

Primordial Black Holes as Dark Matter

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Basics of Primordial Black Holes (PBHs)

Primordial Black Holes – Why are they interesting?

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- Can they constitute all DM?



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- In a matter-dominated Universe, perturbations grow as δ ∝ a ⇒ if there was enough time, even small perturbations can grow large, δ ~ 1. Most importantly, in MD the Jeans pressure does not prevent PBHs from forming









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- Q: Why do not PBHs form in today's Universe?
 A: They could, but their fraction would be negligible.
- Q: What is the origin of these perturbations?
 A: That is THE question!

Inflaton(s) (typically requires features in the inflaton potential)

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Rightmost image credit: David Weir / 1504.03291

PBHs as Dark Matter (1705.05567)

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normalised so that the fraction of the DM in PBHs is

$$f_{\rm PBH} \equiv rac{\Omega_{
m PBH}}{\Omega_{
m DM}} = \int {
m d} {\pmb M} \, \psi({\pmb M})$$

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- Monochromatic mass function: $\psi_{\text{mon}}(M) = f_{\text{PBH}}(M_c)\delta(M M_c)$
- All mass functions that are not monochromatic are extended mass functions
- In most scenarios, one can expect the mass function to be extended rather than a monochromatic one

Examples of extended mass functions I

A power-law mass function

$$\psi(M) \propto M^{\gamma-1} \quad (M_{\min} < M < M_{\max})$$

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Examples of extended mass functions II

A lognormal mass function

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Left: $\psi(M)$ with $M_c = 1$, $\sigma = 1$ (blue) and $\sigma = 1.4$ (orange) Right: $M\psi(M)$ with the same parameters as above

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Constraints on PBHs



Constraints: evaporations; femtolensing of gamma-ray bursts (FL); neutron star capture (NS); white dwarf explosions (WD); the microlensing results from Subaru (HSC), Kepler (K), EROS and MACHO (M); Planck; survival of stars in Segue I (Seg I) and Eridanus II (Eri II); distribution of wide binaries (WB).

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Constraints on PBHs with an extended mass function



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PBHs as Dark Matter



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 \Rightarrow Too faint to be detected by aLIGO–Virgo?

Can we learn something from the constraints on PBHs? (1706.03746)

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- Let us assume that the are two components that contribute to the curvature power spectrum: the inflaton φ and a spectator field s

$$\mathcal{P}_{\mathcal{R}}(k) = \mathcal{P}_{\mathcal{R},\varphi}(k) + \mathcal{P}_{\mathcal{R},s}(k)$$

Primordial Power Spectrum with two components

The inflaton perturbations produce a nearly flat spectrum at small k,

$$\mathcal{P}_{\mathcal{R},\varphi}(k) = A\left(\frac{k}{k_*}\right)^{n-1+\frac{1}{2}\mathrm{d}n/\mathrm{d}\ln k \ln\left(\frac{k}{k_*}\right)},$$

where k_* is a pivot scale, $A \simeq 10^{-9}$, and $n \simeq 0.968$.

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Perturbations in the s field dominate at large k

$$\mathcal{P}_{\mathcal{R},s}(k) = A_s \left(\frac{k}{k_*}\right)^{n_s - 1 + \frac{1}{2} \mathrm{d}n_s/\mathrm{d}\ln k \ln\left(\frac{k}{k_*}\right)}$$

unless the running of the inflaton field's spectral index is considerable.

The total power spectrum



The total curvature power spectrum (black solid lines).

Black dashed line: inflaton. Grey dashed lines: spectator field for different choices of parameters.

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- Possible cause: reheating, massive metastable particles...
- ▶ PBH formation starts when δ grows large enough ($\delta \sim$ 1) and ends when the MD ends

Constraints on running of the spectral tilt

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For constraints on spectral features of the spectator field and PBH DM, see 1706.03746.

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Matter or radiation domination?

- When does the transition occur?
- ► What is the effect of a subdominant radiation bath? ⇒ 1804.08639





Conclusions

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- Constraints on monochromatic PBH mass functions have to be carefully adopted for extended mass functions
- PBHs provide for an effective way to constrain curvature perturbations at small scales
- Subdominant radiation bath can have a larger effect on PBH formation than previously thought