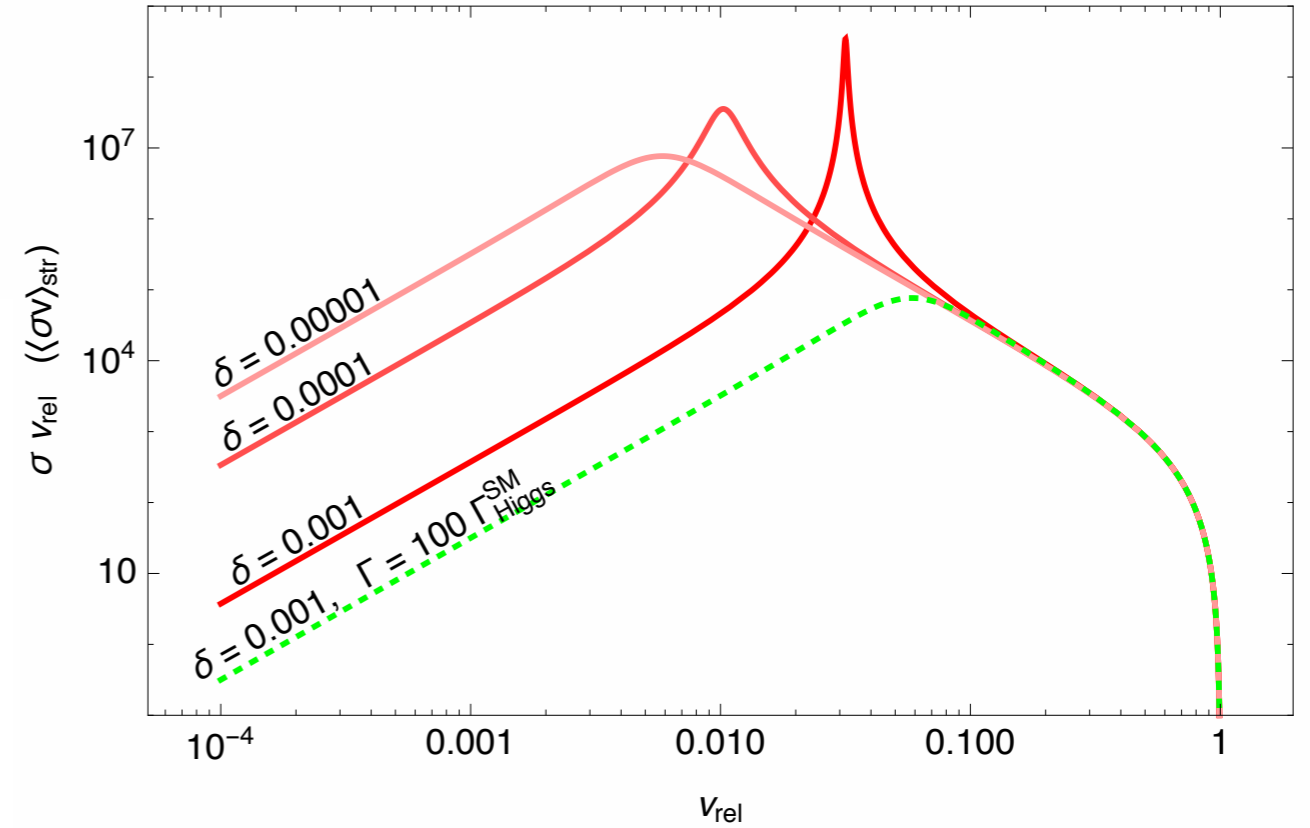
$$\sigma_{XX \rightarrow A \rightarrow ff} = \frac{N_c (g_\chi A y_f)^2}{16 \pi} \frac{m^2 \left(1 - \frac{m_f^2}{m^2}\right)^{1/2}}{(M^2 - s[m^2, v_{rel}])^2 + M^2 \Gamma^2} v_{rel}^{-1}$$

$$m^2 = \frac{M^2}{4(1+\delta)}, \quad s[m^2, v_{rel}] = \frac{4m^2}{1-v_{rel}^2}$$

$$M = m_{higgs}, \quad g_\chi = 1$$

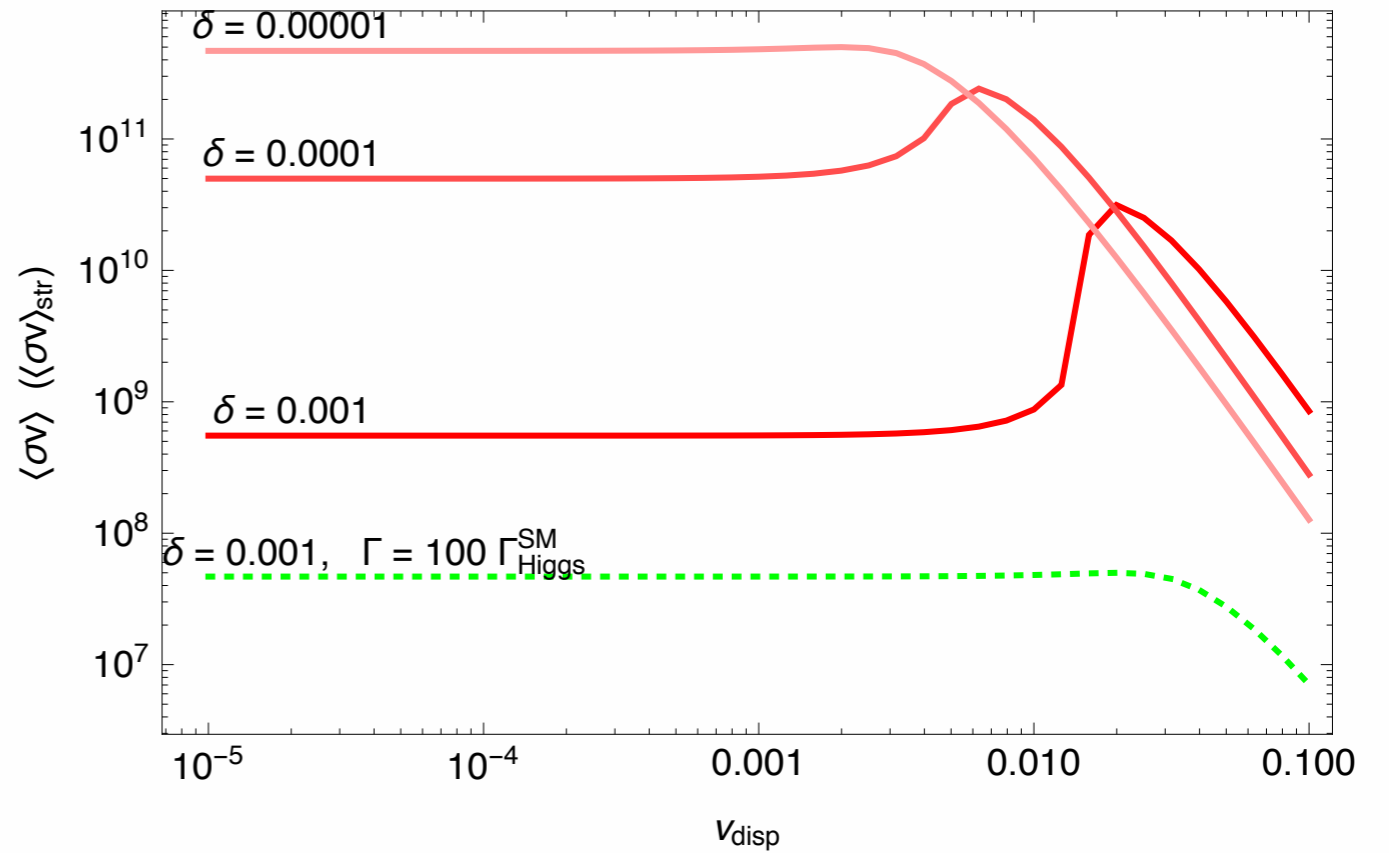
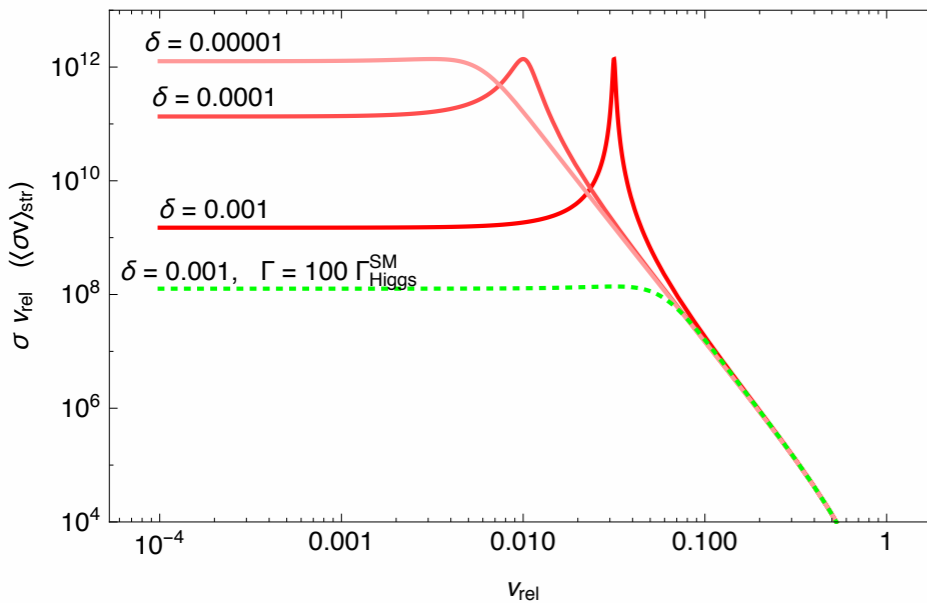
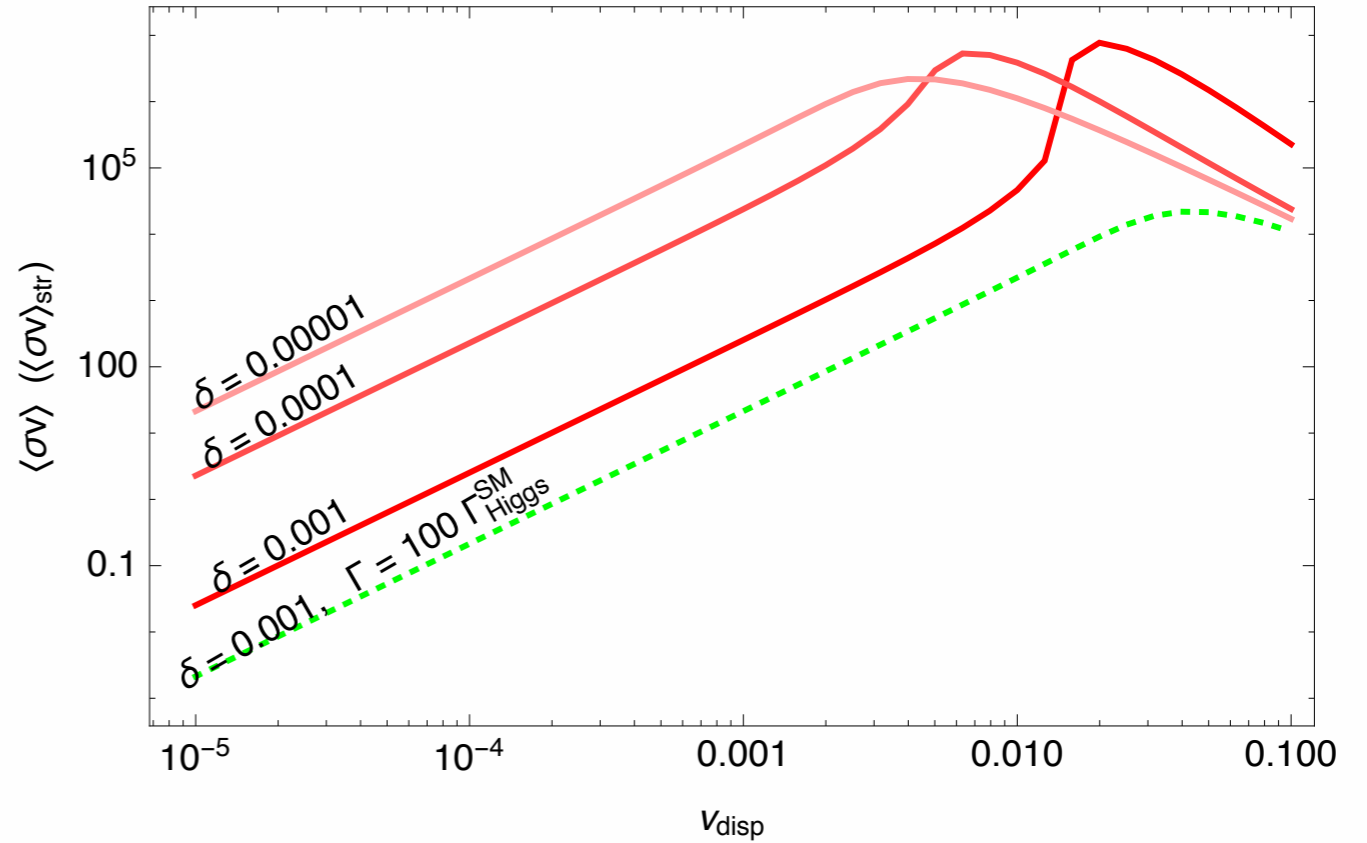
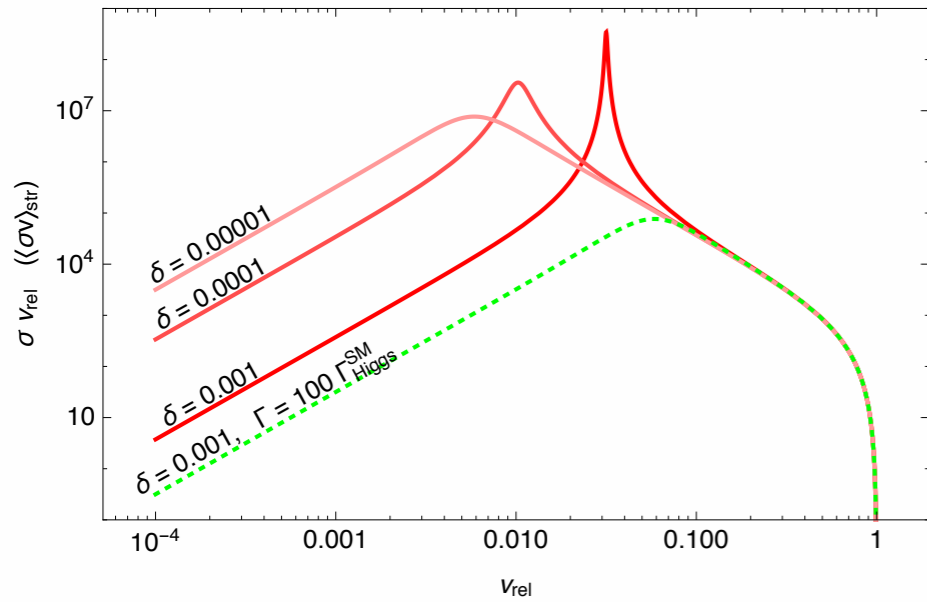
$$A = 1, \quad y_f = 1$$

$$\Gamma = \Gamma_{higgs}$$





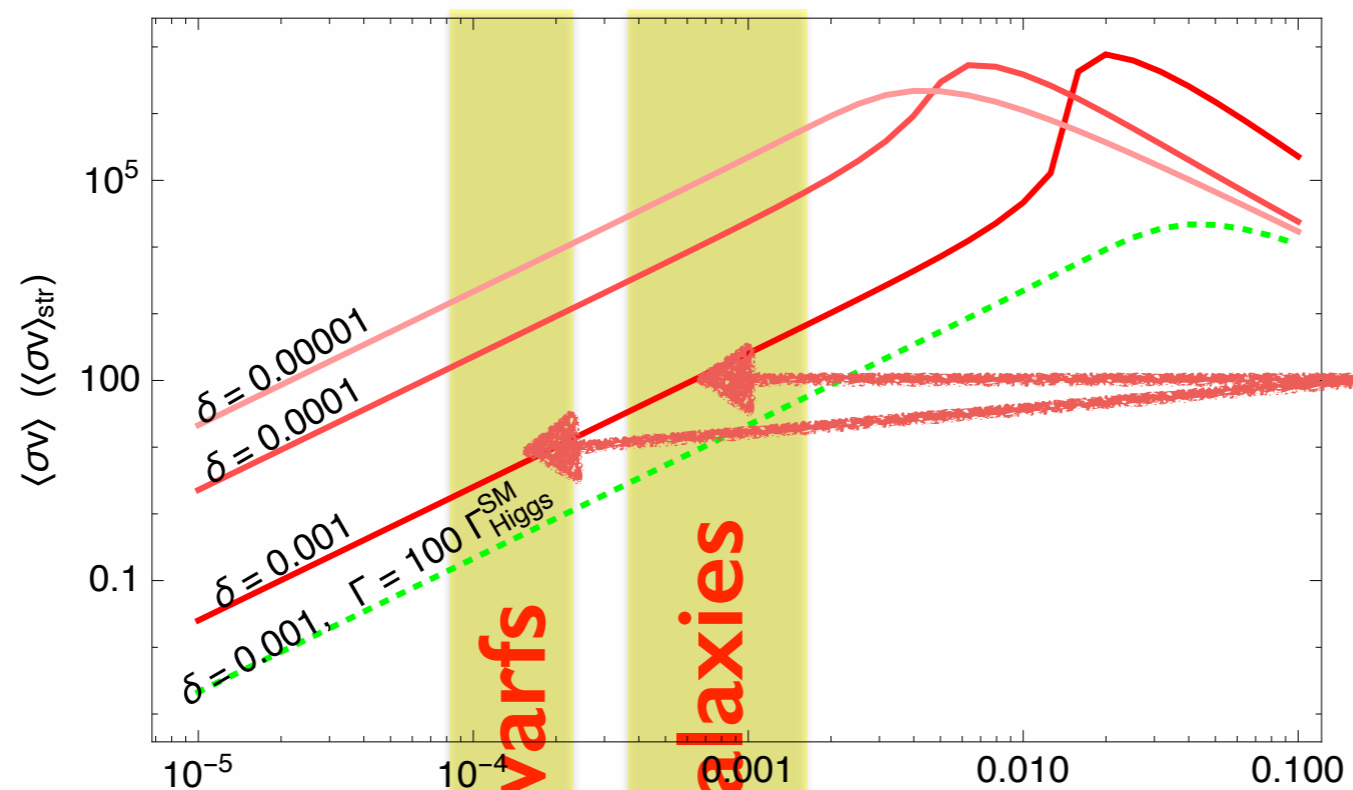
# Interesting phenomenology for indirect section!





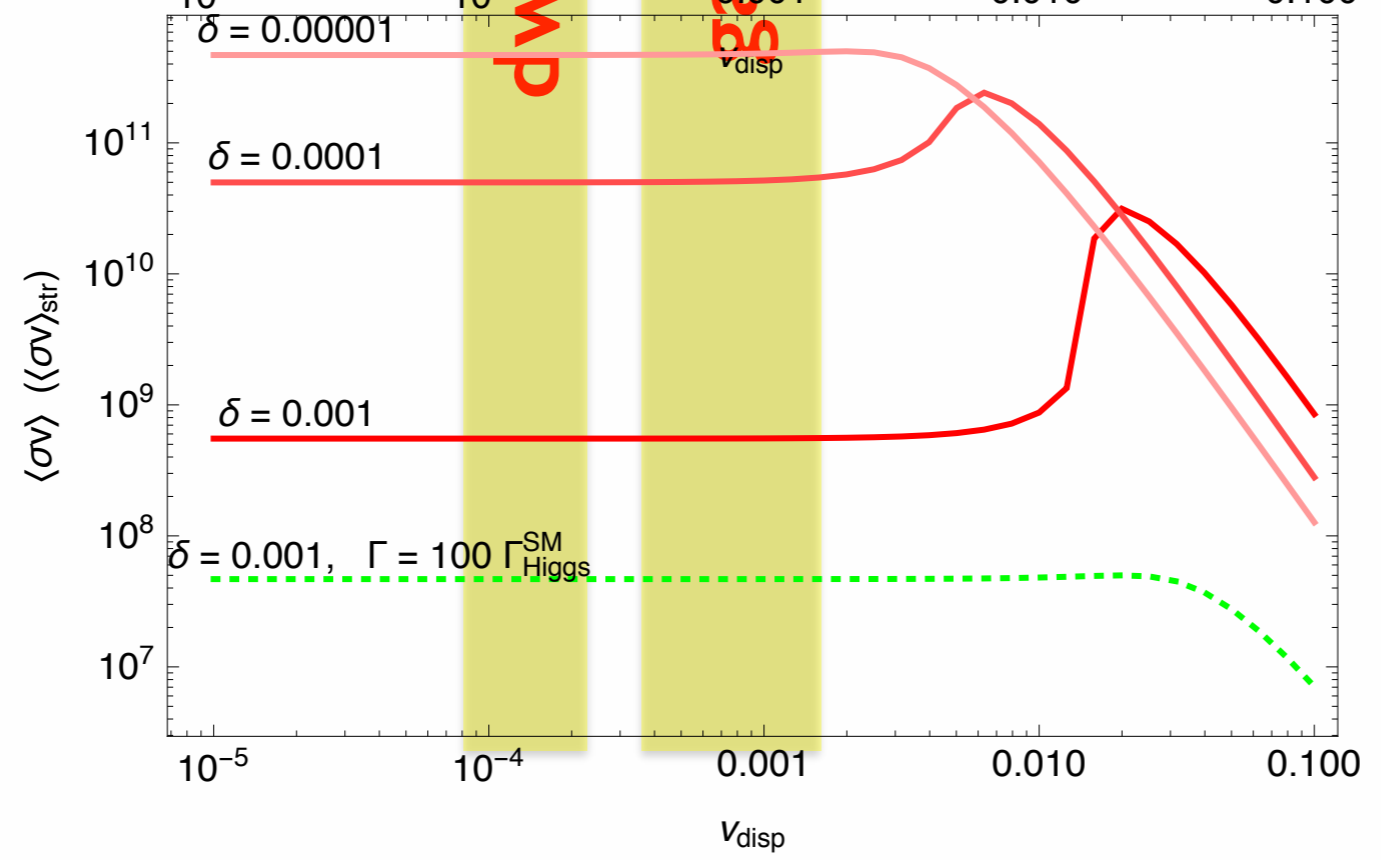
# Interesting phenomenology for indirect section!

scalar



**Difference of two orders magnitude!**

pseudo-scalar



**Real life is more complicated!**





‘Thermal portal’



# Thermal effects on the DM production

Recent interest in thermal effects on dark matter:

'forbidden' channels, D'Agnolo & Ruderman 1505.07107

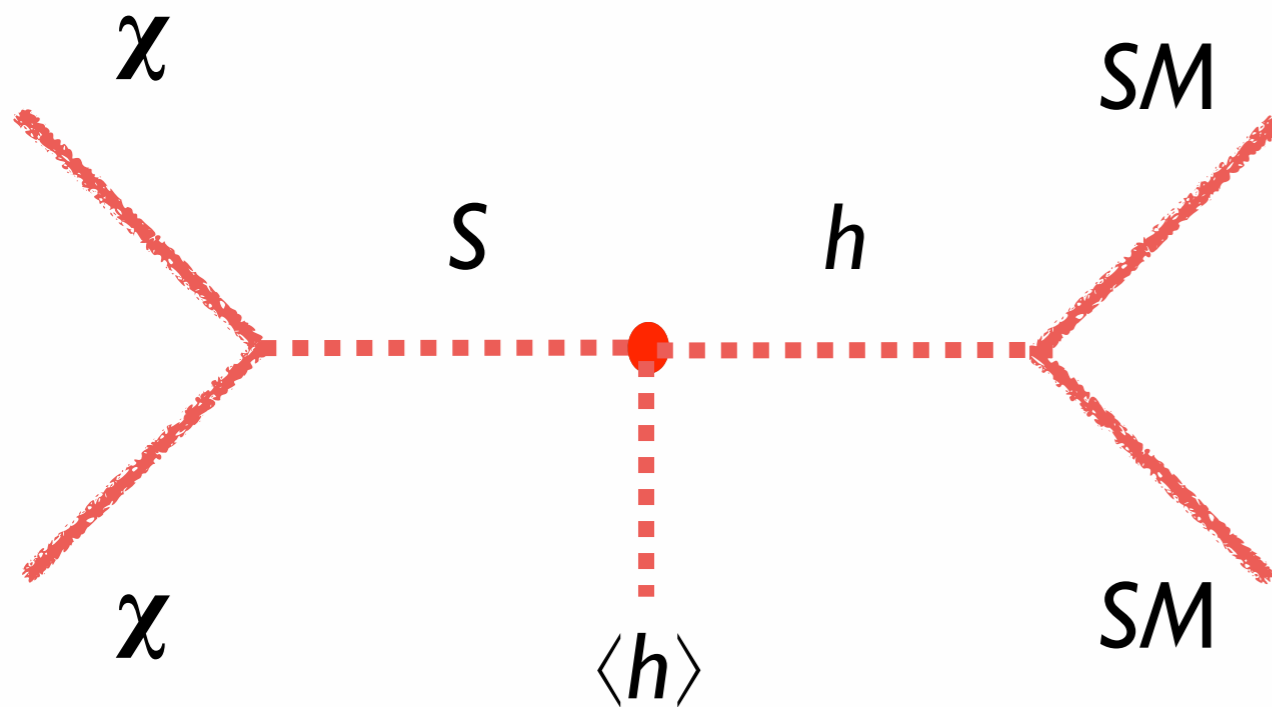
cannibal DM, Pappadopulo, Ruderman & Trevirsan 1602.04219

VEV flip-flop Baker & Kopp 1608.07578

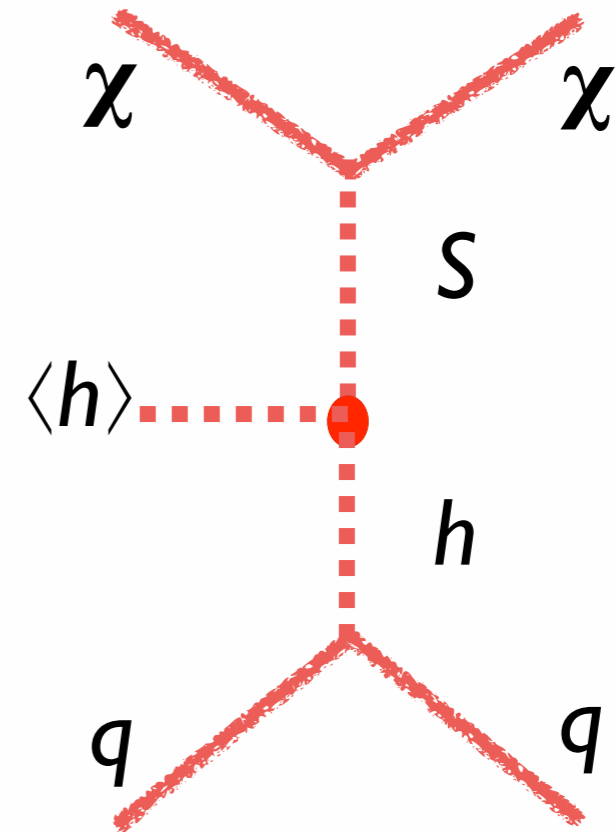


# Thermal effects of the portal?

Higgs-scalar mixing,  
the freeze-out diagram



Direct detection



Typical assumptions at the freeze-out:

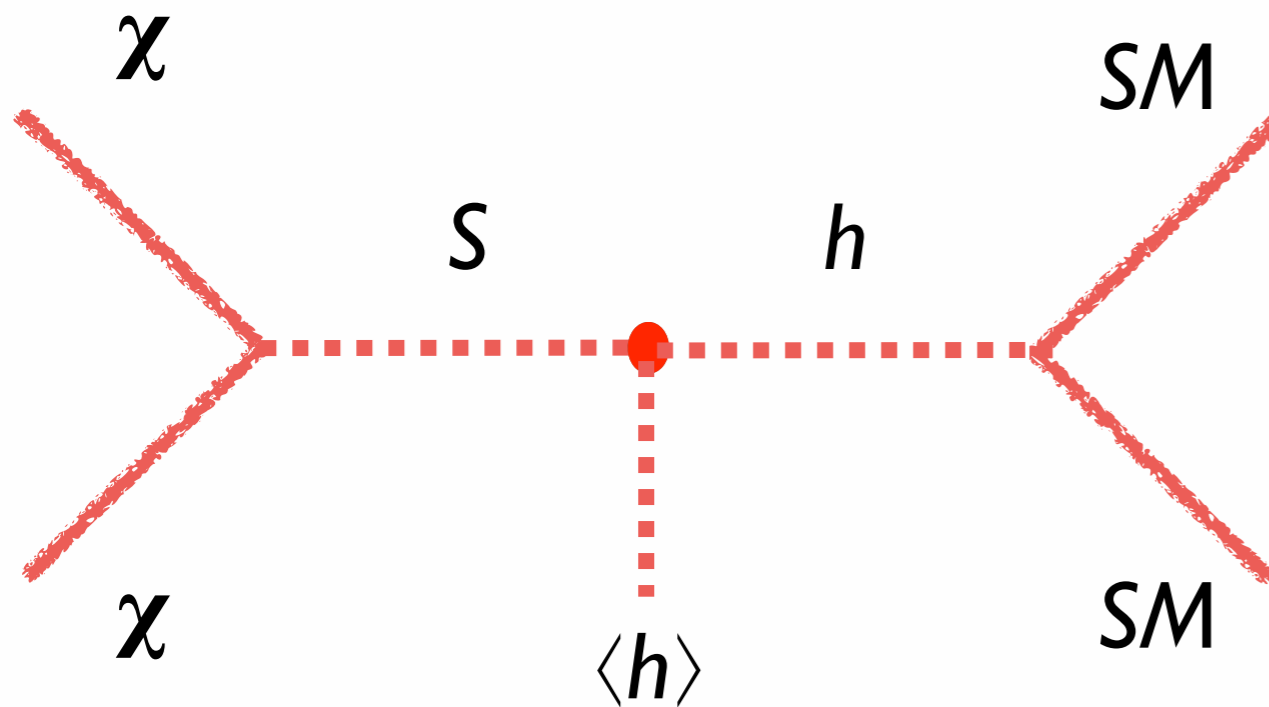
$$\langle h \rangle \simeq \text{const}(\text{vacuum}) > 0$$

$$\langle S \rangle \simeq \text{const}(\text{vacuum}) > 0 \text{ or } = 0$$

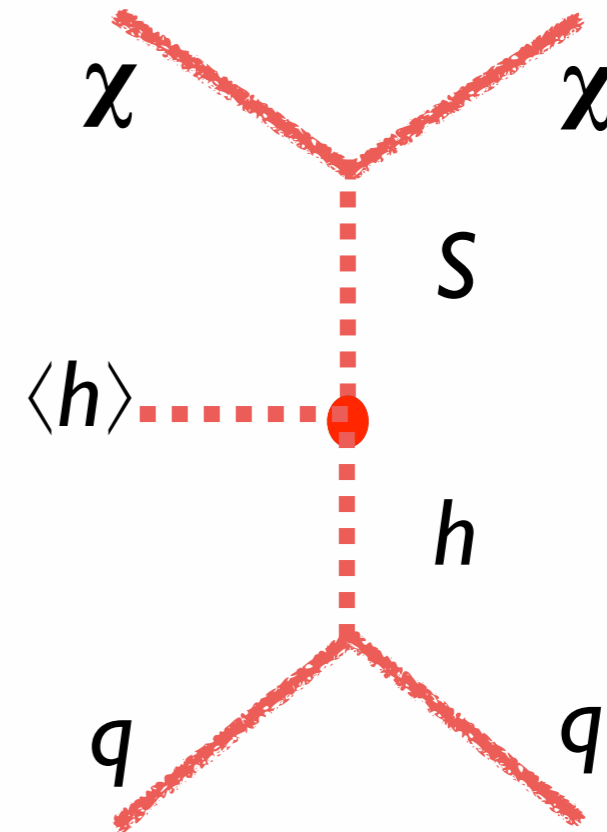


# Thermal effects of the portal?

Higgs-scalar mixing,  
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Typical assumptions at the freeze-out:

$$\langle h \rangle \simeq \text{const}(\text{vacuum}) > 0$$

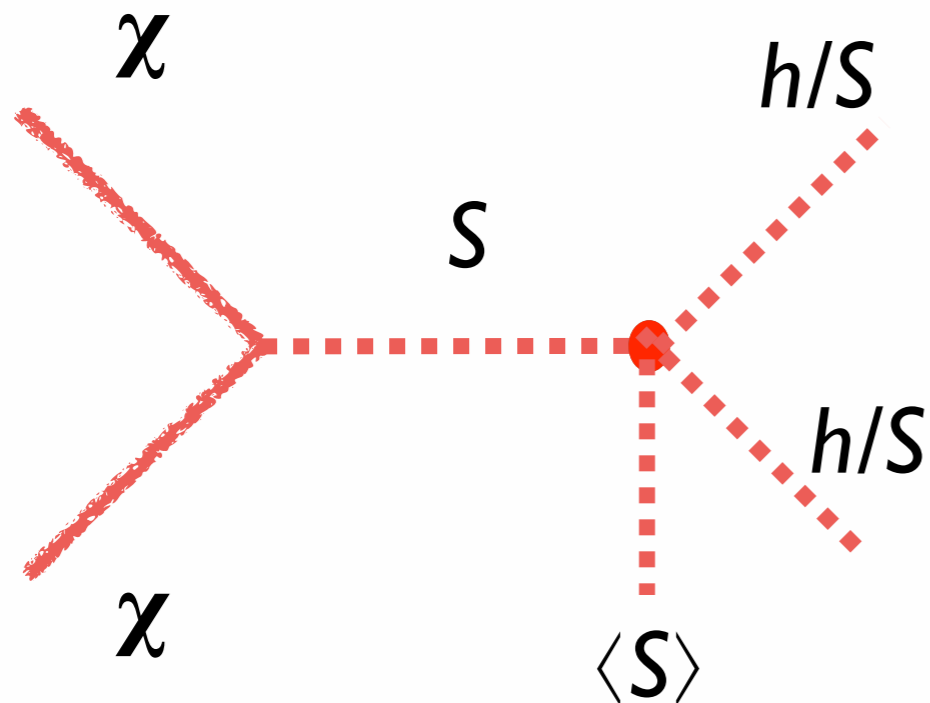
$$\langle S \rangle \simeq \text{const}(\text{vacuum}) > 0$$

*But  $\langle h \rangle$  and  $\langle S \rangle$  are not constants,  
they are the functions of temperature!*

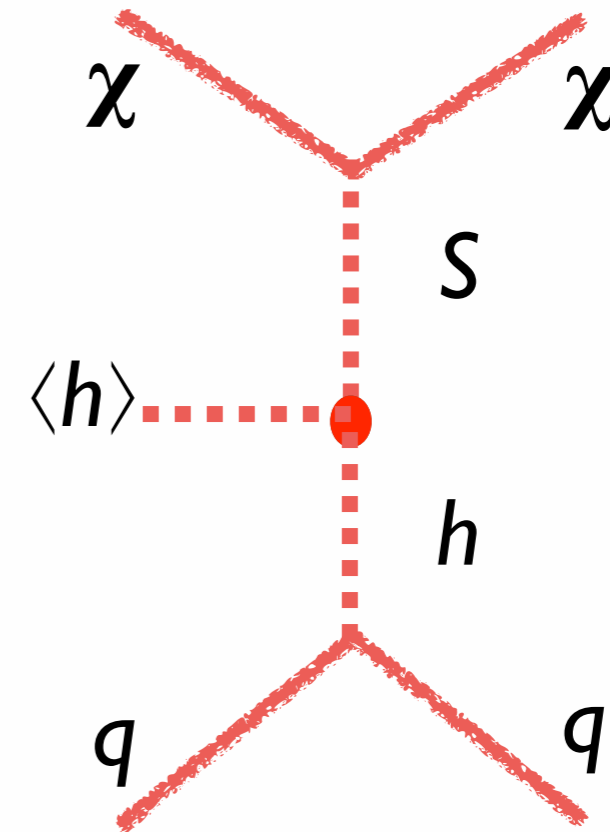


# What is 'thermal portal'?

What about this?



But the direct detection remains same!



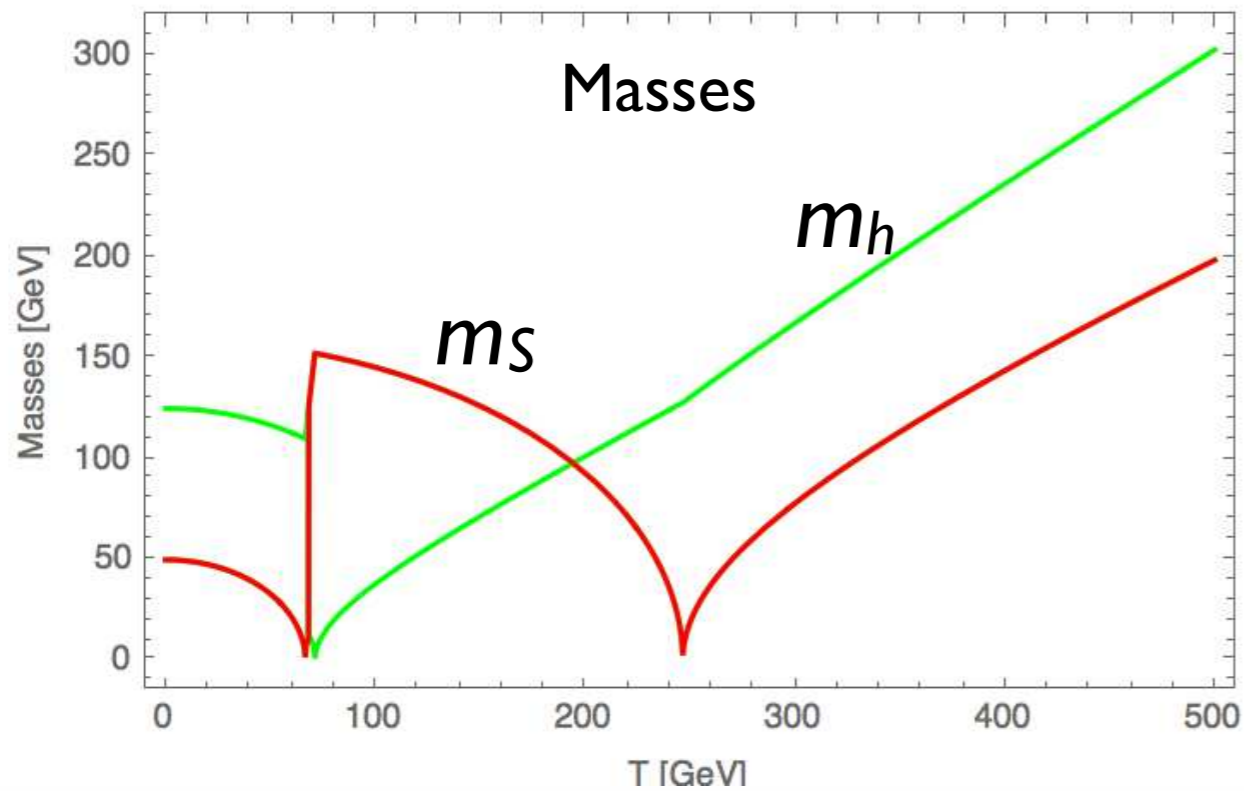
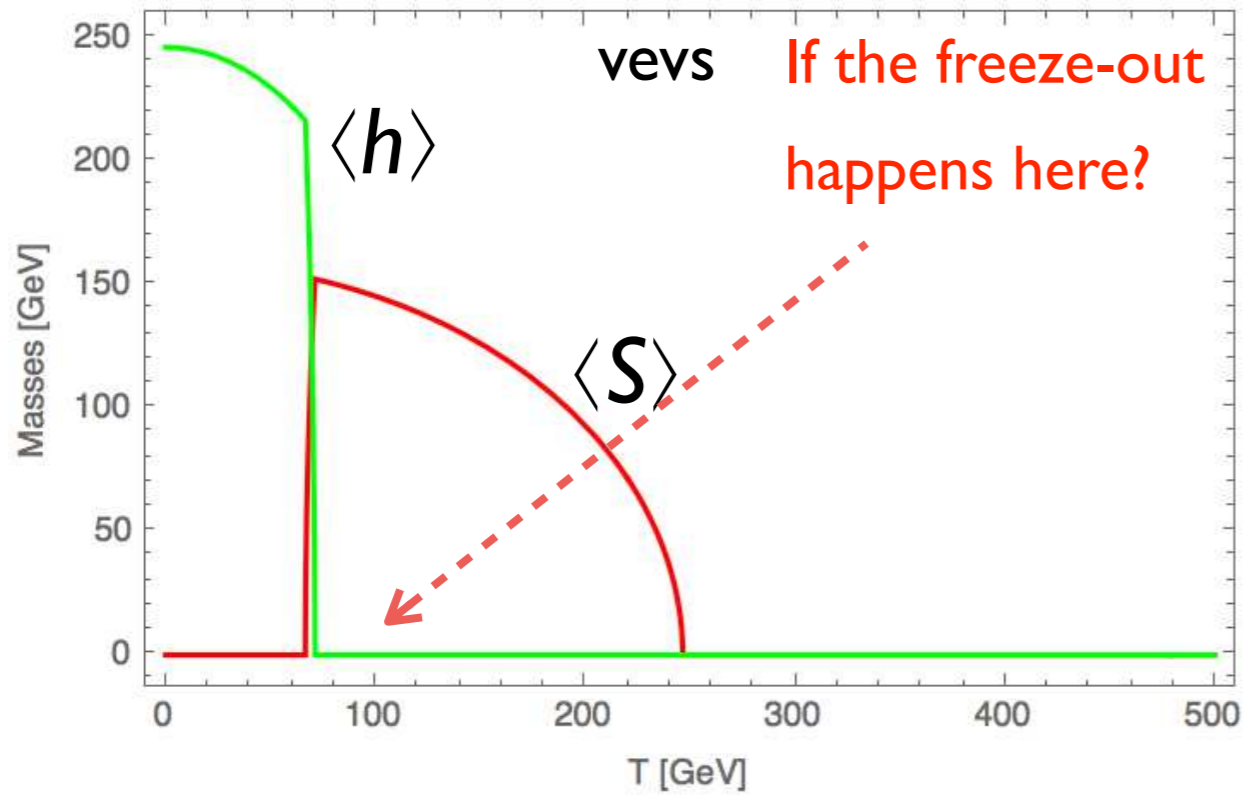
*Due to thermal corrections ( $T \gtrsim m_h, m_S$ )*

**NB!**  $\langle h \rangle_T \approx 0, \langle S \rangle_T > 0$

*(but in vacuum:  $\langle h \rangle_{T=0} > 0, \langle S \rangle_{T=0} = 0$ )*

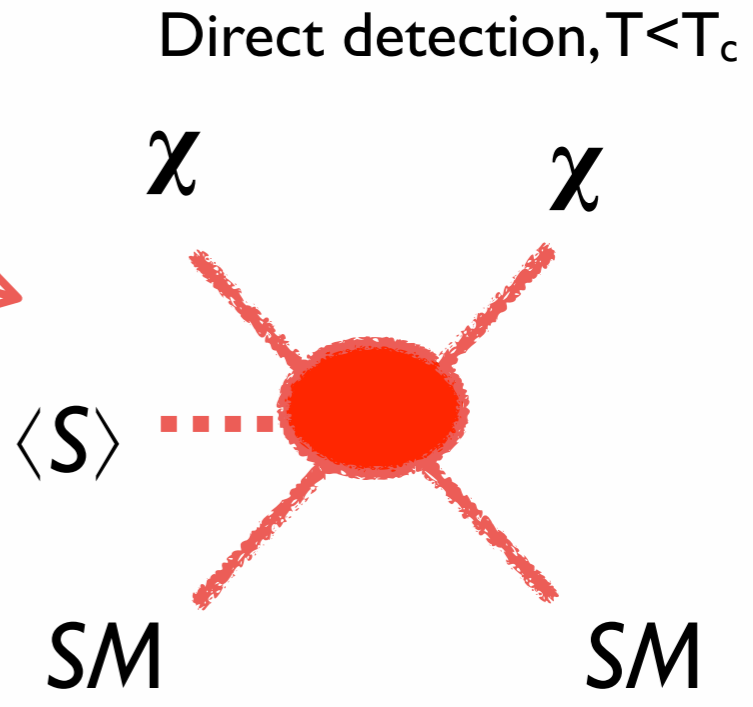
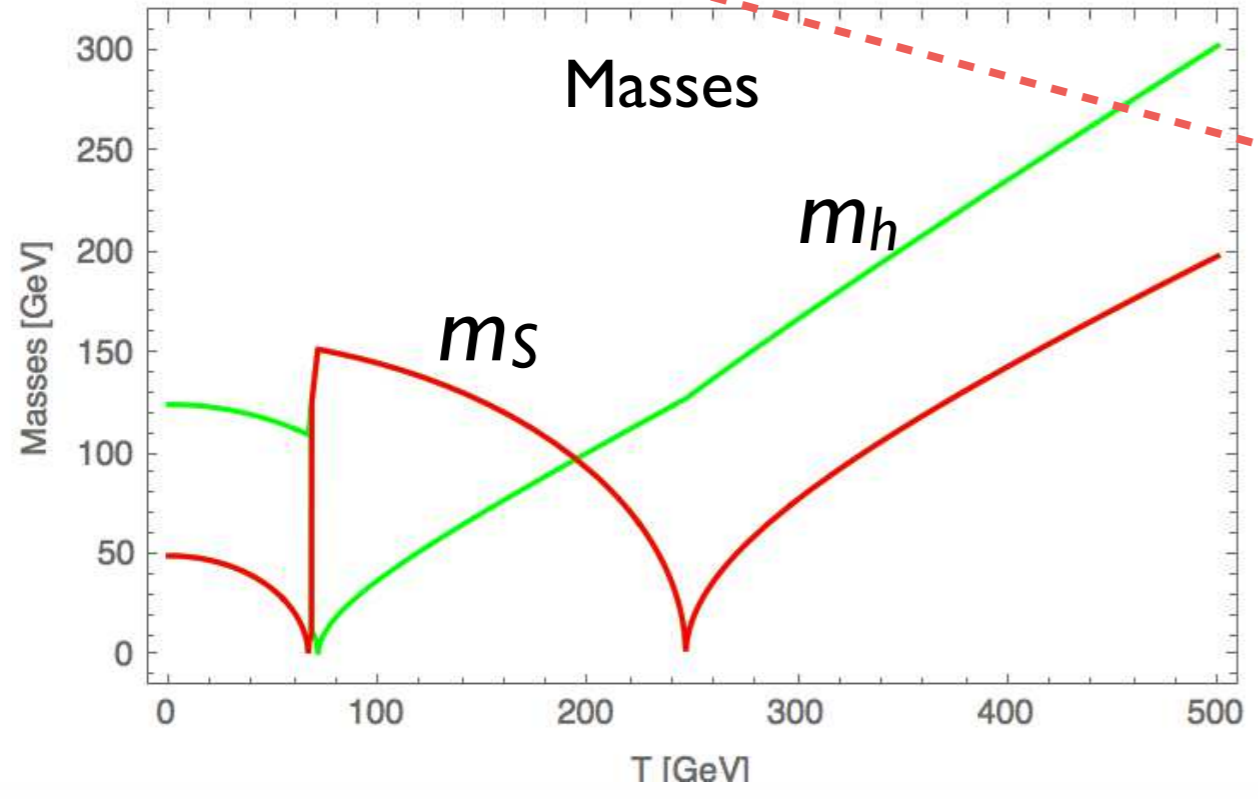
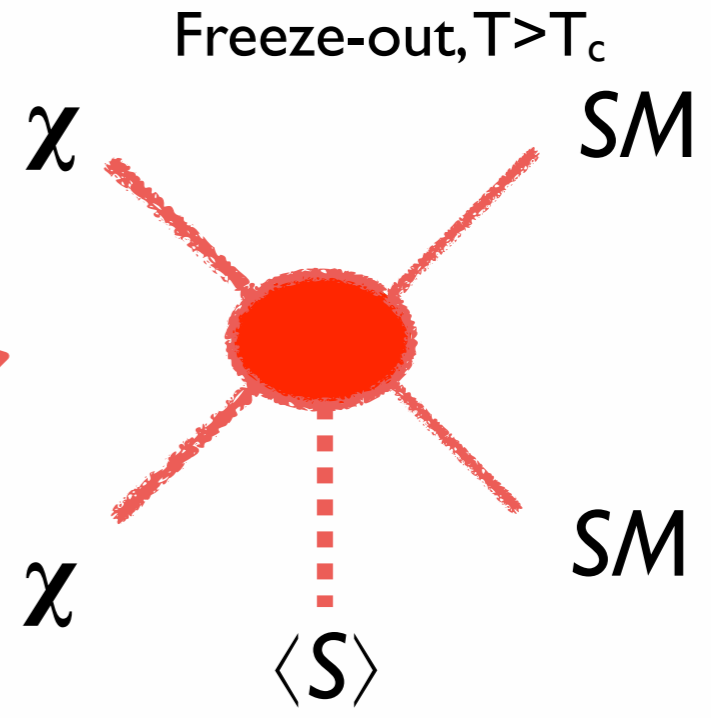
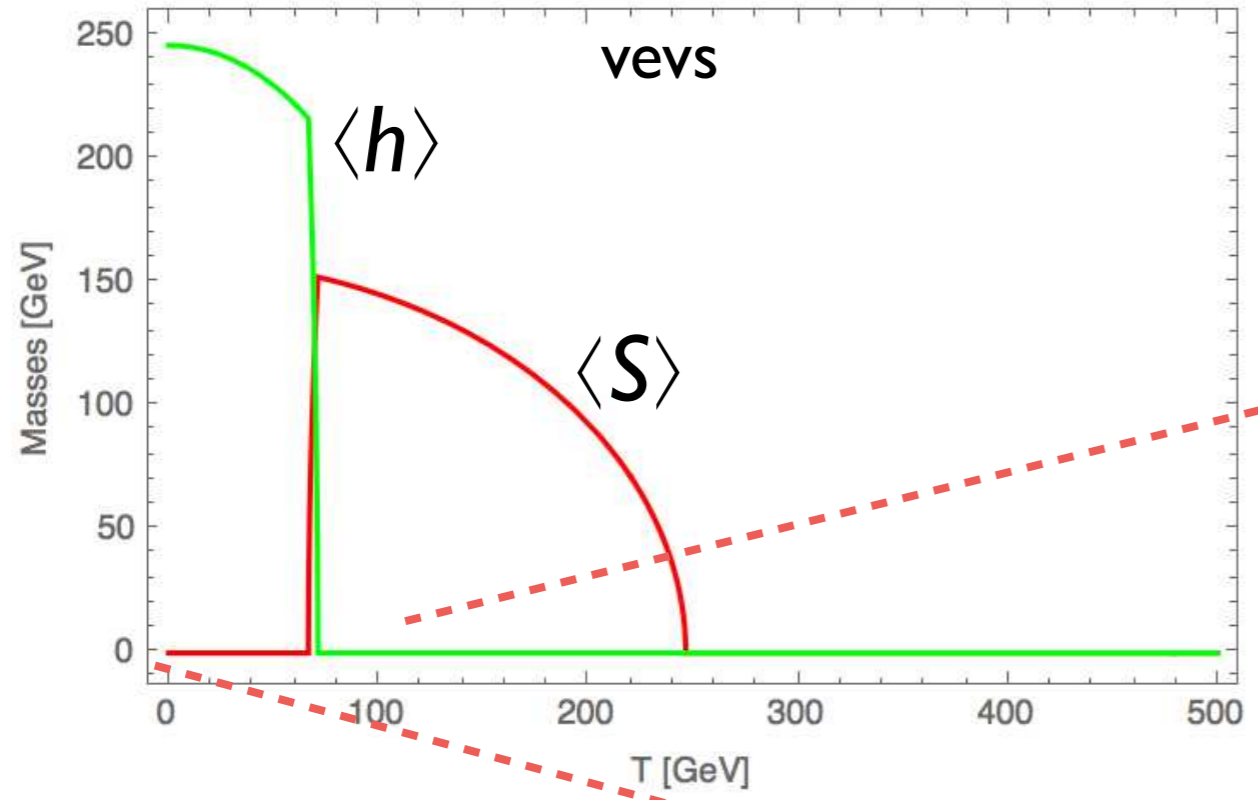


# Thermal running of the masses and vevs





# Thermal running of the masses and vevs





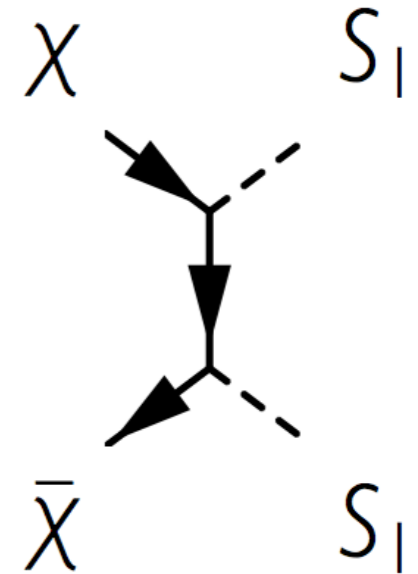
Let's build  
a 'thermal portal' model





# Considerations

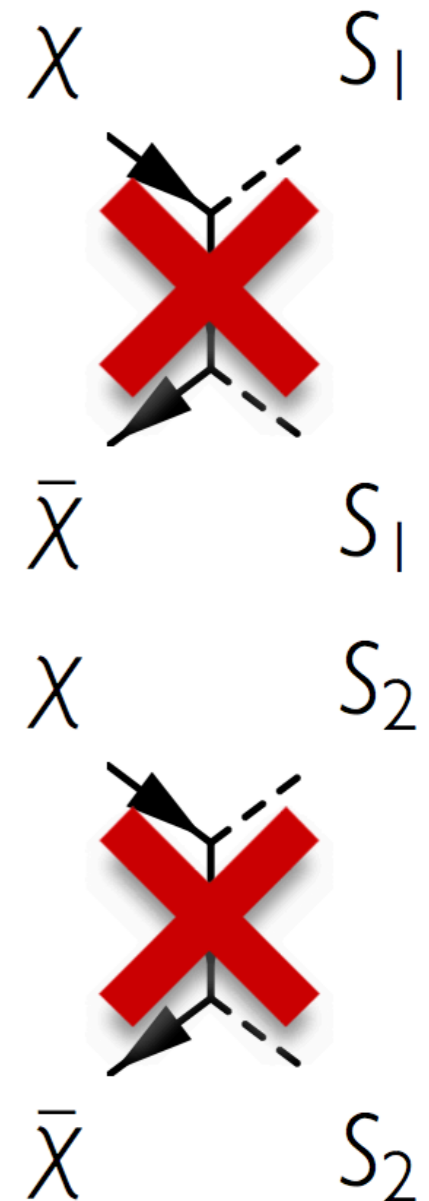
- Dark matter itself cannot be thermal due to its freeze-out:  $m_{\text{DM}} \gg T$
- Fermion dark matter with a thermal scalar
- Only one scalar does not work:  $t$ -channel annihilation works in any phase





# Considerations

- Dark matter itself cannot be thermal: it has to freeze out:  $m_{\text{DM}} \gg T$
- Fermion dark matter with a thermalised scalar singlet
- Only one scalar does not work:  $t$ -channel annihilation works in any phase
- Must have two scalars: light  $S_1$  with  $y_1 = 0$  and heavy  $S_2$  ( $m_2 > m_\chi$ )





# The minimal model

$$\begin{aligned} \mathcal{L} \supset & \bar{\chi} \not{\partial} \chi + |D_\mu H|^2 + \frac{1}{2} (\partial_\mu S_1)^2 + \frac{1}{2} (\partial_\mu S_2)^2 \\ & - m_\chi \bar{\chi} \chi - y_1 S_1 \bar{\chi} \chi - y_2 S_2 \bar{\chi} \chi - V, \end{aligned}$$

where

$$\begin{aligned} V = & \mu_H^2 |H|^2 + \frac{1}{2} \mu_{20}^2 S_1^2 + \frac{1}{2} \mu_{11}^2 S_1 S_2 + \frac{1}{2} \mu_{02}^2 S_2^2 \\ & + \lambda_H |H|^4 + \lambda_{H20} |H|^2 S_1^2 + \lambda_{H11} |H|^2 S_1 S_2 \\ & + \lambda_{H02} |H|^2 S_2^2 + \lambda_{40} S_1^4 + \lambda_{31} S_1^3 S_2 + \lambda_{22} S_1^2 S_2^2 \\ & + \lambda_{13} S_1 S_2^3 + \lambda_{04} S_2^4 \end{aligned}$$



# The thermal corrections

Thermal corrections to mass terms are given by

$$\delta m_{ij}^2 \approx \sum_k \frac{g_k}{24} \frac{\partial m_k^2}{\partial \phi_i \partial \phi_j} T^2,$$

where  $g_k = n_k$  for bosonic degrees of freedom



# The thermal corrections

$$\begin{aligned}\mu_H^2 &\rightarrow \mu_H^2 + c_H T^2, \\ \mu_{20}^2 &\rightarrow \mu_{20}^2 + c_{20} T^2,\end{aligned}$$

where

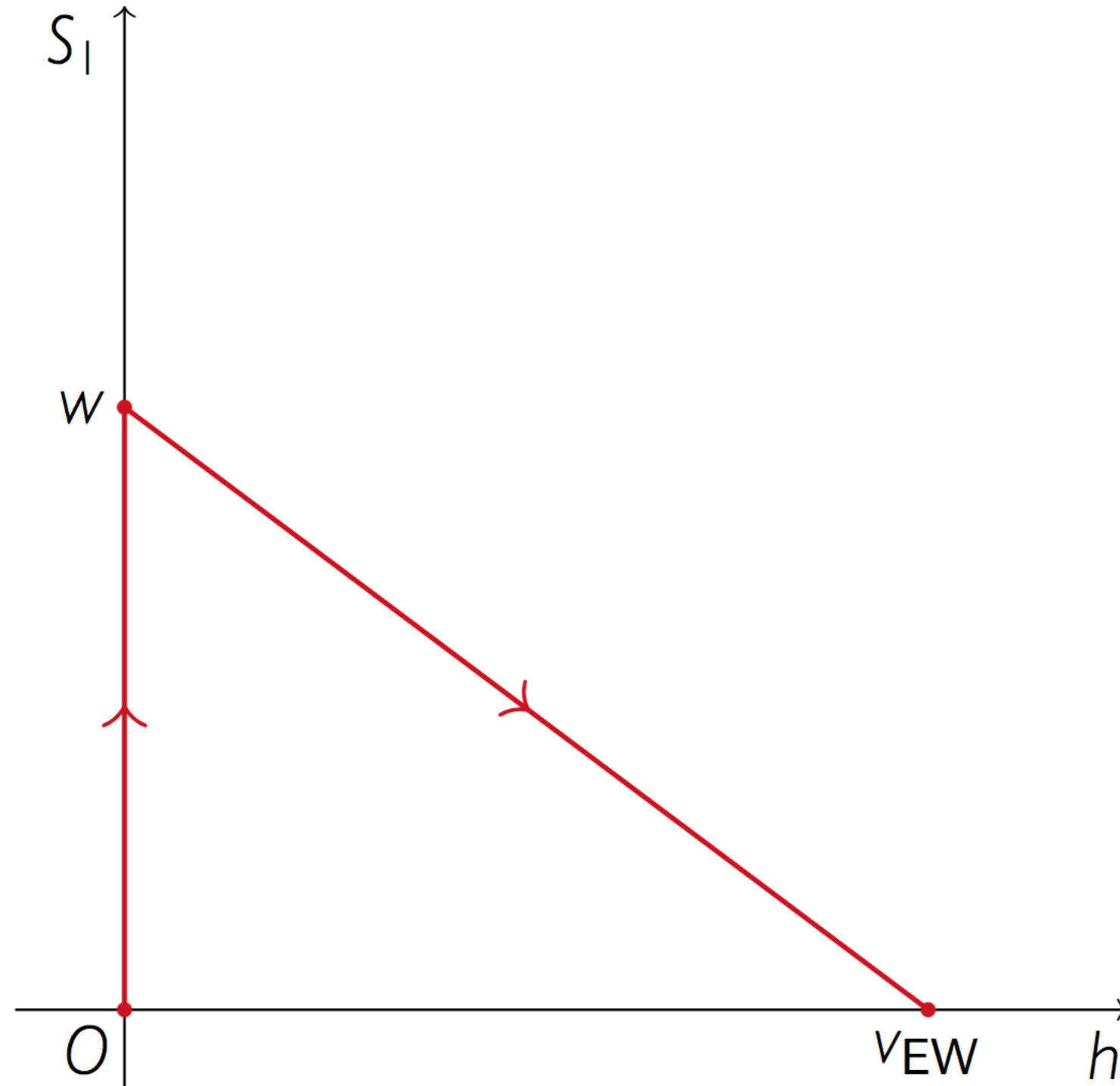
$$c_H = \frac{1}{48} (24\lambda_H + 3g'^2 + 9g^2 + 12y_t^2 + 4\lambda_{H20}),$$

$$c_{20} = \lambda_{40} + \frac{1}{3}\lambda_{H20}$$

- The contributions to the terms  $\mu_{11}^2 S_1 S_2$  and  $\frac{1}{2}\mu_{02}^2 S_2^2$  are approximately zero due to the high mass of  $S_2$

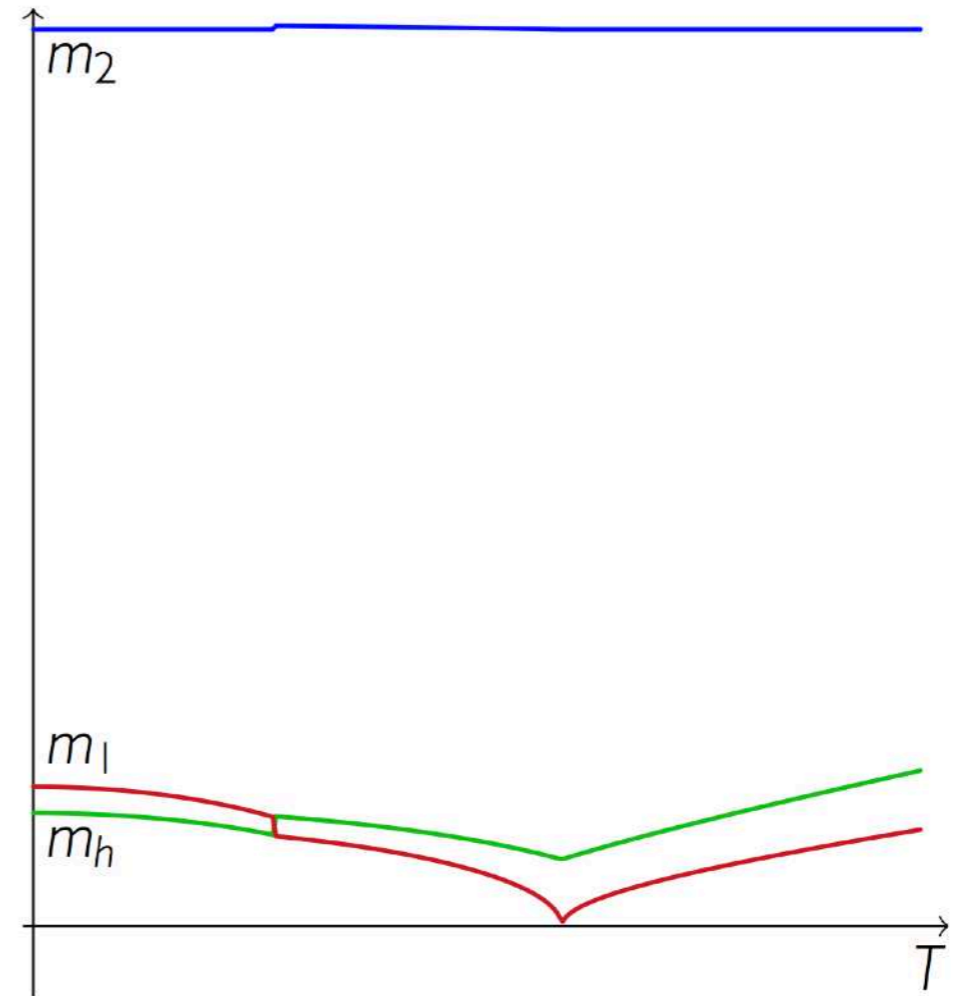
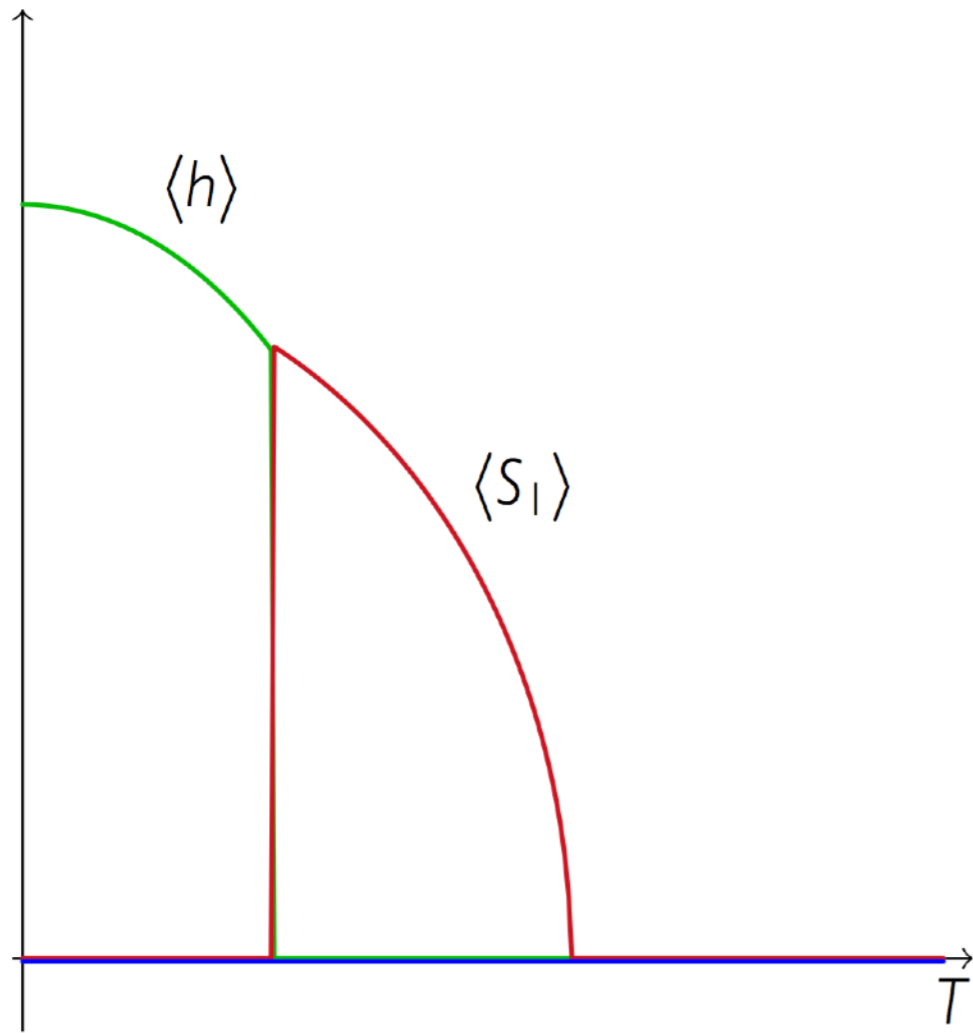


# Phase transition



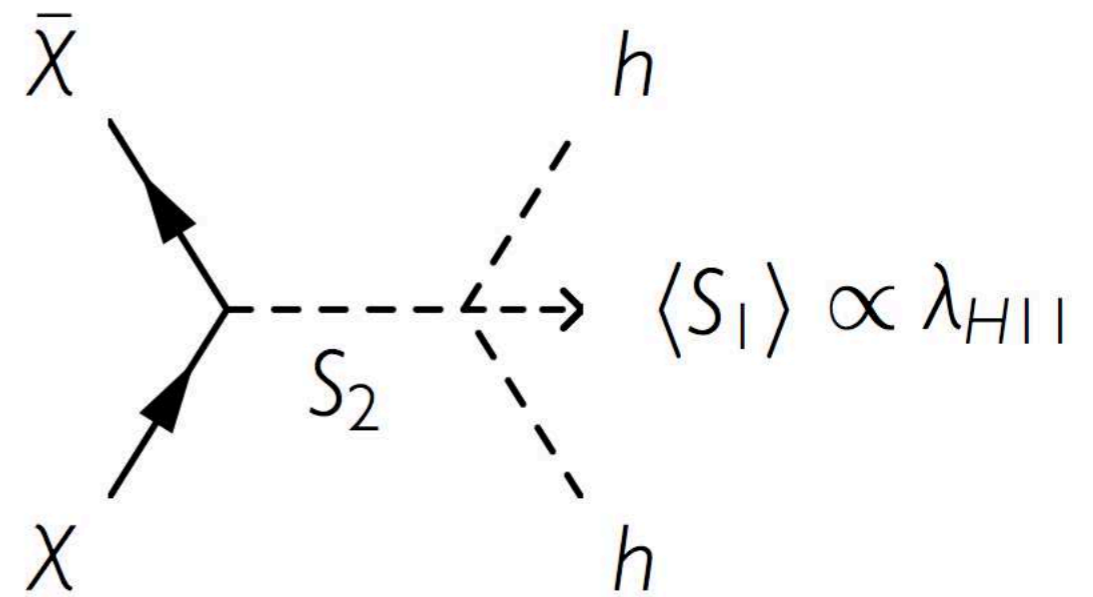
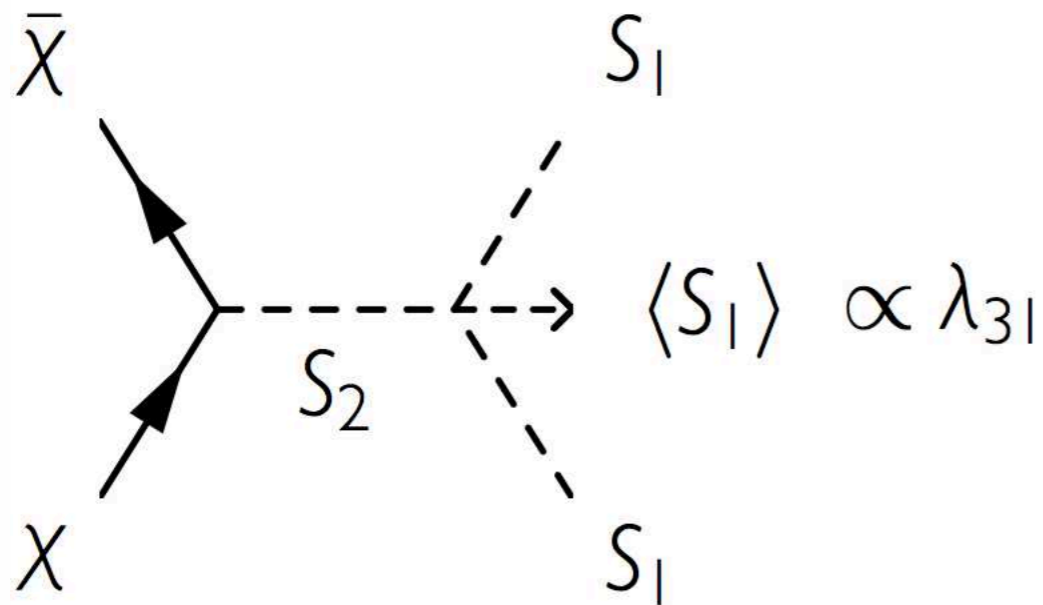
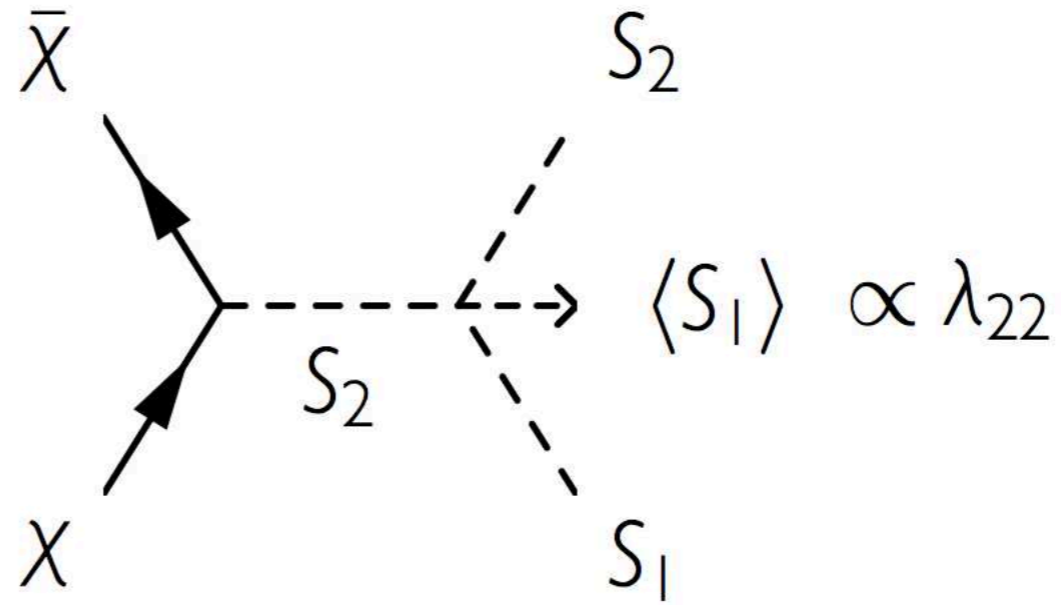


# Phase transition





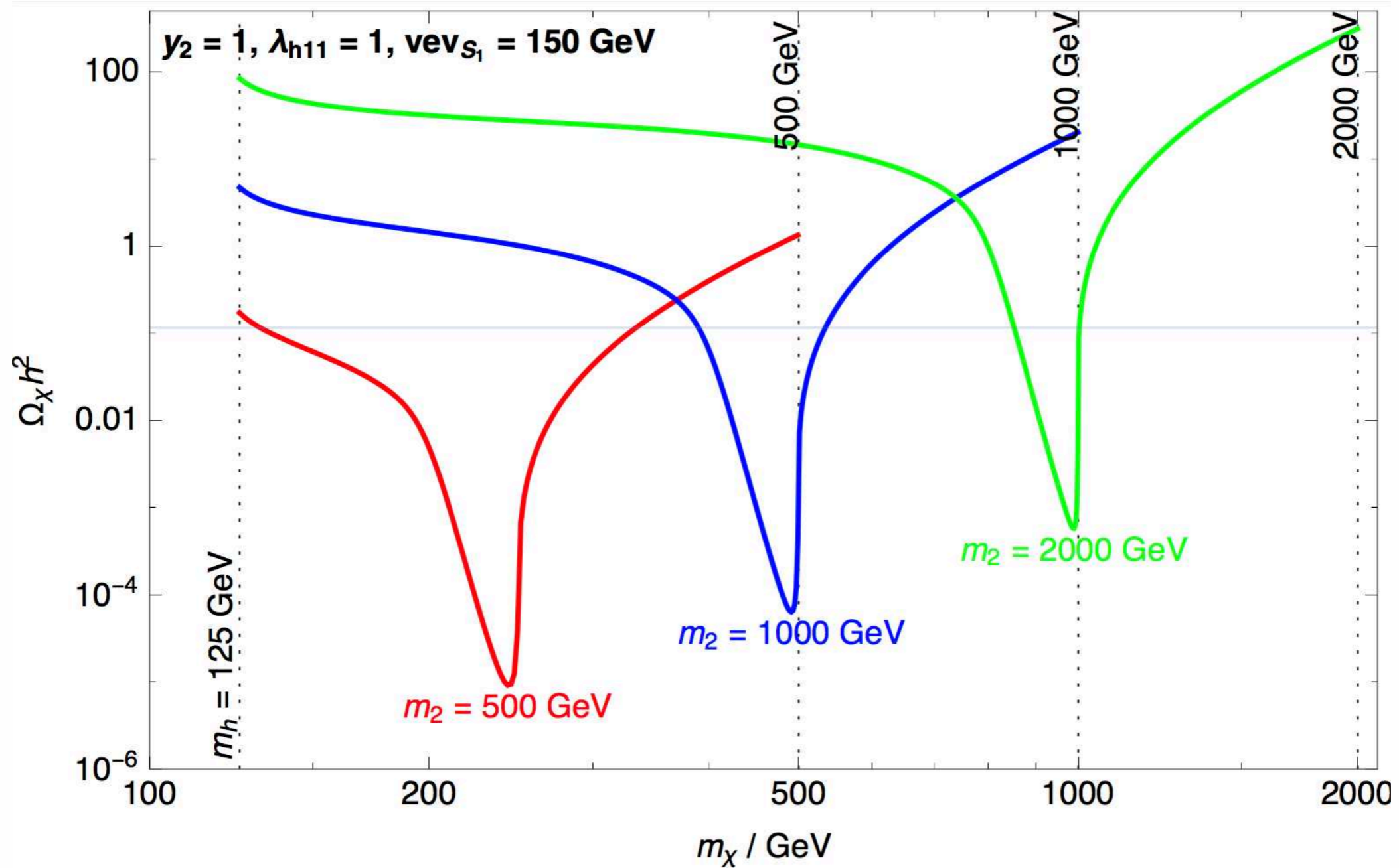
# Annihilation channels





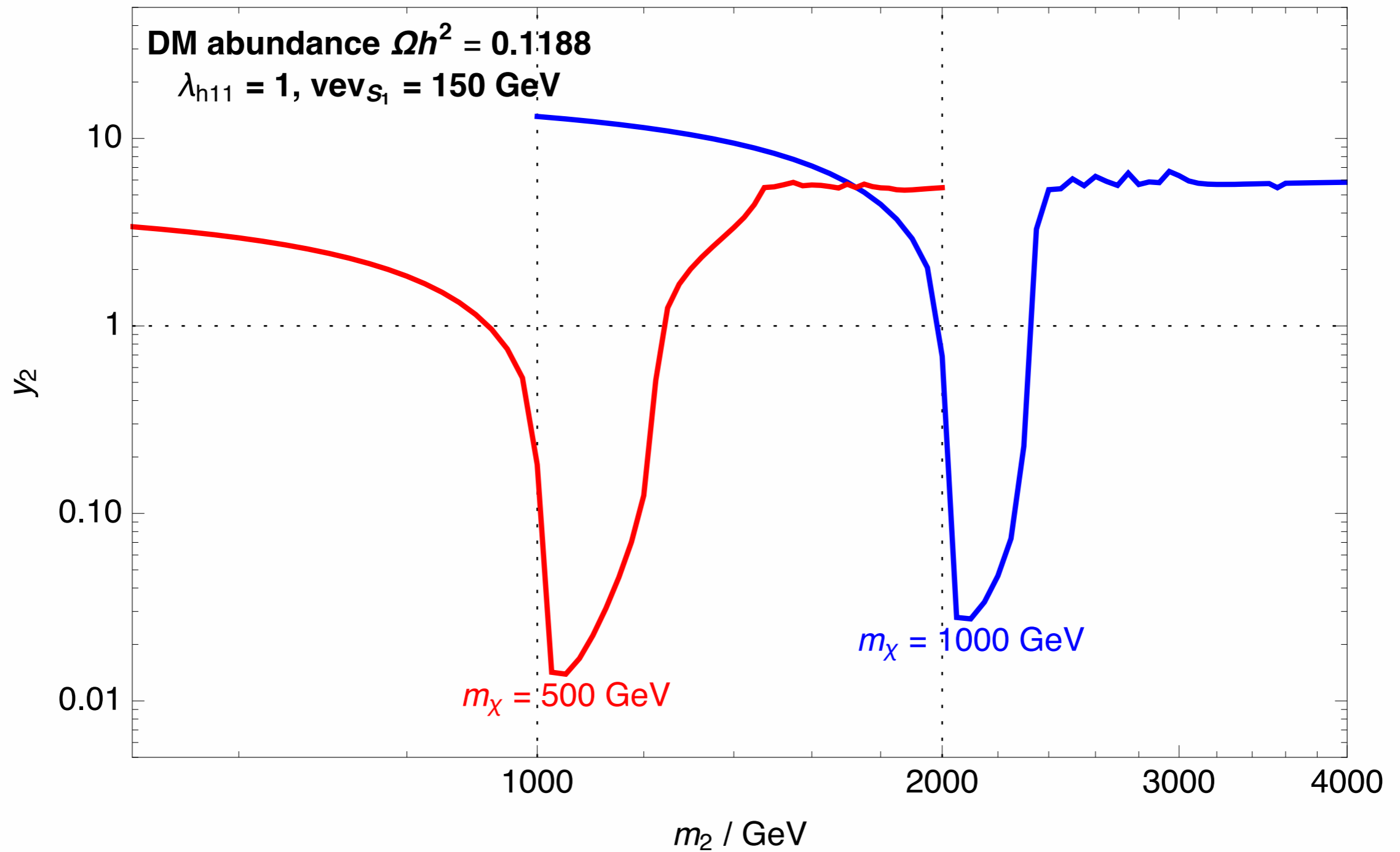


# Relict density & parameter space



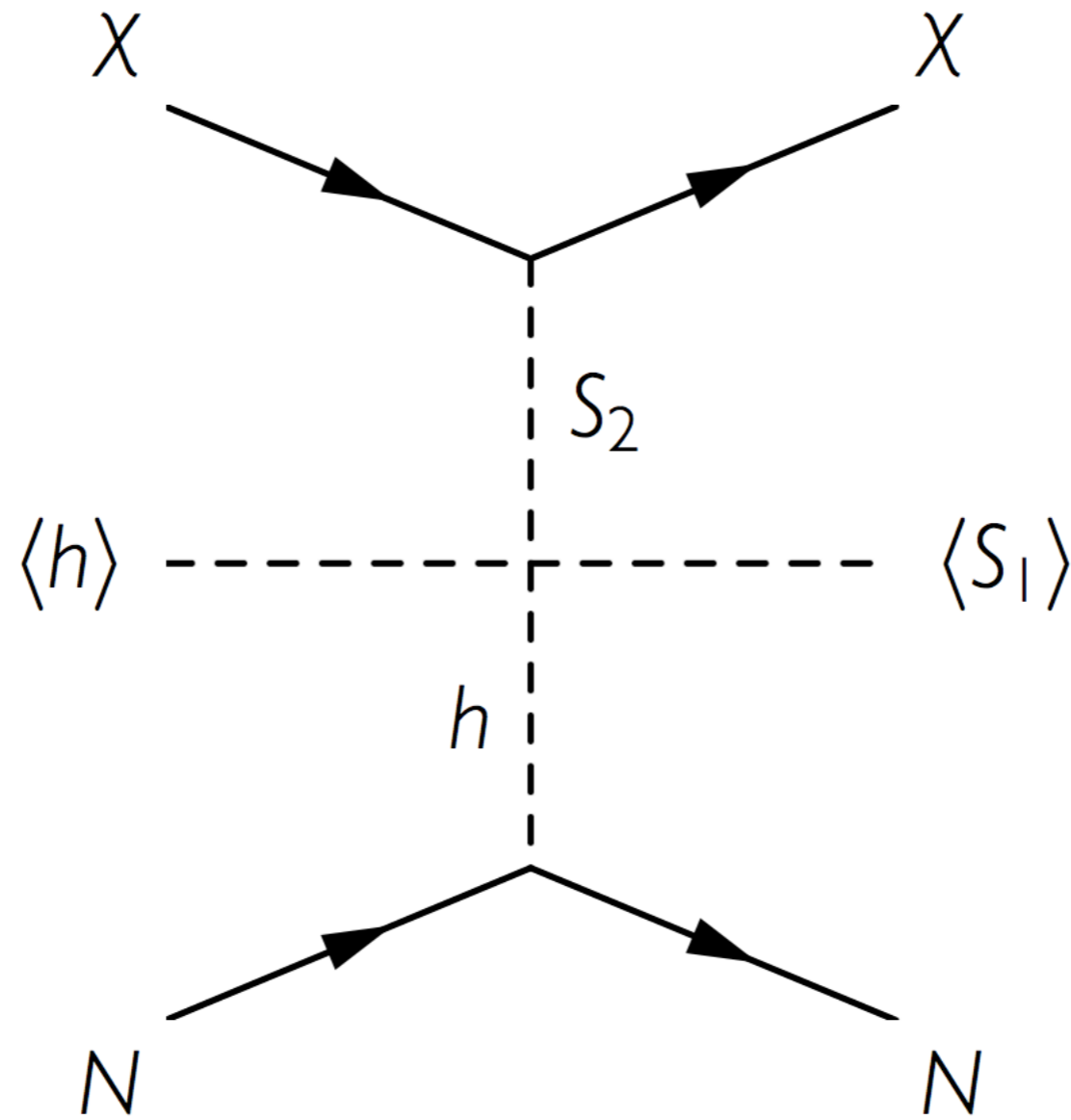


# Relict density & parameter space





# Direct detection





# Gravitational waves

- Gravitational wave production similar to the model of the SM with an EW singlet

Vaskonen 1611.02073 — work in progress



# Barogenesis





# Takeaway messages

- Non-zero temperature can have effect on dark matter freeze-out
- We have constructed a minimal model with freeze-out before phase transition
- While direct detection cross-section is negligible, there can be a gravitational wave signal



**Thank you!**