



Collider prospects for vector Dark Matter coupled with the Higgs

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Based on G.A., A. Djouadi, M. Kado 2001.10750
plus work in preparation

PARTICLE DM

Dark Matter is one of the building blocks of the Standard Cosmological model.
Provides most of the matter contribution to the energy budget of the Universe.
Evidences from astrophysics and cosmology.

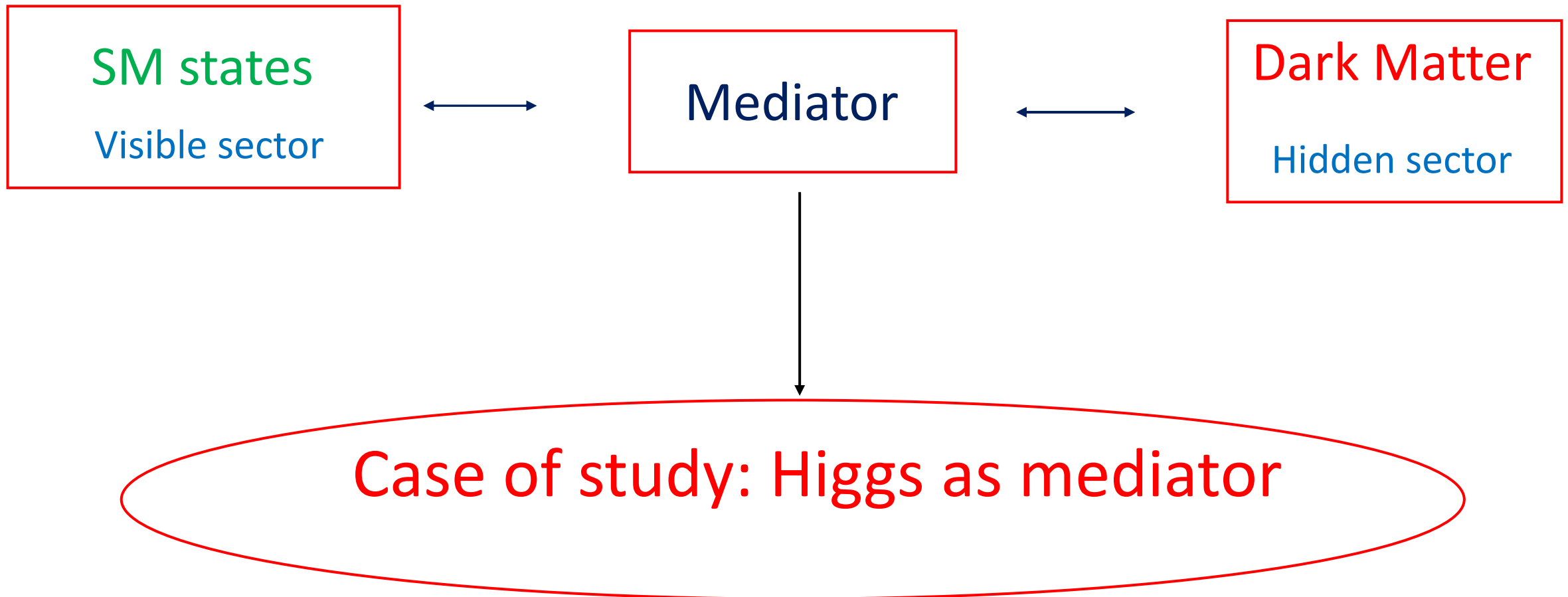
Stable on **cosmological scales**.

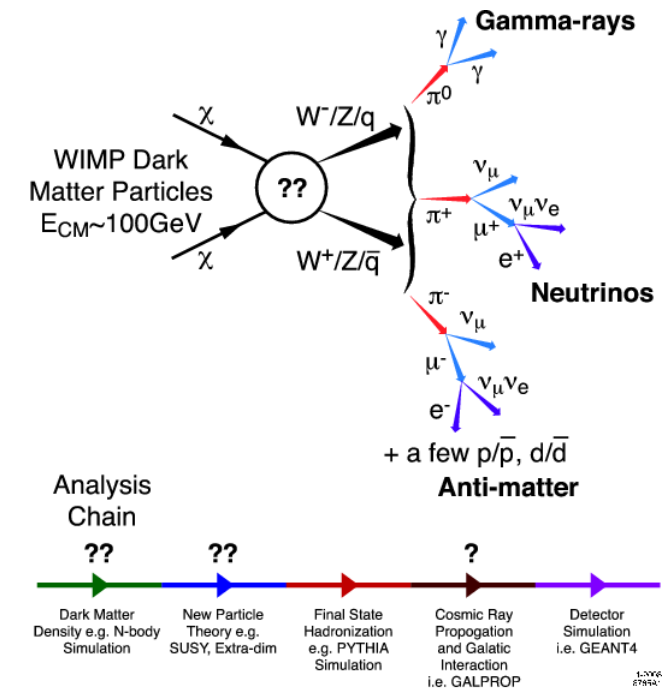
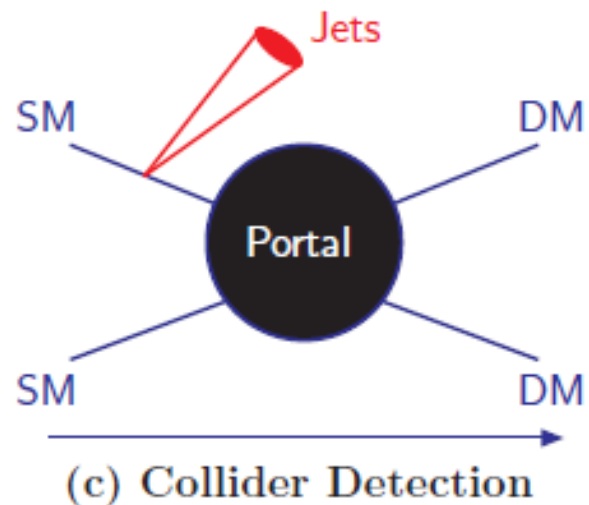
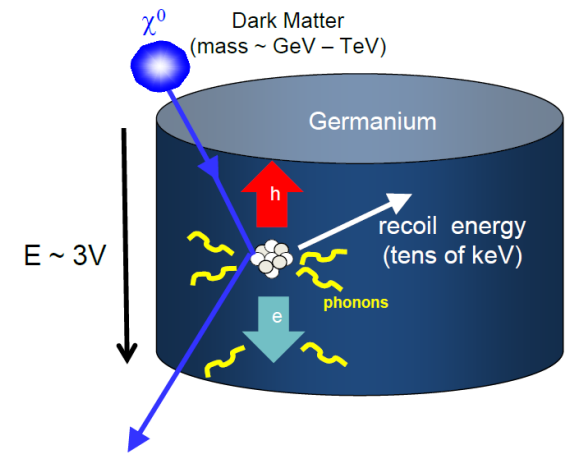
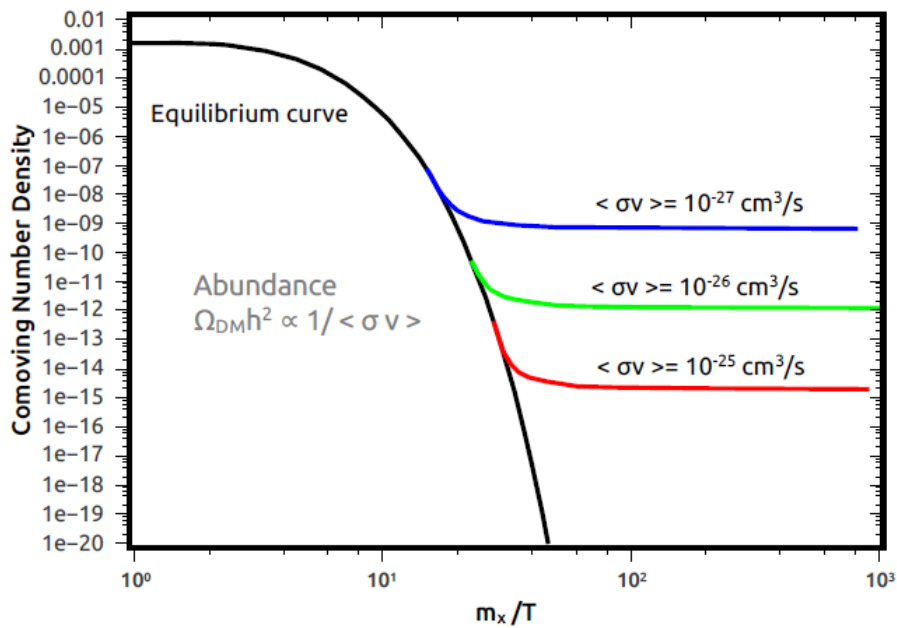
Weakly or SuperWeakly interacting with ordinary matter, photons.

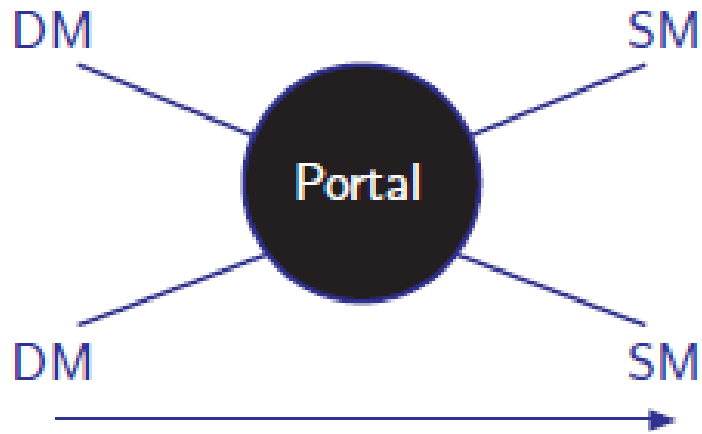
Cold (up to **warm**) as opposed to **hot**.

No (confirmed) detection so far.

“Portals”

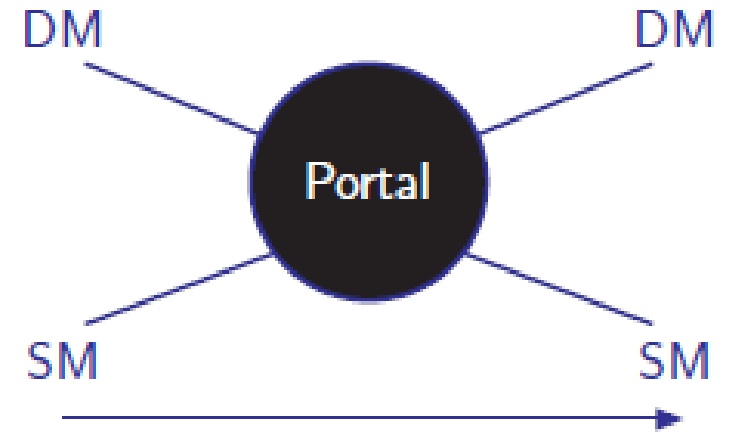






$$\Omega h^2 \approx 0.12$$

$$\langle \sigma v \rangle \approx 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$



Relic Density

$$\langle \sigma v \rangle \approx \frac{\lambda_{DM}^2 \lambda_{SM}^2 m_{DM}^2}{(4 m_{DM}^2 - m_{Portal}^2)^2} \left(a + \frac{b}{x} \right) \approx \frac{\lambda_{DM}^2 \lambda_{SM}^2 m_{DM}^2}{(4 m_{DM}^2 - m_{Portal}^2)^2} \left(a + \frac{1}{2} b v_{f.o.}^2 \right) \longleftrightarrow \sigma_{\chi N} \approx \frac{\mu_{\chi N}^2 \lambda_{DM}^2}{\pi m_{Portal}^4} f(\lambda_q)$$

$$v_{f.o.} \sim 0.3$$

Indirect Detection

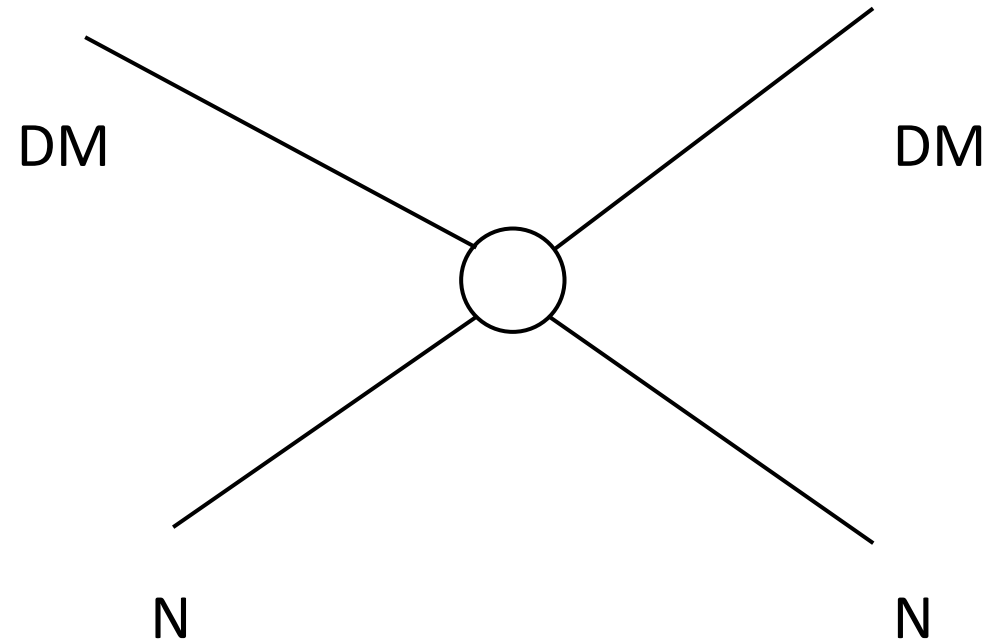
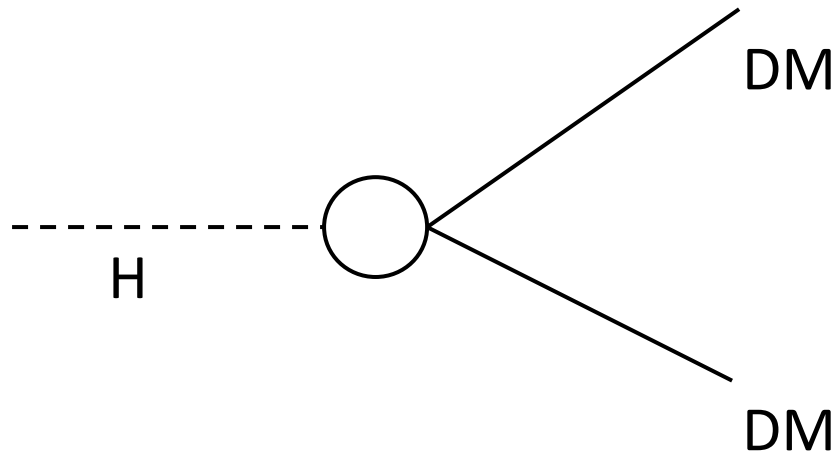
$$\langle \sigma v \rangle \approx \frac{\lambda_{DM}^2 \lambda_{SM}^2 m_{DM}^2}{(4 m_{DM}^2 - m_{Portal}^2)^2} \left(a + \frac{1}{2} b v_{now}^2 \right)$$

$$v_{now} \sim 10^{-3}$$

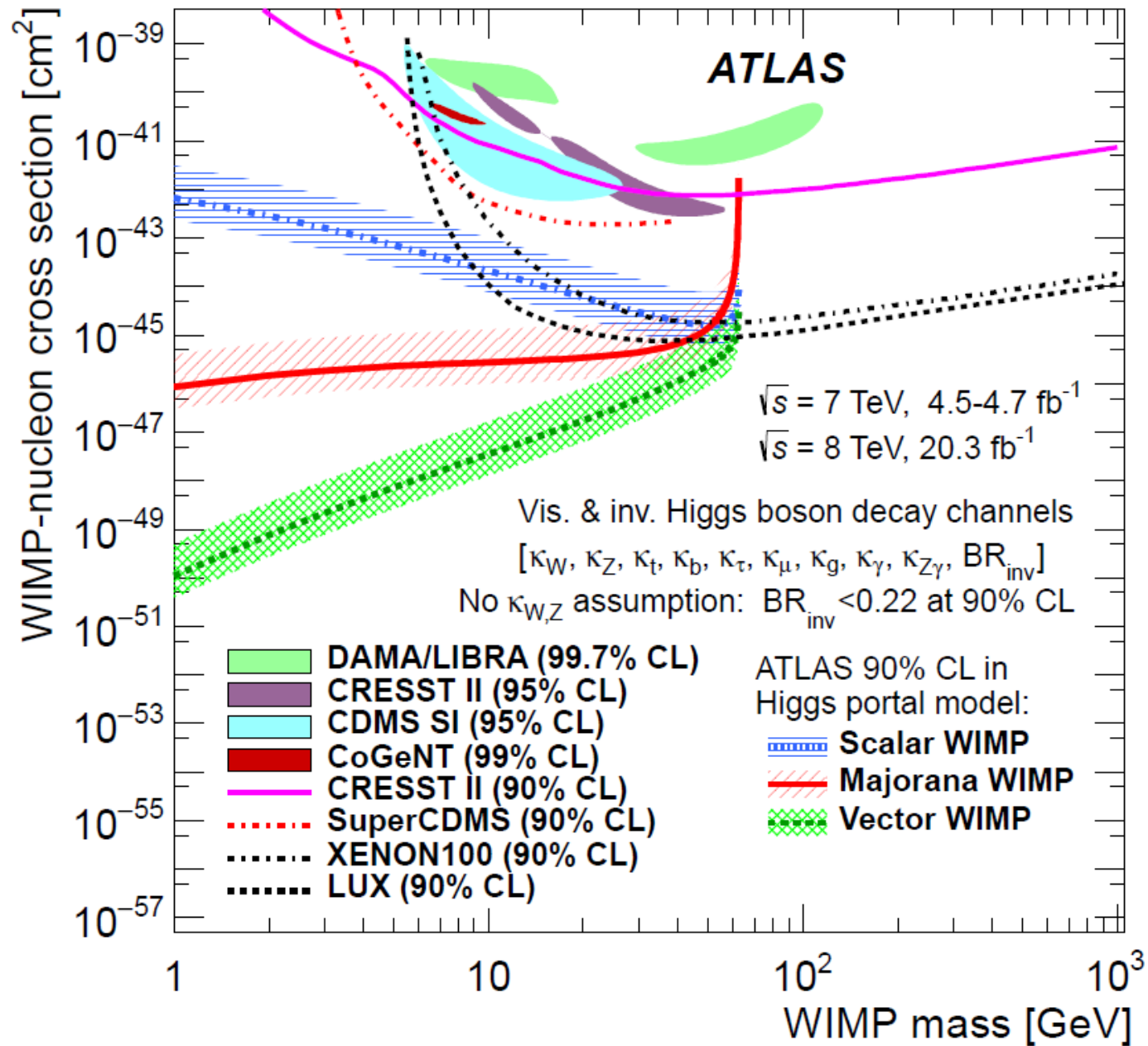
s-wave

p-wave

Invisible H decay vs Direct Detection



$$\sigma_{DM,N}^{SI} = \frac{\mu_{DM,p}^2}{\pi M_H^4} \lambda_{HDMDM}^2 \lambda_{HNN}^2 \rightarrow \sigma_{DM,N}^{SI} \propto Br(H \rightarrow DMDM)$$



The LHC correlation plot appears to be very powerful....

Some relevant questions arise:

- Is the picture full theoretically consistent?
- Can it also describe more complete models?
- Which is the impact of requiring the correct relic density?

(Effective) Higgs portal

D=4, but potential issues of perturbative unitarity

$$\Delta\mathcal{L}_S = -\frac{1}{2}M_S^2 S^2 - \frac{1}{4}\lambda_S S^4 - \frac{1}{4}\lambda_{HSS}\Phi^\dagger\Phi S^2 ,$$

$$\Delta\mathcal{L}_V = \frac{1}{2}M_V^2 V_\mu V^\mu + \frac{1}{4}\lambda_V (V_\mu V^\mu)^2 + \frac{1}{4}\lambda_{HVV}\Phi^\dagger\Phi V_\mu V^\mu$$

$$\Delta\mathcal{L}_\chi = -\frac{1}{2}M_\chi\bar{\chi}\chi - \frac{1}{4}\frac{\lambda_{H\chi\chi}}{\Lambda}\Phi^\dagger\Phi\bar{\chi}\chi .$$

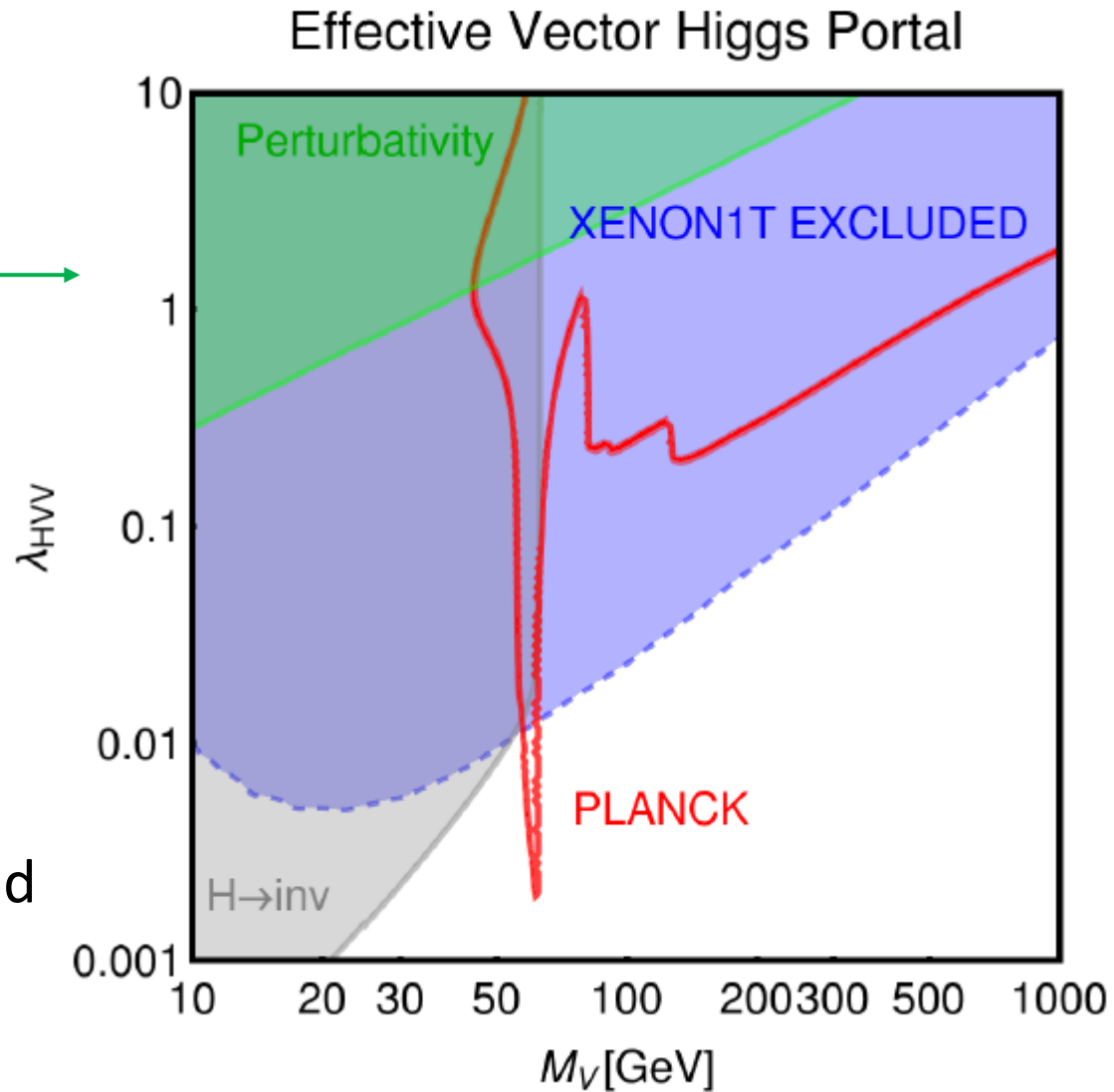
Explicitly $D \geq 4$. Depends on a NP scale Λ .

Perturbative unitarity

$$M_V \geq \frac{\lambda_{HVV} v}{\sqrt{8\pi}}$$

O. Lebedev, Y. M. Lee, Y. Mambrini
Phys.Lett.B 707 (2012) 570-576

Besides experimental constraints,
thermal DM is not viable below around
45 GeV.



Dark U(1)

Let's compare the effective vector Higgs portal with a more complete model.

$$\mathcal{L}_{\text{hidden}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (D_{\mu}\phi)^{\dagger}D^{\mu}\phi - V(\phi)$$

U(1) spontaneously broken \longrightarrow $A_{\mu} \rightarrow -A_{\mu}$

$$\mathcal{L}_{\text{portal}} = -\lambda_{h\phi}|H|^2|\phi|^2$$

Spontaneous breaking leaves a residual Z_2 symmetry

$$\begin{aligned}\rho &= -h_1 \sin \theta + h_2 \cos \theta \\ h &= h_1 \cos \theta + h_2 \sin \theta ,\end{aligned}$$

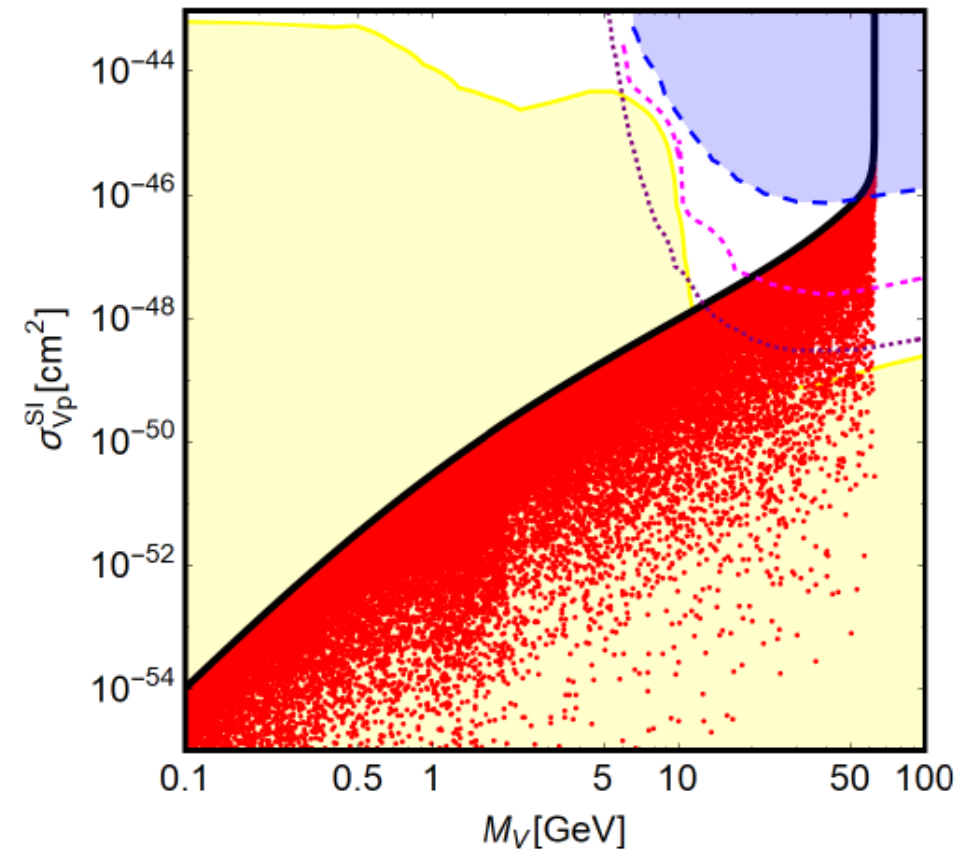
$$\Delta\mathcal{L}_{\text{s-g}} = \frac{\tilde{g}^2}{4}\tilde{v}\rho A_{\mu}A^{\mu} + \frac{\tilde{g}^2}{8}\rho^2 A_{\mu}A^{\mu}$$

As can be imagined, for $M_{H_2} \gg M_{H_1}$

$$\sigma_{Vp}^{SI} \Big|_{EFT} = 32 \mu_{Vp}^2 \frac{M_V^2}{M_H^3} \frac{Br(H \rightarrow VV) \Gamma_H^{\text{tot}}}{\beta_{VH}} \frac{1}{M_H^4} \frac{m_p^2}{v^2} |f_p|^2$$

$$\sigma_{Vp}^{SI} \Big|_{U(1)} = 32 \cos^2 \theta \mu_{Vp}^2 \frac{M_V^2}{M_{H_1}^3} \frac{Br(H \rightarrow VV) \Gamma_{H_1}^{\text{tot}}}{\beta_{VH_1}} \left(\frac{1}{M_{H_2}^2} - \frac{1}{M_{H_1}^2} \right)^2 \frac{m_p^2}{v^2} |f_p|^2$$

$$r = \frac{\sigma_{U(1)}^{SI}}{\sigma_{EFT}^{SI}} = 1 \longrightarrow \cos^2 \theta \left(\frac{1}{M_{H_2}^2} - \frac{1}{M_{H_1}^2} \right) \approx 1$$

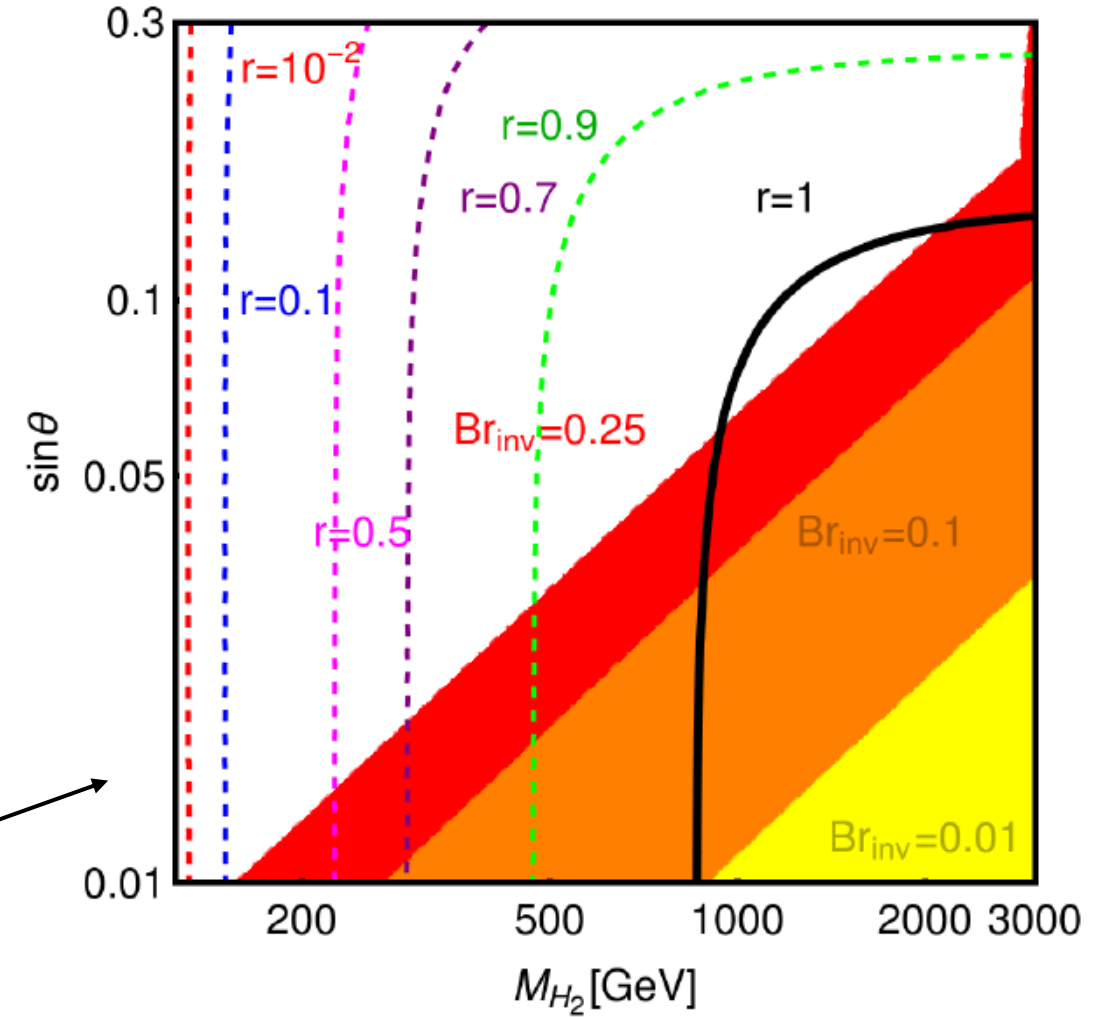


Perturbative Unitarity

$$\lambda_{HS} \leq \frac{4\pi}{3} \implies \text{BR}(H_1 \rightarrow VV) \lesssim 0.25 \left(\frac{3 \text{ TeV}}{M_{H_2}} \right)^4,$$

$$\lambda_S \leq \frac{4\pi}{3} \implies \text{BR}(H_1 \rightarrow VV) \lesssim 0.35 \left(\frac{\sin \theta}{0.1} \right)^2 \left(\frac{3 \text{ TeV}}{M_{H_2}} \right)^2$$

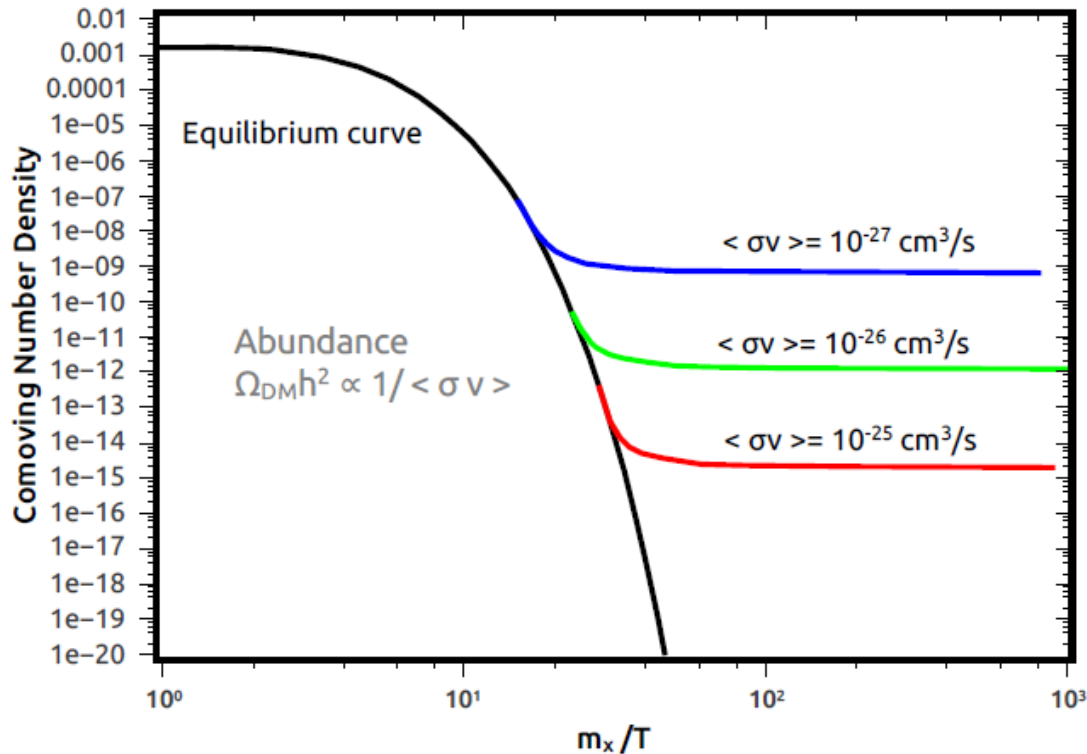
$$r = \frac{\sigma_{U(1)}^{SI}}{\sigma_{EFT}^{SI}}$$



G. Arcadi, A. Djouadi, M. Kado
Phys.Lett.B 805 (2020) 135427

Include DM relic density

We assume the freeze-out paradigm:



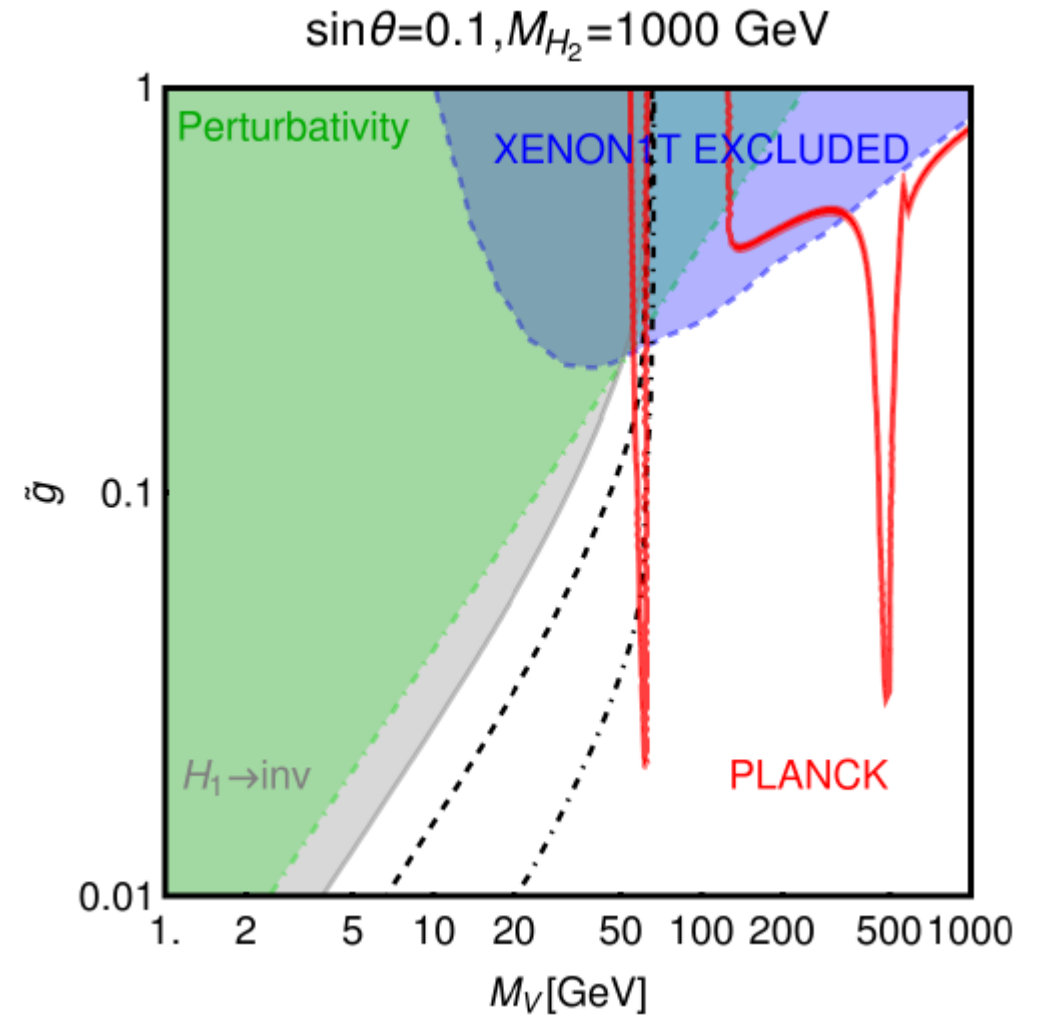
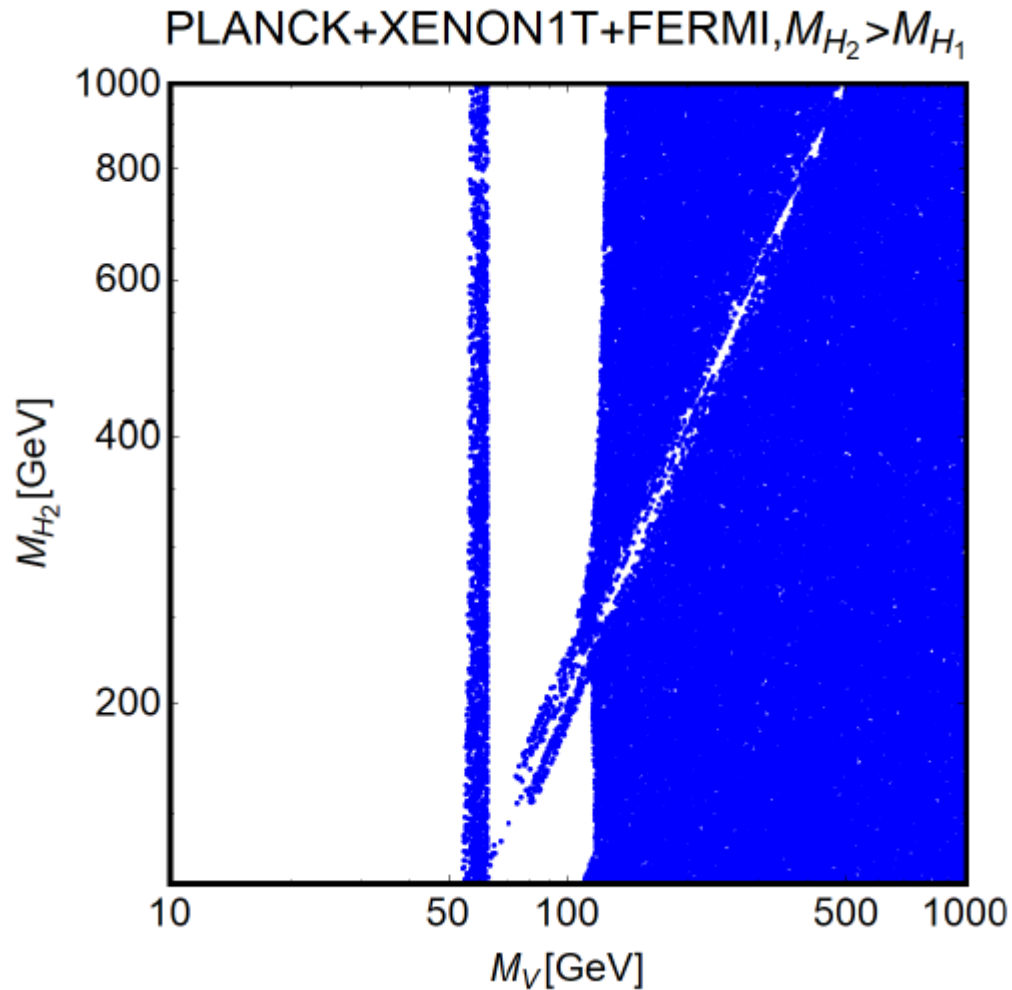
$$\frac{dn_{DM}}{dt} + 3Hn_{DM} = -\langle \sigma v \rangle (n_{DM}^2 - n_{DM,eq}^2)$$

$$\langle \sigma v \rangle = \frac{1}{8 m_{DM}^4 T K_2 \left(\frac{m_{DM}}{T} \right)^2} \int_{4 m_{DM}^2}^{\infty} ds \sqrt{s} (s - 4 m_{DM}^2) \sigma(s) K_1 \left(\frac{\sqrt{s}}{T} \right)$$

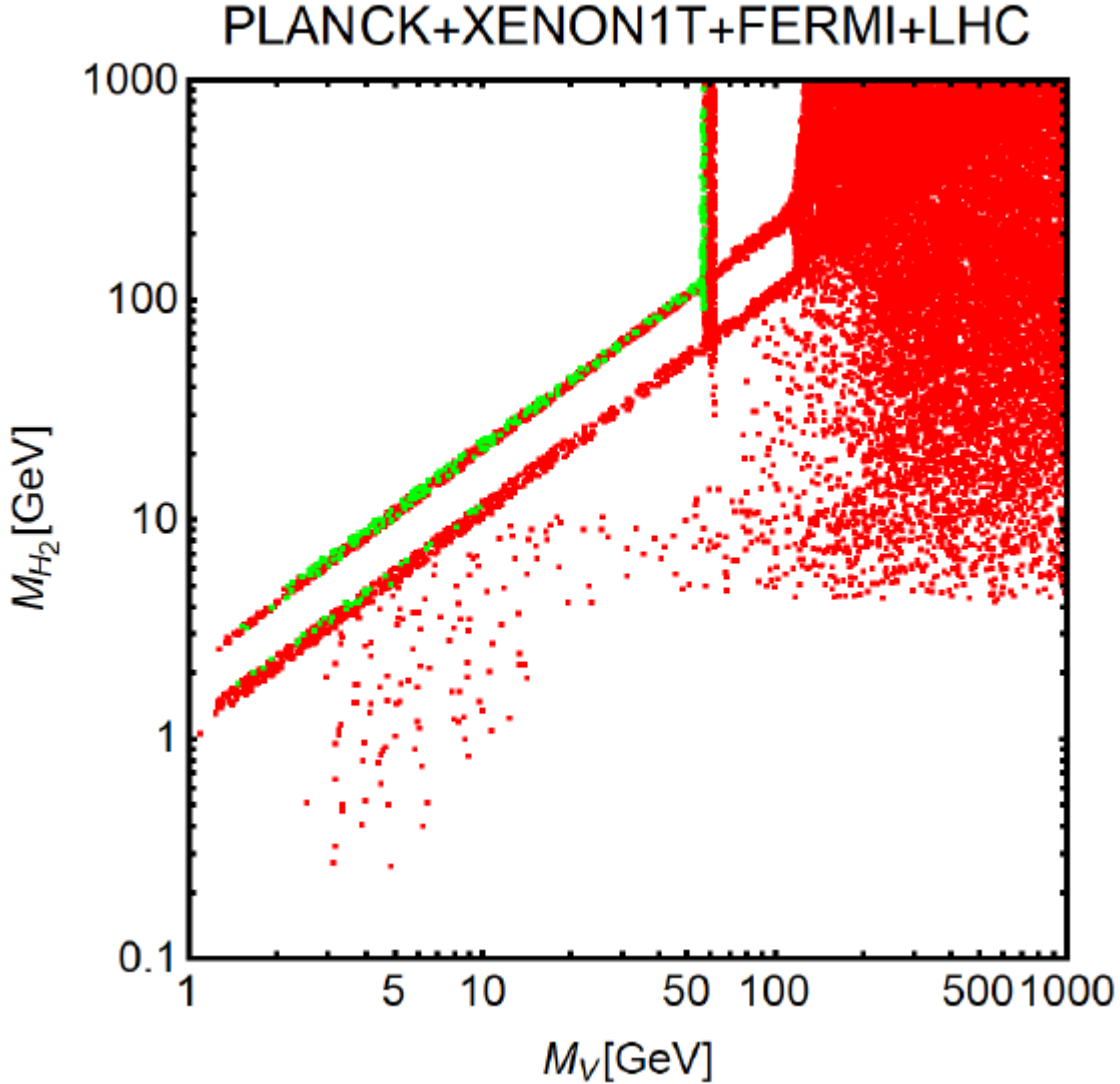
$$\Omega_{DM} h^2 \approx 8.76 \times 10^{-11} \text{ GeV}^{-2} \left[\int_{T_0}^{T_{f.o.}} g_*^{\frac{1}{2}} \langle \sigma v \rangle \frac{dT}{m_{DM}} \right]^{-1}$$

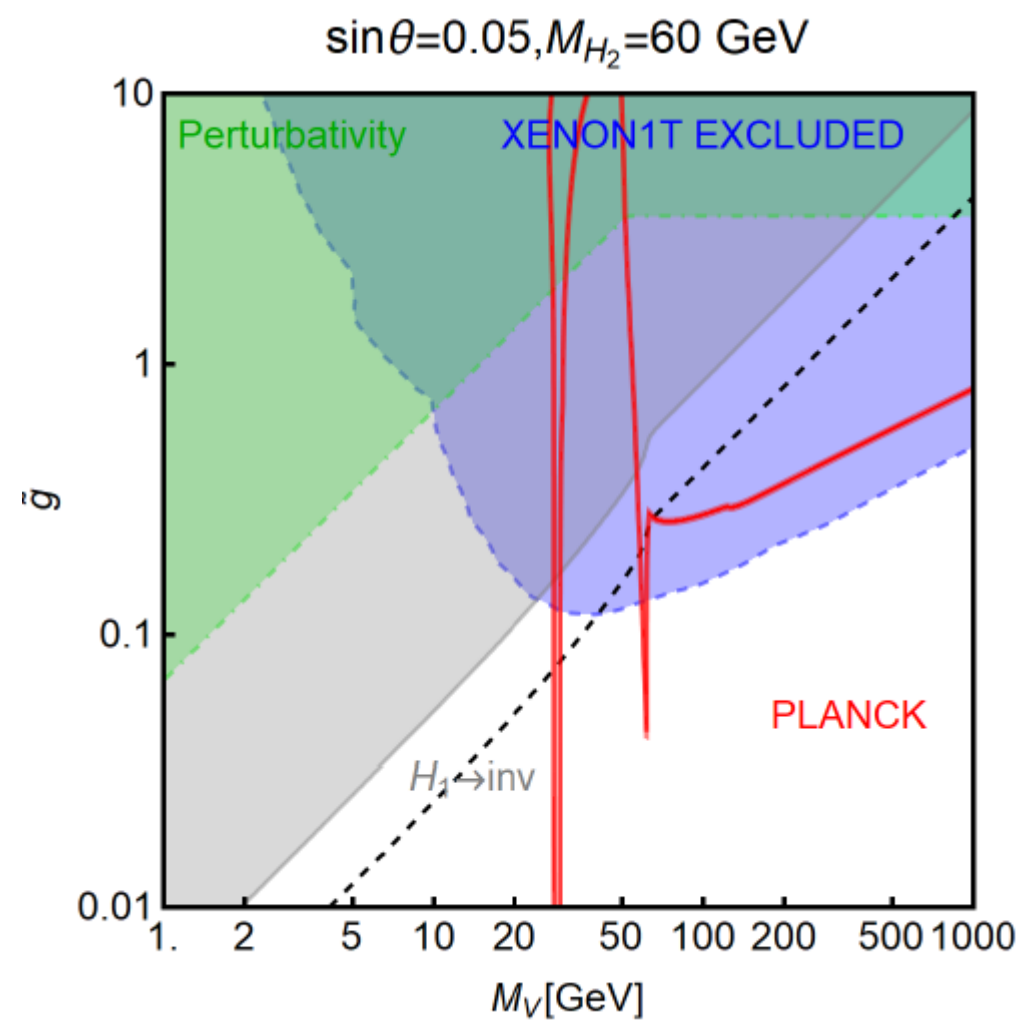
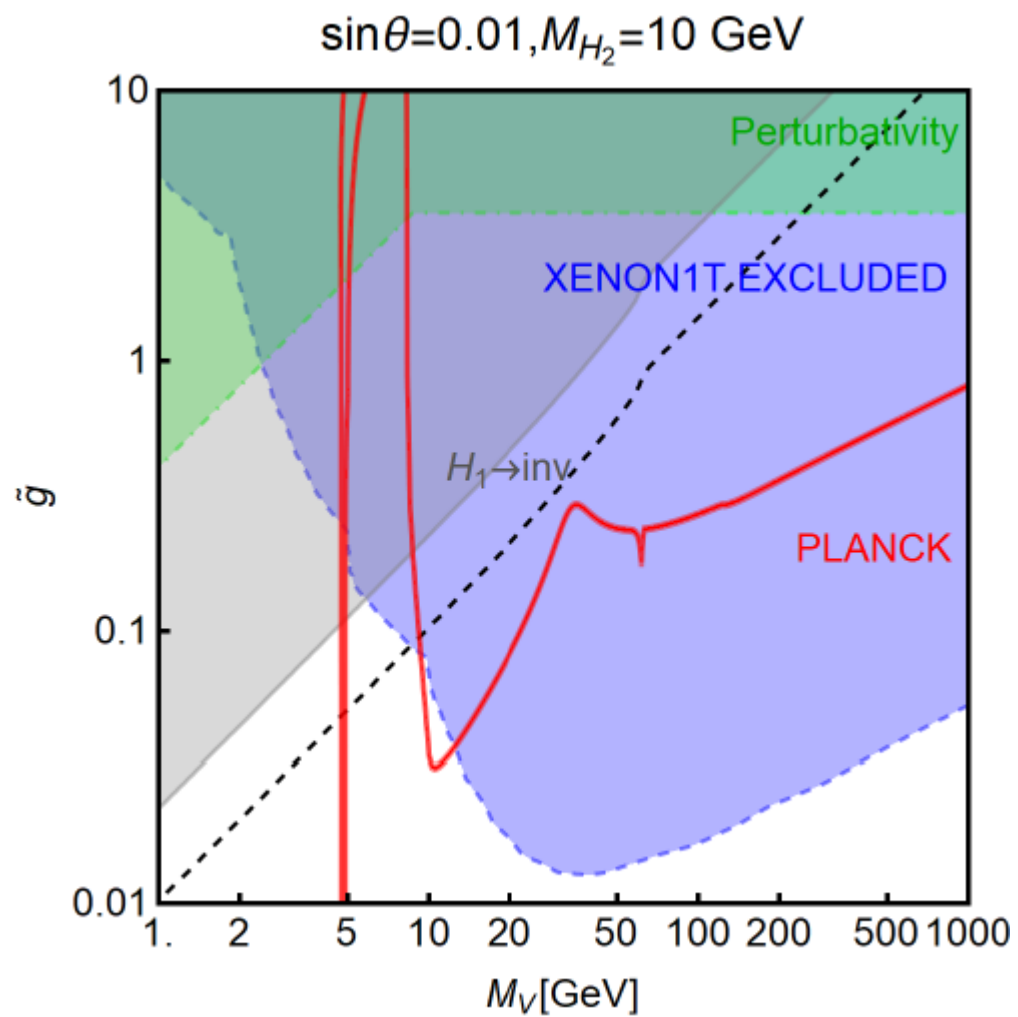
Relic density depends on a single particle physics input

In the 'decoupled' regime we cannot have viable relic density within the kinematical reach of invisible Higgs decays.

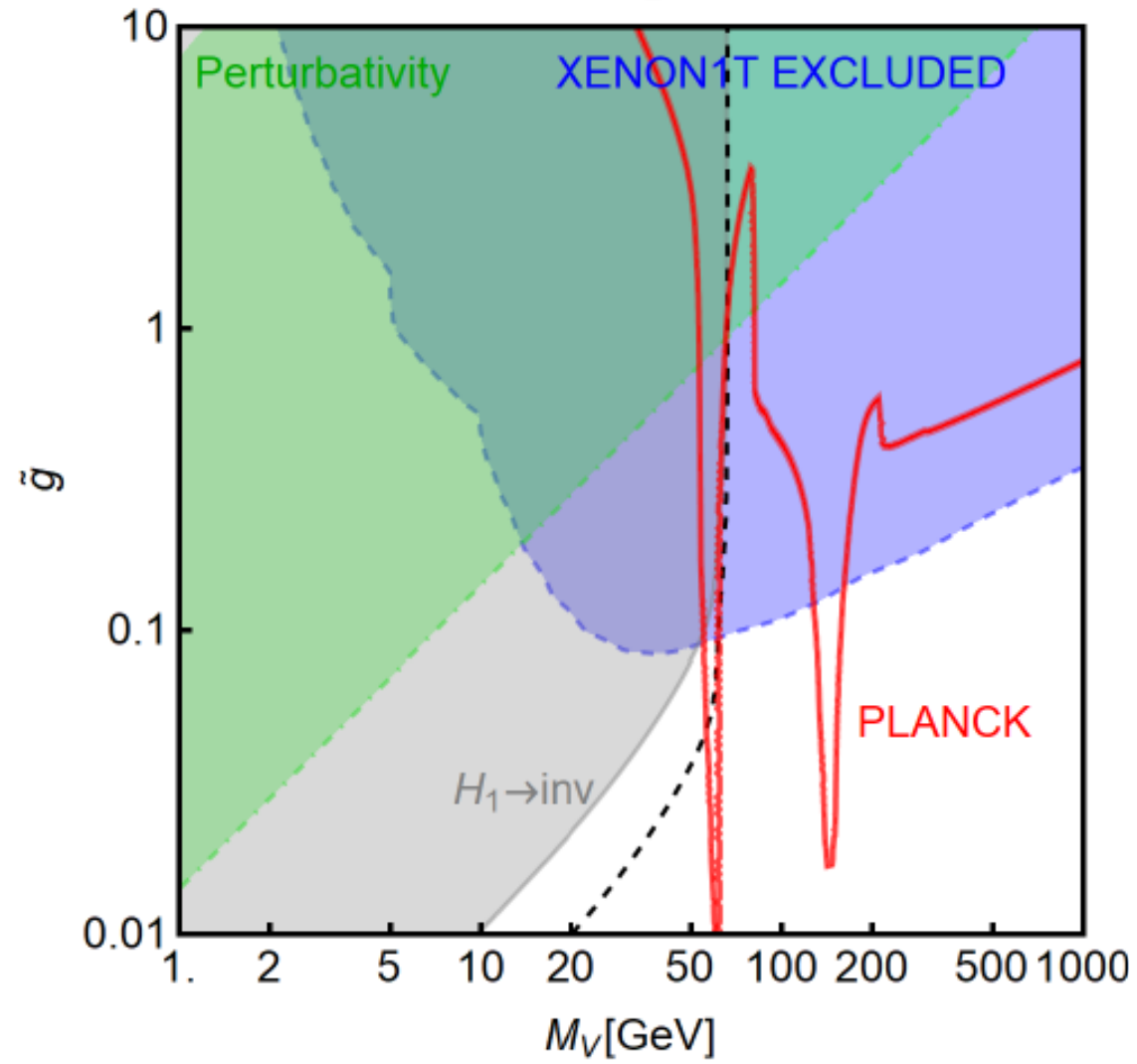


Light viable DM requires extra light degrees of freedom at the scale of the SM-like Higgs or below.

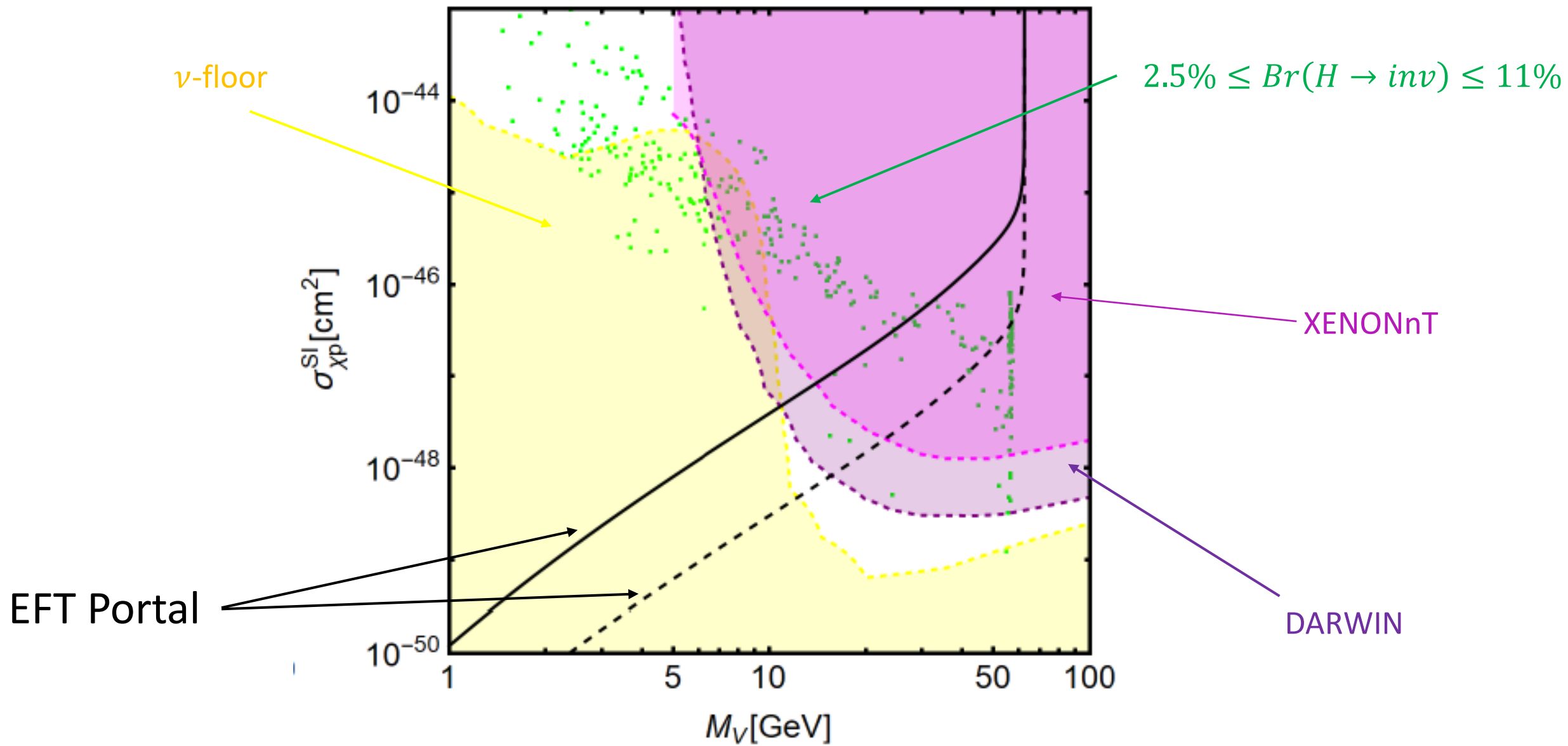




$\sin\theta=0.3, M_{H_2}=300 \text{ GeV}$



PLANCK+XENON1T+LHC+FERMI



Conclusions

The LHC vs DD correlation plot is a powerful tool to highlight the complementarity between different DM search strategies.

The Vector Higgs portal is a theoretically consistent benchmark for the interpretation of the experimental outcome.

However the correlation plot does not properly described candidates achieving the correct relic density through the conventional freeze-out.