

# The origin of primordial non-gaussianity

David H. Lyth

Particle Theory and Cosmology Group  
Physics Department  
Lancaster University

# Non-gaussianity as a discriminator

Model  $\Rightarrow$  Primordial Perturbations  $\Rightarrow$  Observables

# Non-gaussianity as a discriminator

Model  $\Rightarrow$  Primordial Perturbations  $\Rightarrow$  Observables

- From observation, curvature perturbation  $\zeta$  dominates
  - **How much others?** (tensor, mag. field, cosmic string . . .)

# Non-gaussianity as a discriminator

Model  $\Rightarrow$  Primordial Perturbations  $\Rightarrow$  Observables

- From observation, curvature perturbation  $\zeta$  dominates
  - **How much others?** (tensor, mag. field, cosmic string . . .)
- **What is spectral tilt of curvature perturbation?**

# Non-gaussianity as a discriminator

Model  $\Rightarrow$  Primordial Perturbations  $\Rightarrow$  Observables

- From observation, curvature perturbation  $\zeta$  dominates
  - **How much others?** (tensor, mag. field, cosmic string . . .)
- **What is spectral tilt of curvature perturbation?**
- Curvature perturbation almost gaussian ( $\zeta_k$  uncorrelated)
  - **How much non-gaussianity?**
  - **What kind (bispectrum, trispectrum)**

# Non-gaussianity as a discriminator

Model  $\Rightarrow$  Primordial Perturbations  $\Rightarrow$  Observables

- From observation, curvature perturbation  $\zeta$  dominates
  - How much others? (tensor, mag. field, cosmic string . . .)
- What is spectral tilt of curvature perturbation?
- Curvature perturbation almost gaussian ( $\zeta_k$  uncorrelated)
  - How much non-gaussianity?
  - What kind (bispectrum, trispectrum)
- **How may we compute the non-gaussianity?**

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.
2. Consider era  $T \gtrsim 10$  keV  $\Rightarrow$  separate universes

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.
2. Consider era  $T \gtrsim 10$  keV  $\Rightarrow$  separate universes
3. Work on slicing of uniform energy density  $\rho$

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.
2. Consider era  $T \gtrsim 10$  keV  $\Rightarrow$  separate universes
3. Work on slicing of uniform energy density  $\rho$
4. Define local scale factor:  $g_{ij} = \tilde{a}^2(\mathbf{x}, t)\delta_{ij}$

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.
2. Consider era  $T \gtrsim 10$  keV  $\Rightarrow$  separate universes
3. Work on slicing of uniform energy density  $\rho$
4. Define local scale factor:  $g_{ij} = \tilde{a}^2(\mathbf{x}, t)\delta_{ij}$
5. Define  $\zeta(\mathbf{x}, t) \equiv \ln \tilde{a}(\mathbf{x}, t) - \ln a(t) \equiv \delta N$

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.
  2. Consider era  $T \gtrsim 10$  keV  $\Rightarrow$  separate universes
  3. Work on slicing of uniform energy density  $\rho$
  4. Define local scale factor:  $g_{ij} = \tilde{a}^2(\mathbf{x}, t)\delta_{ij}$
  5. Define  $\zeta(\mathbf{x}, t) \equiv \ln \tilde{a}(\mathbf{x}, t) - \ln a(t) \equiv \delta N$
- 
- In words,  $\zeta(\mathbf{x}, t)$  is the perturbation in the number of *e*-folds of expansion, starting from a *flat* slice.

Starobinsky 92; Salopek & Bond 90; DHL, Malik & Sasaki 2005 (non-perturbative refs.)

# Defining the curvature perturbation $\zeta$

1. Smooth Universe on shortest cosmological scale  $10^{-2}$  Mpc.
  2. Consider era  $T \gtrsim 10$  keV  $\Rightarrow$  separate universes
  3. Work on slicing of uniform energy density  $\rho$
  4. Define local scale factor:  $g_{ij} = \tilde{a}^2(\mathbf{x}, t)\delta_{ij}$
  5. Define  $\zeta(\mathbf{x}, t) \equiv \ln \tilde{a}(\mathbf{x}, t) - \ln a(t) \equiv \delta N$
- 
- In words,  $\zeta(\mathbf{x}, t)$  is the perturbation in the number of *e-folds* of expansion, starting from a *flat* slice.
- Starobinsky 92; Salopek & Bond 90; DHL, Malik & Sasaki 2005 (non-perturbative refs.)
- 
- Constant value  $\zeta(\mathbf{x})$  at  $T \sim 10$  keV provides the initial condition for adiabatic perturbations.

# The correlators

Spectrum  $\mathcal{P}_\zeta$ , bispectrum $^\dagger f_{\text{NL}}$ , trispectrum $^{\dagger\dagger} \tau_{\text{NL}}$ :

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}') K_1 \mathcal{P}_\zeta$$

$$\frac{5}{3} \langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \zeta_{\mathbf{k}''} \rangle = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}' + \mathbf{k}'') K_2 \mathcal{P}_\zeta^2 f_{\text{NL}}$$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \zeta_{\mathbf{k}''} \zeta_{\mathbf{k}'''}\rangle_c = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}' + \mathbf{k}'' + \mathbf{k'''}) K_3 \mathcal{P}_\zeta^3 \tau_{\text{NL}}$$

# The correlators

Spectrum  $\mathcal{P}_\zeta$ , bispectrum $^\dagger f_{\text{NL}}$ , trispectrum $^{\dagger\dagger} \tau_{\text{NL}}$ :

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}') K_1 \mathcal{P}_\zeta$$

$$\frac{5}{3} \langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \zeta_{\mathbf{k}''} \rangle = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}' + \mathbf{k}'') K_2 \mathcal{P}_\zeta^2 f_{\text{NL}}$$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \zeta_{\mathbf{k}''} \zeta_{\mathbf{k}'''}\rangle_c = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}' + \mathbf{k}'' + \mathbf{k'''}) K_3 \mathcal{P}_\zeta^3 \tau_{\text{NL}}$$

where the kinematic factors depend on the wave-vectors:

$$K_1 \equiv 2\pi^2/k^3$$

$$K_2 \equiv K_1(k)K_1(k') + 5\text{perms}$$

$$K_3 \equiv K_2 K_1(|\mathbf{k} + \mathbf{k}''|) + 23\text{perms}$$

$\dagger$  Komatsu/Spergel 2000; Maldacena 2003     $\dagger\dagger$  Boubeker/DHL 2005

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)
- $\tau_{\text{NL}} \lesssim 10^4$  (WMAP)

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)
- $\tau_{\text{NL}} \lesssim 10^4$  (WMAP)
- Eventual bounds:  $|f_{\text{NL}}| \lesssim 1$  and  $|\tau_{\text{NL}}| \lesssim 300$

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)
- $\tau_{\text{NL}} \lesssim 10^4$  (WMAP)
- **Eventual bounds:**  $|f_{\text{NL}}| \lesssim 1$  and  $|\tau_{\text{NL}}| \lesssim 300$
- Or  $|f_{\text{NL}}| \sim 0.01$  from 21 cm anisotropy? (Cooray 06)

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)
- $\tau_{\text{NL}} \lesssim 10^4$  (WMAP)
- **Eventual bounds:**  $|f_{\text{NL}}| \lesssim 1$  and  $|\tau_{\text{NL}}| \lesssim 300$
- Or  $|f_{\text{NL}}| \sim 0.01$  from 21 cm anisotropy? (Cooray 06)
- Theory gives  $|f_{\text{NL}}| \sim 0.01$  (standard paradigm)

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)
- $\tau_{\text{NL}} \lesssim 10^4$  (WMAP)
- **Eventual bounds:**  $|f_{\text{NL}}| \lesssim 1$  and  $|\tau_{\text{NL}}| \lesssim 300$
- Or  $|f_{\text{NL}}| \sim 0.01$  from 21 cm anisotropy? (Cooray 06)
- Theory gives  $|f_{\text{NL}}| \sim 0.01$  (standard paradigm)
- Or  $|f_{\text{NL}}| \gtrsim 1$  (curvaton & inhomogeneous reheating paradigms)

# Observation

- $-54 < f_{\text{NL}} < 114$  (WMAP+SDSS)
- $\tau_{\text{NL}} \lesssim 10^4$  (WMAP)
- **Eventual bounds:**  $|f_{\text{NL}}| \lesssim 1$  and  $|\tau_{\text{NL}}| \lesssim 300$
- Or  $|f_{\text{NL}}| \sim 0.01$  from 21 cm anisotropy? (Cooray 06)
- Theory gives  $|f_{\text{NL}}| \sim 0.01$  (standard paradigm)
- Or  $|f_{\text{NL}}| \gtrsim 1$  (curvaton & inhomogeneous reheating paradigms)

# Inflationary origin of $\zeta$

All light fields get perturbation  $\sim H/2\pi$  from vacuum fluctuation

# Inflationary origin of $\zeta$

All light fields get perturbation  $\sim H/2\pi$  from vacuum fluctuation

- Invoke separate universe assumption
  - Local evolution is that of an unperturbed universe
  - Zeroth order gradient expansion plus local isotropy

# Inflationary origin of $\zeta$

All light fields get perturbation  $\sim H/2\pi$  from vacuum fluctuation

- Invoke separate universe assumption
  - Local evolution is that of an unperturbed universe
  - Zeroth order gradient expansion plus local isotropy
- Assume some light fields  $\phi_i(\mathbf{x}, t_1)$  define subsequent expansion  $N(\mathbf{x}, t)$ 
  - Choose  $c_s a_1 H_1/k \sim$  a few, so that that  $\delta\phi_i$  is classical

# Inflationary origin of $\zeta$

All light fields get perturbation  $\sim H/2\pi$  from vacuum fluctuation

- Invoke separate universe assumption
  - Local evolution is that of an unperturbed universe
  - Zeroth order gradient expansion plus local isotropy
- Assume some light fields  $\phi_i(\mathbf{x}, t_1)$  define subsequent expansion  $N(\mathbf{x}, t)$ 
  - Choose  $c_s a_1 H_1/k \sim \text{a few}$ , so that that  $\delta\phi_i$  is classical

$$\begin{aligned}\zeta(\mathbf{x}, t) &= N_i(\phi_i(\mathbf{x}), \rho(t)) - N(\phi_i, \rho(t)) \\ &= N_i(t)\delta\phi_i(\mathbf{x}) + \frac{1}{2}N_{ij}(t)\delta\phi_i(\mathbf{x})\delta\phi_j(\mathbf{x}) + \dots\end{aligned}$$

DHL, Malik & Sasaki 05; DHL & Rodriguez 05 (non-perturbative)

# Calculating $f_{\text{NL}}$

Simplest case

$$\zeta(\mathbf{x}) = b\phi(\mathbf{x}) + (\phi(\mathbf{x}) - \bar{\phi})^2$$

with  $\phi(\mathbf{x})$  gaussian and dropping  $-b\bar{\phi}$

# Calculating $f_{\text{NL}}$

Simplest case

$$\zeta(\mathbf{x}) = b\phi(\mathbf{x}) + (\phi(\mathbf{x}) - \bar{\phi})^2$$

with  $\phi(\mathbf{x})$  gaussian and dropping  $-b\bar{\phi}$

$$\begin{aligned}\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle &= b^2 \langle \phi_{\mathbf{k}} \phi_{\mathbf{k}'} \rangle \\ \langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle &= \frac{b^2}{(2\pi)^3} \langle \phi_{\mathbf{k}_1} \phi_{\mathbf{k}_2} \left[ (\phi - \bar{\phi})^2 \right]_{\mathbf{k}_3} \rangle + \text{cyclic} \\ &= \frac{b^2}{(2\pi)^3} \int \langle \phi_{\mathbf{k}_1} \phi_{\mathbf{k}_2} \phi_{\mathbf{q}} \phi_{\mathbf{k}_3 - \mathbf{q}} \rangle d^3 q + \text{cyclic} \\ &= \frac{b^2}{(2\pi)^3} \int \langle \phi_{\mathbf{k}_1} \phi_{\mathbf{q}} \rangle \langle \phi_{\mathbf{k}_2} \phi_{\mathbf{k}_3 - \mathbf{q}} \rangle d^3 q + 5 \text{ perms.}\end{aligned}$$

Gives  $\boxed{(3/5)f_{\text{NL}} = 1/b^2}$

# Loop correction

Include quadratic term of  $\zeta$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = b^2 \langle \phi_{\mathbf{k}} \phi_{\mathbf{k}'} \rangle + \langle [(\phi - \bar{\phi})^2]_{\mathbf{k}} [(\phi - \bar{\phi})^2]_{\mathbf{k}'} \rangle$$

# Loop correction

Include quadratic term of  $\zeta$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = b^2 \langle \phi_{\mathbf{k}} \phi_{\mathbf{k}'} \rangle + \langle [(\phi - \bar{\phi})^2]_{\mathbf{k}} [(\phi - \bar{\phi})^2]_{\mathbf{k}'} \rangle$$

Infrared divergent, use box size  $L$  then

$$\boxed{\mathcal{P}_\zeta = b^2 \mathcal{P}_\phi + 4 \mathcal{P}_\phi^2 \ln(kL)}$$

# Loop correction

Include quadratic term of  $\zeta$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = b^2 \langle \phi_{\mathbf{k}} \phi_{\mathbf{k}'} \rangle + \langle [(\phi - \bar{\phi})^2]_{\mathbf{k}} [(\phi - \bar{\phi})^2]_{\mathbf{k}'} \rangle$$

Infrared divergent, use box size  $L$  then

$$\boxed{\mathcal{P}_\zeta = b^2 \mathcal{P}_\phi + 4 \mathcal{P}_\phi^2 \ln(kL)}$$

Now go to box size  $M \ll L$

$$\begin{aligned}\zeta(\mathbf{x}) &= [b + 2(\bar{\phi}_M - \bar{\phi})] \phi(\mathbf{x}) + (\phi(\mathbf{x}) - \bar{\phi}_M)^2 \\ \mathcal{P}_{\zeta M} &= [b + 2(\bar{\phi}_M - \bar{\phi})]^2 \mathcal{P}_\phi + 4 \mathcal{P}_\phi^2 \ln(kM)\end{aligned}$$

# Loop correction

Include quadratic term of  $\zeta$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = b^2 \langle \phi_{\mathbf{k}} \phi_{\mathbf{k}'} \rangle + \langle [(\phi - \bar{\phi})^2]_{\mathbf{k}} [(\phi - \bar{\phi})^2]_{\mathbf{k}'} \rangle$$

Infrared divergent, use box size  $L$  then

$$\boxed{\mathcal{P}_\zeta = b^2 \mathcal{P}_\phi + 4 \mathcal{P}_\phi^2 \ln(kL)}$$

Now go to box size  $M \ll L$

$$\begin{aligned}\zeta(\mathbf{x}) &= [b + 2(\bar{\phi}_M - \bar{\phi})] \phi(\mathbf{x}) + (\phi(\mathbf{x}) - \bar{\phi}_M)^2 \\ \mathcal{P}_{\zeta M} &= [b + 2(\bar{\phi}_M - \bar{\phi})]^2 \mathcal{P}_\phi + 4 \mathcal{P}_\phi^2 \ln(kM)\end{aligned}$$

But  $\overline{(\bar{\phi}_M - \bar{\phi})^2} = \mathcal{P}_\phi \ln L/M$  hence  $\boxed{\overline{\mathcal{P}_{\zeta M}} = \mathcal{P}_\zeta}$

# The standard scenario

Only relevant light field is slow-roll inflaton  $\phi$

# The standard scenario

Only relevant light field is slow-roll inflaton  $\phi$

$$\zeta = \frac{\partial N}{\partial \phi} \delta \phi + \frac{1}{2} \frac{\partial^2 N}{\partial \phi^2} (\delta \phi)^2 + \dots$$

# The standard scenario

Only relevant light field is slow-roll inflaton  $\phi$

$$\zeta = \frac{\partial N}{\partial \phi} \delta\phi + \frac{1}{2} \frac{\partial^2 N}{\partial \phi^2} (\delta\phi)^2 + \dots$$

But  $3H\dot{\phi} = -V'(\phi)$  and  $3M_{\text{P}}^2 H^2 = V$  so

# The standard scenario

Only relevant light field is slow-roll inflaton  $\phi$

$$\zeta = \frac{\partial N}{\partial \phi} \delta\phi + \frac{1}{2} \frac{\partial^2 N}{\partial \phi^2} (\delta\phi)^2 + \dots$$

But  $3H\dot{\phi} = -V'(\phi)$  and  $3M_{\text{P}}^2 H^2 = V$  so

$$dN = -Hdt = -\frac{Hd\phi}{\dot{\phi}} = \frac{3H^2}{V'} d\phi$$

# The standard scenario

Only relevant light field is slow-roll inflaton  $\phi$

$$\zeta = \frac{\partial N}{\partial \phi} \delta \phi + \frac{1}{2} \frac{\partial^2 N}{\partial \phi^2} (\delta \phi)^2 + \dots$$

But  $3H\dot{\phi} = -V'(\phi)$  and  $3M_{\text{P}}^2 H^2 = V$  so

$$dN = -Hdt = -\frac{Hd\phi}{\dot{\phi}} = \frac{3H^2}{V'} d\phi$$

so  $\partial N / \partial \phi = M_{\text{P}}^{-2} V / V'$  giving

$$\zeta = \left( \frac{\partial N}{\partial \phi} \delta \phi \right) + \left( \epsilon - \frac{1}{2} \eta \right) \left( \frac{\partial N}{\partial \phi} \delta \phi \right)^2$$

where  $\epsilon \equiv (M_{\text{P}}^2/2)(V'/V)^2$  and  $\eta \equiv M_{\text{P}}^2 V''/V$

# Non-gaussianity of $\delta\phi$

If  $\delta\phi$  gaussian,  $\frac{3}{5}f_{\text{NL}} = \epsilon - \eta/2 \sim \pm 10^{-2}$

# Non-gaussianity of $\delta\phi$

If  $\delta\phi$  gaussian,  $\frac{3}{5}f_{\text{NL}} = \epsilon - \eta/2 \sim \pm 10^{-2}$

But non-gaussianity of  $\delta\phi$  gives extra contribution

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle_{\text{ngphi}} = (\partial N / \partial \phi)^3 \langle \delta\phi_{\mathbf{k}_1} \delta\phi_{\mathbf{k}_2} \delta\phi_{\mathbf{k}_3} \rangle$$

# Non-gaussianity of $\delta\phi$

If  $\delta\phi$  gaussian,  $\boxed{\frac{3}{5}f_{\text{NL}} = \epsilon - \eta/2 \sim \pm 10^{-2}}$

But non-gaussianity of  $\delta\phi$  gives extra contribution

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle_{\text{ngphi}} = (\partial N / \partial \phi)^3 \langle \delta\phi_{\mathbf{k}_1} \delta\phi_{\mathbf{k}_2} \delta\phi_{\mathbf{k}_3} \rangle$$

How to calculate  $\langle \delta\phi_{\mathbf{k}_1} \delta\phi_{\mathbf{k}_2} \delta\phi_{\mathbf{k}_3} \rangle$

- (i) Calculate action to at least third order in  $\delta\phi$
- (ii) Apply in-in Feynman rules

# Non-gaussianity of $\delta\phi$

If  $\delta\phi$  gaussian,  $\boxed{\frac{3}{5}f_{\text{NL}} = \epsilon - \eta/2 \sim \pm 10^{-2}}$

But non-gaussianity of  $\delta\phi$  gives extra contribution

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle_{\text{ngphi}} = (\partial N / \partial \phi)^3 \langle \delta\phi_{\mathbf{k}_1} \delta\phi_{\mathbf{k}_2} \delta\phi_{\mathbf{k}_3} \rangle$$

How to calculate  $\langle \delta\phi_{\mathbf{k}_1} \delta\phi_{\mathbf{k}_2} \delta\phi_{\mathbf{k}_3} \rangle$

(i) Calculate action to at least third order in  $\delta\phi$

(ii) Apply in-in Feynman rules

Keeping just tree level, total  $f_{\text{NL}}$  is [with  $0 < y < 5/6$ ]

$$\boxed{\frac{3}{5}f_{\text{NL}} = -\frac{1}{4} [n - 1 + (r/8)y(k_1, k_2, k_3)] \sim \pm 10^{-2}}$$

Maldacena 2003; Seery/Lidsey 2005

# The curvaton model

Mollerach 90; Linde/Mukhanov 96; DHL/Wands 01; Moroi/Takahashi 2001

- Negligible curvature perturbation generated during inflation
  - Just need  $|\dot{H}/H^2| \ll 1$ , inflation model irrelevant

# The curvaton model

Mollerach 90; Linde/Mukhanov 96; DHL/Wands 01; Moroi/Takahashi 2001

- Negligible curvature perturbation generated during inflation
  - Just need  $|\dot{H}/H^2| \ll 1$ , inflation model irrelevant
- Curvaton field  $\sigma$  light during inflation, value  $\sigma_*(\mathbf{x}, t)$ 
  - Near-gaussian perturbation  $\delta\sigma_*$ , spectrum  $(H/2\pi)^2$

# The curvaton model

Mollerach 90; Linde/Mukhanov 96; DHL/Wands 01; Moroi/Takahashi 2001

- Negligible curvature perturbation generated during inflation
  - Just need  $|\dot{H}/H^2| \ll 1$ , inflation model irrelevant
- Curvaton field  $\sigma$  light during inflation, value  $\sigma_*(\mathbf{x}, t)$ 
  - Near-gaussian perturbation  $\delta\sigma_*$ , spectrum  $(H/2\pi)^2$
- Curvaton oscillates when  $H \sim m_\sigma$  with amplitude  $\sigma_{\text{os}}(\sigma_*)$ 
  - Curvature perturbation still negligible

# The curvaton model

Mollerach 90; Linde/Mukhanov 96; DHL/Wands 01; Moroi/Takahashi 2001

- Negligible curvature perturbation generated during inflation
  - Just need  $|\dot{H}/H^2| \ll 1$ , inflation model irrelevant
- Curvaton field  $\sigma$  light during inflation, value  $\sigma_*(\mathbf{x}, t)$ 
  - Near-gaussian perturbation  $\delta\sigma_*$ , spectrum  $(H/2\pi)^2$
- Curvaton oscillates when  $H \sim m_\sigma$  with amplitude  $\sigma_{\text{os}}(\sigma_*)$ 
  - Curvature perturbation still negligible
- But  $\rho_\sigma(\mathbf{x}, t)/\rho_{\text{rad}}(t) \propto \tilde{a}(\mathbf{x}, t)$ 
  - Eventually generates  $\zeta = \delta N$ , then curvaton decays

# The curvaton model

Mollerach 90; Linde/Mukhanov 96; DHL/Wands 01; Moroi/Takahashi 2001

- Negligible curvature perturbation generated during inflation
  - Just need  $|\dot{H}/H^2| \ll 1$ , inflation model irrelevant
- Curvaton field  $\sigma$  light during inflation, value  $\sigma_*(\mathbf{x}, t)$ 
  - Near-gaussian perturbation  $\delta\sigma_*$ , spectrum  $(H/2\pi)^2$
- Curvaton oscillates when  $H \sim m_\sigma$  with amplitude  $\sigma_{\text{os}}(\sigma_*)$ 
  - Curvature perturbation still negligible
- But  $\rho_\sigma(\mathbf{x}, t)/\rho_{\text{rad}}(t) \propto \tilde{a}(\mathbf{x}, t)$ 
  - Eventually generates  $\zeta = \delta N$ , then curvaton decays
- Curvaton model is not an epicycal!
  - Candidate  $sRH_\nu$  discovered serendiptiously  
(Hamaguchi/Murayama/Yanagida 01)

# Prediction of curvaton model

$$\zeta = \frac{\partial N}{\partial \sigma} \delta \sigma + \frac{1}{2} \frac{\partial^2 N}{\partial \sigma^2} \delta \sigma^2$$

# Prediction of curvaton model

$$\zeta = \frac{\partial N}{\partial \sigma} \delta \sigma + \frac{1}{2} \frac{\partial^2 N}{\partial \sigma^2} \delta \sigma^2$$

- Allow evolution of curvaton,  $\sigma_{\text{OS}}(\sigma_*)$ 
  - At decay  $\rho_\sigma / \rho \equiv \Omega_\sigma$

# Prediction of curvaton model

$$\zeta = \frac{\partial N}{\partial \sigma} \delta \sigma + \frac{1}{2} \frac{\partial^2 N}{\partial \sigma^2} \delta \sigma^2$$

- Allow evolution of curvaton,  $\sigma_{\text{os}}(\sigma_*)$ 
  - At decay  $\rho_\sigma / \rho \equiv \Omega_\sigma$

$$\frac{3}{5} f_{\text{NL}} = \frac{3}{4\Omega_\sigma} \left( 1 + \frac{\sigma_{\text{os}} \sigma''_{\text{os}}}{(\sigma'_{\text{os}})^2} \right) - 1 - \frac{1}{2} \Omega_\sigma$$

# Prediction of curvaton model

$$\zeta = \frac{\partial N}{\partial \sigma} \delta \sigma + \frac{1}{2} \frac{\partial^2 N}{\partial \sigma^2} \delta \sigma^2$$

- Allow evolution of curvaton,  $\sigma_{\text{os}}(\sigma_*)$ 
  - At decay  $\rho_\sigma/\rho \equiv \Omega_\sigma$

$$\frac{3}{5} f_{\text{NL}} = \frac{3}{4\Omega_\sigma} \left( 1 + \frac{\sigma_{\text{os}} \sigma''_{\text{os}}}{(\sigma'_{\text{os}})^2} \right) - 1 - \frac{1}{2}\Omega_\sigma$$

- Also from cosmological perturbation theory

DHL/Ungarelli/Wands 02; Bartolo/Mataresse/Riotto 03

# Prediction of curvaton model

$$\zeta = \frac{\partial N}{\partial \sigma} \delta \sigma + \frac{1}{2} \frac{\partial^2 N}{\partial \sigma^2} \delta \sigma^2$$

- Allow evolution of curvaton,  $\sigma_{\text{os}}(\sigma_*)$ 
  - At decay  $\rho_\sigma/\rho \equiv \Omega_\sigma$

$$\frac{3}{5} f_{\text{NL}} = \frac{3}{4\Omega_\sigma} \left( 1 + \frac{\sigma_{\text{os}} \sigma''_{\text{os}}}{(\sigma'_{\text{os}})^2} \right) - 1 - \frac{1}{2}\Omega_\sigma$$

- Also from cosmological perturbation theory

DHL/Ungarelli/Wands 02; Bartolo/Mataresse/Riotto 03

- If  $\sigma_{\text{os}}$  linear and  $\Omega_\sigma \ll 1$  need  $\Omega_\sigma \gtrsim 10^{-2}$

# Prediction of curvaton model

$$\zeta = \frac{\partial N}{\partial \sigma} \delta \sigma + \frac{1}{2} \frac{\partial^2 N}{\partial \sigma^2} \delta \sigma^2$$

- Allow evolution of curvaton,  $\sigma_{\text{os}}(\sigma_*)$ 
  - At decay  $\rho_\sigma/\rho \equiv \Omega_\sigma$

$$\frac{3}{5} f_{\text{NL}} = \frac{3}{4\Omega_\sigma} \left( 1 + \frac{\sigma_{\text{os}} \sigma''_{\text{os}}}{(\sigma'_{\text{os}})^2} \right) - 1 - \frac{1}{2}\Omega_\sigma$$

- Also from cosmological perturbation theory

DHL/Ungarelli/Wands 02; Bartolo/Matarrese/Riotto 03

- If  $\sigma_{\text{os}}$  linear and  $\Omega_\sigma \ll 1$  need  $\Omega_\sigma \gtrsim 10^{-2}$
- If  $\sigma_{\text{os}}$  linear and  $\Omega_\sigma \simeq 1$  then  $f_{\text{NL}} = -\frac{4}{5}$

# My main messages

## To particle theorists

- User-friendly formula for primordial non-gaussianity

# My main messages

## To particle theorists

- User-friendly formula for primordial non-gaussianity

## To astronomers

- Go for  $|f_{\text{NL}}| < 1$  (using 21 cm anisotropy)?

# My main messages

## To particle theorists

- User-friendly formula for primordial non-gaussianity

## To astronomers

- Go for  $|f_{\text{NL}}| < 1$  (using 21 cm anisotropy)?

## To all

- Prediction generally depends on mean values of light scalar fields in observable Universe

# My main messages

## To particle theorists

- User-friendly formula for primordial non-gaussianity

## To astronomers

- Go for  $|f_{\text{NL}}| < 1$  (using 21 cm anisotropy)?

## To all

- Prediction generally depends on mean values of light scalar fields in observable Universe
  - Anthropic discussion mandatory

# My main messages

## To particle theorists

- User-friendly formula for primordial non-gaussianity

## To astronomers

- Go for  $|f_{\text{NL}}| < 1$  (using 21 cm anisotropy)?

## To all

- Prediction generally depends on mean values of light scalar fields in observable Universe
  - Anthropic discussion mandatory
    - Exception ??: single-component inflation