

Quintessential Inflation with a Trap

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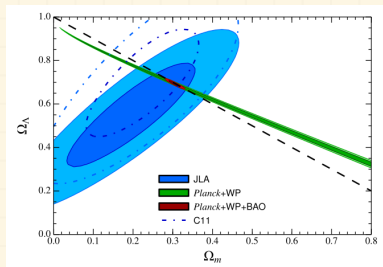
Dimopoulos, MK, Owen (2019)

Gonzalvo, MK, Rusak

Introduction

- ▶ Introduction
- ▶ Tachyonic trap
- ▶ A concrete example

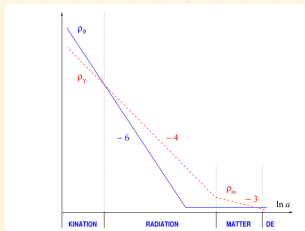
Cosmological Constant



Betoule+ (2014)

- ▶ Cosmological Constant Λ
 - ▶ ρ_λ is ~ 120 orders of magnitude smaller than m_{Pl}
 - ▶ The coincidence: $\Omega_\Lambda \simeq \Omega_m$

Quintessence

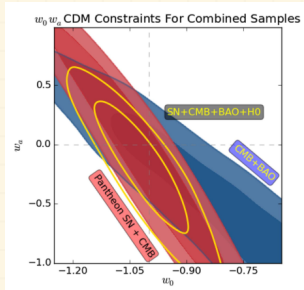


- ▶ Assume $\Lambda = 0$
- ▶ A dynamical mechanism: quintessence
 - ▶ Slowly rolling scalar field ϕ

$$\ddot{\phi} + 3H\dot{\phi} + V_{,\phi} = 0$$

- ▶ Reminiscent of inflation, but matter and radiation cannot be neglected

Quintessence



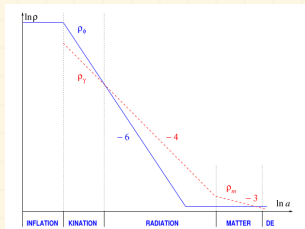
Scolnic+ (2018)

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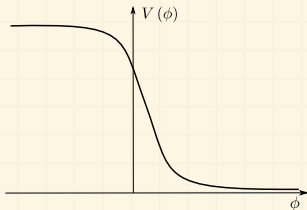
Quintessential Inflation



Borrowed from: Dimopoulos & Owen (2017)

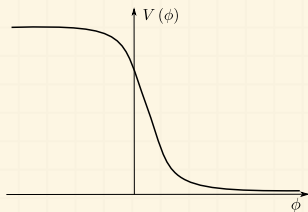
- ▶ Inflaton and quintessence are the same field
 - ▶ Economy: no new degrees of freedom
 - ▶ Quintessence initial conditions given by inflation
- ▶ Difficulty: a huge range of energies between ρ_{inf} and ρ_{Λ}
 - ▶ The potential has to have a huge gradient

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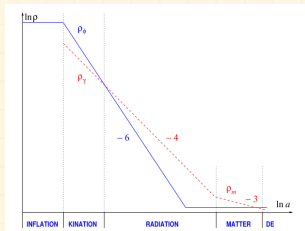
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Reheating



- ▶ Reheating
 - ▶ Non-oscillating part
- ▶ Other options
 - ▶ gravitational reheating, but it is very inefficient
 - ▶ instant preheating
 - ▶ curvaton preheating

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Other Problems

- ▶ The mass of the quintessence field $m_\phi \ll H_0$ very light, also [Carroll (1998)]

$$\beta_i \left(\frac{\phi}{M} \right) \mathcal{L}_i$$

- ▶ 5th force: Eötvös-type experiments constraint violations of equivalence principle
- ▶ Varying fundamental constants: e.g. fine structure constant α from $\beta_{F^2} \left(\frac{\phi}{M} \right) F_{\mu\nu} F^{\mu\nu}$
- ▶ The excursion of the scalar field $\Delta\phi > m_{\text{Pl}}$
 - ▶ Radiative corrections
 - ▶ Non-renormalisable terms

- ▶ Introduction
- ▶ Tachyonic trap
- ▶ A concrete example

The Trapping Mechanism

- Simple Model: $H = 0$ and no (classical) potential
[Kofman et al. (2004)]

$$\mathcal{L} = -\partial_\mu \phi \partial^\mu \phi - \partial_\mu \chi \partial^\mu \chi - \frac{1}{2} g^2 \chi^2 (\phi - \phi_{\text{ESP}})^2$$

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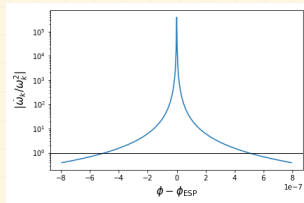
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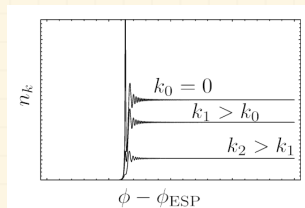
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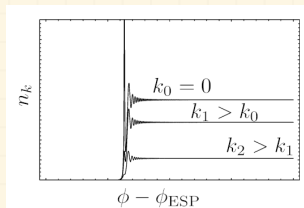
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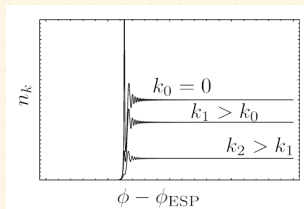
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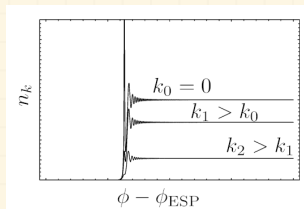
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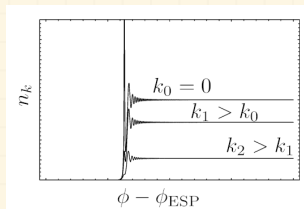
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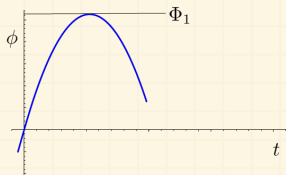
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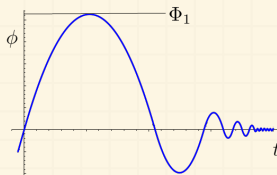
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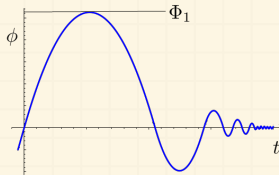
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- ▶ End of resonance

$$|\Delta\phi| < \sqrt{v/g} \Rightarrow \rho_\chi \simeq \rho_{\text{kin}}$$



Tachyonic Trap

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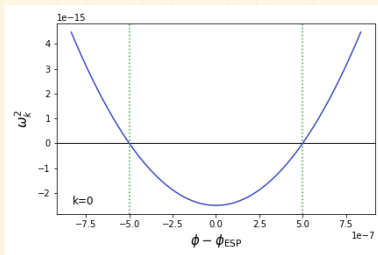
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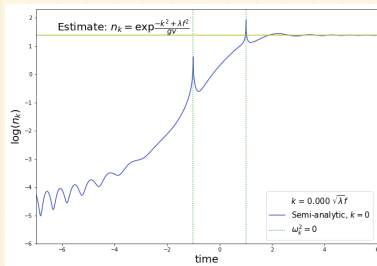
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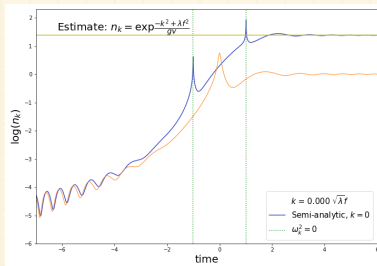
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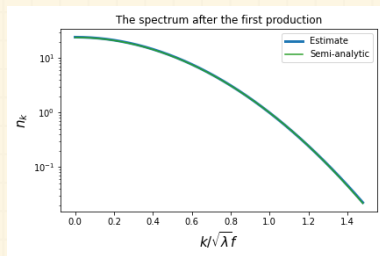
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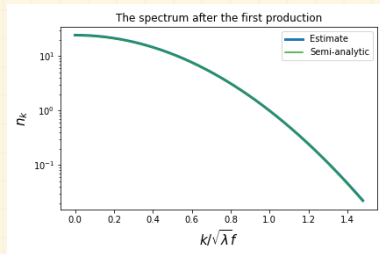
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$$n_{k1} = e^{\pi \frac{-k^2 + \lambda f^2}{g v}}$$

- ▶ Particle production stops when either

- ▶ $|\Delta\phi| < \text{non-adiabatic region};$
or
- ▶ χ becomes heavy



Prompt Trapping

- ▶ Field excursion $\phi \ll m_{\text{Pl}} \Rightarrow \phi_{\text{ESP}} \ll m_{\text{Pl}}$ and $\Phi_1 \ll m_{\text{Pl}}$

1. Trapping is prompt

$$H_{\text{ESP}} \simeq v \quad \Rightarrow \quad t_1 H_{\text{ESP}} \ll 1$$

2. Radiative corrections small
3. Non-renormalisable terms in the perturbative potential are small

- ▶ In the vacuum ϕ is heavy and fixed

$$m_\phi^2 = g^2 f^2 \gg H^2$$

1. No 5th force
2. No variation of fundamental constants

- ▶ Introduction
- ▶ Tachyonic trap
- ▶ A concrete example

Non-Perturbative Part

- Non-perturbative + perturbative part

$$\mathcal{L} = \mathcal{L}_{\text{kin}} + \mathcal{L}_V + \mathcal{L}_{\text{prt b}}$$

Non-Perturbative Part

► Kinetic term

$$\mathcal{L} = \frac{\alpha}{2} \left(\frac{m_{\text{Pl}}}{\phi} \right)^2 \partial_\mu \phi \partial^\mu \phi + \mathcal{L}_V + \mathcal{L}_{\text{prt b}}$$

Non-Perturbative Part

- The potential (e.g. gaugino condensation)

$$\mathcal{L} = \frac{\alpha}{2} \left(\frac{m_{\text{Pl}}}{\phi} \right)^2 \partial_\mu \phi \partial^\mu \phi + V_0 e^{-\kappa \phi / m_{\text{Pl}}} + \mathcal{L}_{\text{prtb}}$$

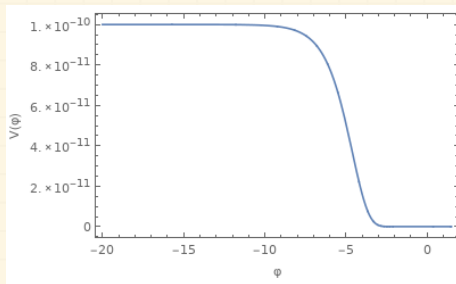
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$$\frac{\varphi}{m_{\text{Pl}}} \equiv \sqrt{\alpha} \ln \left(\frac{\phi}{m_{\text{Pl}}} \right)$$

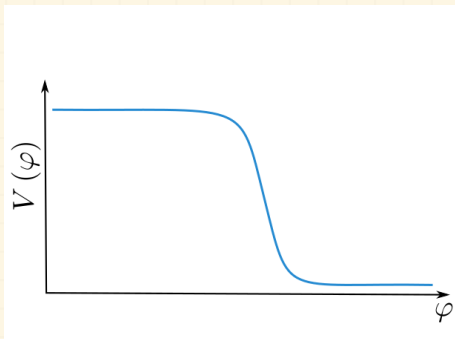
The pole is sent to $-\infty$



Perturbative Lagrangian

► Perturbative part

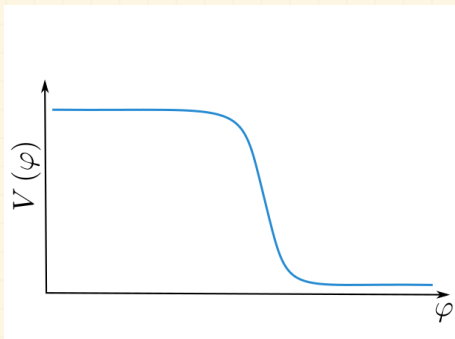
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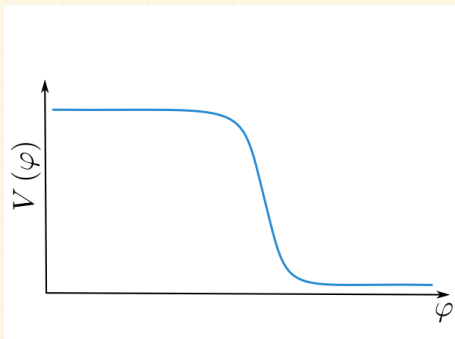
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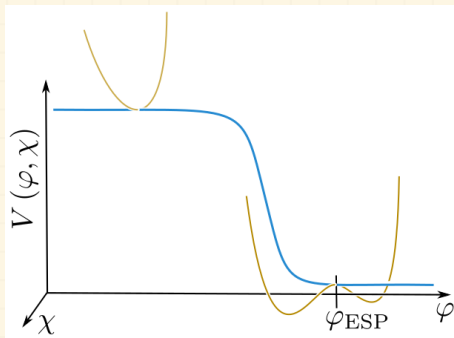
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- ▶ Peccei-Quinn field

- ▶ Economy: axion DM $f \sim 10^{12} \text{ GeV}$

- ▶ χ heavy during inflation $\Rightarrow$ no isocurvature perturbations

- ▶ κ : the vacuum energy

$$\kappa \frac{\phi_{\text{ESP}}}{m_{\text{Pl}}} = \ln \frac{V_0}{V_{\text{vac}}} \sim 10^2$$

Constraints

- ▶ Inflation
- ▶ Kination
- ▶ ESP
 - ▶ reheating $\Rightarrow N_*$ $\Rightarrow$ Planck Constraints

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$$m_{\text{Pl}} = 1$$

$$\varphi \equiv \sqrt{\alpha} \ln \phi$$

Inflation

- ▶ Double exponential potential $V(\varphi) = V_0 \exp[-\kappa \exp(\varphi/\sqrt{\alpha})]$
- ▶ Observables

$$n_s \simeq 1 - \frac{2}{N_*} - \frac{\alpha}{N_*^2}; \quad r \simeq \frac{8\alpha}{N_*^2}; \quad A_s = \frac{1}{24\pi^2} \frac{V}{\epsilon};$$

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- ▶ Observational (Planck) constraints

$$\begin{array}{ll} n_s = 0.968 \pm 0.006 & 53 \leq N_* < 70 \\ r < 0.06 & \Rightarrow 1 \leq \alpha < 37 \\ \ln(10^{10} A_s) = 3.094 \pm 0.034 & \kappa \phi_{\text{ESP}} < 236 \end{array}$$

Inflation

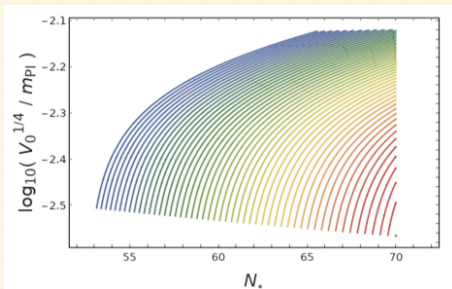
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- ▶ Inflationary energy scale (within 2σ of n_s)



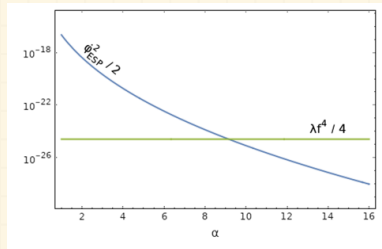
Kination

- Potential is irrelevant

$$\ddot{\phi} + 3H\dot{\phi} \simeq 0$$

$\Downarrow$

$$\dot{\phi}^2 \propto a^{-6}$$



$$\alpha < 10$$

- Velocity at ESP

$$\dot{\phi}_{\text{ESP}} = \frac{\sqrt{2V_0(\alpha)/3}}{\exp\left(\sqrt{\alpha/2}\right)} \left(\frac{\sqrt{2\alpha}}{\kappa(\alpha)\phi_{\text{ESP}}} \right)^{\sqrt{3\alpha/2}} < \sqrt{V(\varphi)}$$

At the enhanced symmetry point

- ▶ With $|\varphi - \varphi_{\text{ESP}}| \ll 1$

$$V \simeq V_{\text{vac}} \left(\frac{V_{\text{vac}}}{V_0} \right)^{(\varphi - \varphi_{\text{ESP}})/\sqrt{\alpha}} + \frac{1}{2} \gamma^2 (\varphi - \varphi_{\text{ESP}})^2 \chi^2 + V(\chi)$$

where

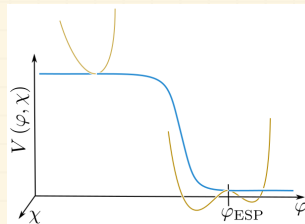
$$\gamma \equiv \frac{g\phi_{\text{ESP}}}{\sqrt{\alpha}}$$

- ▶ Dispersion relation of χ_k

$$\omega_k^2 = k^2 + \gamma^2 (\varphi - \varphi_{\text{ESP}})^2 + \lambda (3 \langle \chi^2 \rangle - f^2)$$

- ▶ Constraint

$$\Phi_1(\alpha, \gamma, \phi_{\text{ESP}}) < 1$$



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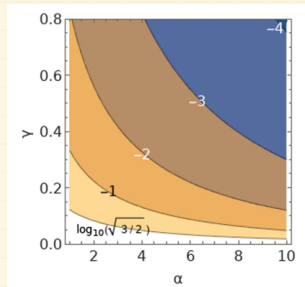
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$$\phi_{\text{ESP}} = 1$$

Trapping

- ▶ Particle production ends when

1. $|\Delta\varphi| < \text{non-adiabtic window} \quad \Rightarrow \quad \rho_r \simeq \frac{1}{2}\dot{\varphi}_{\text{ESP}}^2$

2. $m_\chi^2 \simeq \lambda (3\langle\chi^2\rangle - f^2)$ becomes large $\Rightarrow \quad \rho_r < \frac{1}{2}\dot{\varphi}_{\text{ESP}}^2$

- ▶ 1st case if

$$\frac{\lambda}{\gamma^2} < \frac{1}{7}$$

- ▶ Stronger χ 's self-coupling less effective trapping.

Reheating

- ▶ If $\frac{\lambda}{\gamma^2} < \frac{1}{7}$ is satisfied
- ▶ Reheating temperature (instant thermalisation)

$$\frac{1}{2}\dot{\phi}_{\text{ESP}}^2 \simeq \frac{\pi^2}{30}g_* T_{\text{reh}}^4$$

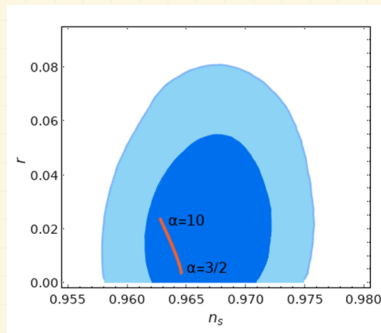
- ▶ An estimate

$$\begin{array}{rcl} \phi_{\text{ESP}} & \simeq & 1 \\ V_{\text{end}}^{1/4} & \sim & 10^{-2} \\ \kappa & \sim & 10^2 \\ & \Downarrow & \\ T_{\text{reh}} & \sim & 10^{11-13} \text{ GeV} \end{array}$$

Planck Constraints

- Number of e-folds of observable inflation

$$N_* \simeq 56.3 + \sqrt{\frac{\alpha}{24}} \ln \left(\frac{244}{\sqrt{2\alpha}} \phi_{\text{ESP}}^2 \right)$$



Conclusions

- ▶ Quintessential inflation
- ▶ Trapping mechanism helps solving several traditional problems
- ▶ A concrete example

Kinetic Term

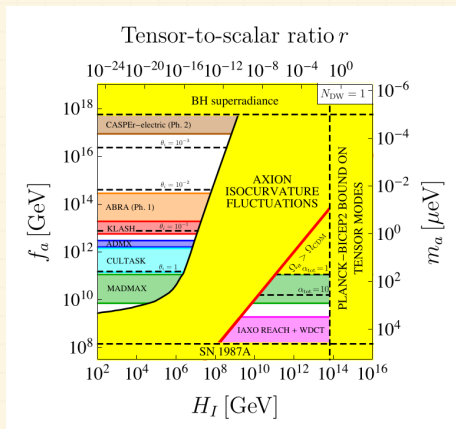
- Kähler Potential for a string modulus $T = \frac{\phi + i\sigma}{\sqrt{2}m_{\text{Pl}}}$

$$\begin{aligned}\frac{K}{m_{\text{Pl}}^2} &= -3 \ln (T + \bar{T}) \\ &= -3 \ln \left(\sqrt{2} \frac{\phi}{m_{\text{Pl}}} \right)\end{aligned}$$

This gives the kinetic term

$$\begin{aligned}\mathcal{L}_{\text{kin}} &= K_{T\bar{T}} \partial_\mu T \partial^\mu \bar{T} \\ &= \underbrace{\frac{3}{2}}_\alpha \left(\frac{m_{\text{Pl}}}{\phi} \right)^2 \left[\frac{1}{2} (\partial\phi)^2 + \frac{1}{2} (\partial\sigma)^2 \right]\end{aligned}$$

Axion DM Constraints



Di Luzio+ (2020)